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ANALYTIC INTEGRABILITY OF BIANCHI CLASS A COSMOLOGICAL MODELS WITH k = 1

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ABSTRACT. We complete the study of the analytic integrability of the Class A of Bianchi cosmological models with k = 1, characterizing the analytic first integrals of the Bianchi types VI₀ and VIII₀.

1. INTRODUCTION AND STATEMENT OF THE MAIN RESULT

A cosmological model describes the universe and is defined by the space time geometry (determined by a metric); the presence of the matter and its physical behaviour (determined by the energy-momentum tensor); and the interaction between geometry and matter (described through the Einstein's equations).

Friedmann in 1922 introduced the study of homogeneous cosmologies. For spatially homogeneous cosmologies the Einstein field equations can be written as an autonomous system of first order differential equations, see [19] and references therein.

Bianchi models describe space-times which are foliated by homogeneous hypersurfaces of constant time. Homogeneity requires a three dimensional isometry group (and so a three dimensional Lie algebra). Bianchi [2, 3] was the first to solve the problem of classifying three dimensional Lie algebras which are non-isomorphic. The classification is determined by the dimension n of the algebra. There are nine types of models according to n:

(a)
$$n = 0$$
: type I;

(b) n = 1: types II, III;

- (c) n = 2: types IV, V, VI, VII;
- (d) n = 3: types VIII, IX.

The types I,V and IX contain as special cases the Friedmann-Robertson-Walker universes, see [16].

If we consider X_1, X_2, X_3 an appropriate basis of the three dimensional Lie algebra, then the classification depends on a scalar $a \in \mathbb{R}$ and a vector (n_1, n_2, n_3) , with

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 $n_i \in \{-1, 0, 1\}$, such that

$$[X_1, X_2] = n_3 X_3, \quad [X_2, X_3] = n_1 X_1 - a X_2, \quad [X_3, X_1] = n_2 X_2 + a X_1,$$

where [,] is the Lie bracket. In particular for a = 0 we obtain the models of Class A and for $a \neq 0$ we obtain the models of Class B. For more details see Bogoyavlensky [4]. The Bianchi models of Class A are detailed in Table 1.

Type	Ι	Π	VI ₀	VII ₀	VIII	IX
a	0	0	0	0	0	0
n_1	0	1	1	1	1	1
n_2	0	0	-1	1	1	1
n_3	0	0	0	0	-1	1

TABLE 1. The classification of Bianchi Class A cosmologies.

Bianchi models as dynamical systems were first studied in 1971 by Collins [5]. A systematic study is due to Bogoyavlensky and collaborators [4]. According to Bogoyavlensky's work, the dynamics of Class A types can be reduced into a six dimensional autonomous differential system defined on a compact manifold with a boundary. Since then we find in the literature a big amount of works studying the dynamical behaviour of these cosmologies.

It is worth to mention that the most contraversory model seems to be the Bianchi IX model because of its chaotic (in some sense) behaviour. Many authors studied the integrability of this model. We must mention that the notion of integrability can have different meanings. Cushman and Sniatycki [6] proved that Bianchi IX is a locally integrable Hamiltonian system. Latifi, Musette and Conte [7] proved that this model is not integrable in the sense that it does not satisfy the Painlevé property. Morales-Ruiz and Ramis [15], using techniques of differential Galois theory, showed that the Bianchi IX model (as a complex Hamiltonian system) is not completely integrable (in the Liouville sense) with rational first integrals. Llibre and Valls in [8], using techniques of formal series, proved that any analytic first integral of Bianchi IX is a function of its Hamiltonian. The same authors in [9] proved that Bianchi VIII does not admit an analytic first integral different from the Hamiltonian using the Darboux theory of integrability. Maciejewski, Strelcyn and Szydłowski in [12] applied Morales-Ramis approach (and also Kovacic algorithm for a reduced system) to prove the non-meromorphic integrability of Bianchi VIII.

In [14] Class A was studied reducing the dimension of the system. In [17] the authors studied the vacuum Bianchi models of Class A, and using the zero level set of the Hamiltonian they reduced the dimension of the dynamical system (preserving the polynomial structure). In [18] there is a study about the algebraic integrability of Class A Bianchi models in vacuum and with dust and radiative matter.

The Bianchi Class A cosmological models that we consider here are defined by the Hamiltonian

$$H = \frac{1}{(q_1 q_2 q_3)^{(1-k)/2}} \ \mathcal{H},$$

where

$$\mathcal{H} = -(p_1^2 q_1^2 + p_2^2 q_2^2 + p_3^2 q_3^2) + 2(p_1 p_2 q_1 q_2 + p_1 p_3 q_1 q_3 + p_2 p_3 q_2 q_3) + \frac{1}{2}(n_1 n_2 q_1 q_2 + n_1 n_3 q_1 q_3 + n_2 n_3 q_2 q_3) - \frac{1}{4}(n_1^2 q_1^2 + n_2^2 q_2^2 + n_3^2 q_3^2).$$
(1)

Here we consider the case k = 1 also analyzed in the works of Maciejewski and Szydłowski [13, 14], and of Llibre and Valls [8, 9]. See the first two papers for a good physical introduction to these systems.

The differential system associated to \mathcal{H} is $\dot{q}_i = \mathcal{H}_{p_i}$, $\dot{p}_i = -\mathcal{H}_{q_i}$, for i = 1, 2, 3, which writes

$$\begin{split} \dot{q}_{1} &= 2q_{1}(-p_{1}q_{1} + p_{2}q_{2} + p_{3}q_{3}), \\ \dot{q}_{2} &= 2q_{2}(p_{1}q_{1} - p_{2}q_{2} + p_{3}q_{3}), \\ \dot{q}_{3} &= 2q_{3}(p_{1}q_{1} + p_{2}q_{2} - p_{3}q_{3}), \\ \dot{p}_{1} &= -\frac{n_{1}}{2}(-n_{1}q_{1} + n_{2}q_{2} + n_{3}q_{3}) - 2p_{1}(-p_{1}q_{1} + p_{2}q_{2} + p_{3}q_{3}), \\ \dot{p}_{2} &= -\frac{n_{2}}{2}(n_{1}q_{1} - n_{2}q_{2} + n_{3}q_{3}) - 2p_{2}(p_{1}q_{1} - p_{2}q_{2} + p_{3}q_{3}), \\ \dot{p}_{3} &= -\frac{n_{3}}{2}(n_{1}q_{1} + n_{2}q_{2} - n_{3}q_{3}) - 2p_{3}(p_{1}q_{1} + p_{2}q_{2} - p_{3}q_{3}). \end{split}$$

$$(2)$$

System (2) is Hamiltonian with three degrees of freedom. Due to its Hamiltonian structure it is *completely integrable* if it admits three independent first integrals in involution, see for more details [1].

Under the two changes of variables $q_i \to x_i$, $p_i \to x_{i+3}/x_i$, for i = 1, 2, 3, first, and $x_i \to 6x_i$, for $i = 1, 2, 3, t \to t/2$, afterwards, system (2) becomes

$$\dot{x}_{1} = x_{1}(-x_{4} + x_{5} + x_{6}),
\dot{x}_{2} = x_{2}(x_{4} - x_{5} + x_{6}),
\dot{x}_{3} = x_{3}(x_{4} + x_{5} - x_{6}),
\dot{x}_{4} = 9n_{1}x_{1}(n_{1}x_{1} - n_{2}x_{2} - n_{3}x_{3}),
\dot{x}_{5} = 9n_{2}x_{2}(-n_{1}x_{1} + n_{2}x_{2} - n_{3}x_{3}),
\dot{x}_{6} = 9n_{3}x_{3}(-n_{1}x_{1} - n_{2}x_{2} + n_{3}x_{3}).$$
(3)

This system has lost the Hamiltonian structure. The first integral corresponding to system (3) coming from (1) is

$$\mathcal{H}^* = -9(n_1^2 x_1^2 + n_2^2 x_2^2 + n_3^2 x_3^2) - (x_4^2 - x_5^2 - x_6^2) + 18(n_1 n_2 x_1 x_2 + n_1 n_3 x_1 x_3 + n_2 n_3 x_2 x_3) + 2(x_4 x_5 + x_4 x_6 + x_5 x_6).$$
(4)

Our aim in this work is to complete the characterization of the analytic integrability of all Bianchi cosmologies of Class A with k = 1 described by systems (2) or (3).

Theorem 1. The following statements hold for system (3).

- (a) The Bianchi type I model has three independent polynomial first integrals, namely x_4, x_5 and x_6 .
- (b) The Bianchi type II model has three independent polynomial first integrals, namely x_5, x_6 and $\mathcal{H}^* = -9x_1^2 - x_4(x_4 - 2x_5 - 2x_6)$.
- (c) The Bianchi types VI_0 and VII_0 models have the polynomial first integrals x_6 and $\mathcal{H}^* = -9(x_1 - n_2x_2)^2 - (x_4^2 + x_5^2 + x_6^2) + 2(x_4x_5 + x_4x_6 + x_5x_6)$. These models do not admit any additional analytic first integral independent from the two previous ones.
- (d) The Bianchi type VIII and IX models have the polynomial first integral $\mathcal{H}^* = -9(x_1^2 + x_2^2 + n_3^2 x_3^2) (x_4^2 + x_5^2 + x_6^2) + 18(x_1 x_2 + n_3 x_1 x_3 + n_3 x_2 x_3) + 2(x_4 x_5 + x_4 x_6 + x_5 x_6)$. These models do not admit any additional analytic first integral independent from the previous one.

Due to the Theorem of Liouville-Arnold (see for details [1]) and since system (2) is Hamiltonian, we only need three independent first integrals in involution for describing their dynamics. We provide the independent analytical first integrals of the equivalent system (3), from these we can compute the corresponding ones of system (2). Of course, systems (2) and (3) can have more than three independent first integrals.

Statements (a) and (b) of Theorem 1 are easy to check by direct computation. In fact the Bianchi type I model has the two additional non-polynomial independent first integrals

$$x_1^{x_4-x_5+x_6}x_2^{x_4-x_5-x_6}, \quad x_1^{x_4+x_5-x_6}x_3^{x_4-x_5-x_6};$$

and the Bianchi type II model has the two additional non-polynomial independent first integrals

$$x_1^{\sqrt{A}-2x_6}x_2^{\sqrt{A}}(A+\sqrt{A}(-x_4+x_5+x_6))^{2x_6}, \quad x_1^{\sqrt{A}-2x_5}x_3^{\sqrt{A}}(A+\sqrt{A}(-x_4+x_5+x_6))^{2x_5},$$

where $A = (x_5 + x_6)^2 - \mathcal{H}^*$.

Statement (d) of Theorem 1 is proved in [8, 9].

The main result of this paper is statement (c) of Theorem 1 on the Bianchi types VI_0 and VII_0 models. Related with the integrability of these two systems there are some related results when the dynamics of these systems are reduced to dimension three, see [14, 10].

2. Proof of statement (c) of Theorem 1

System (3) for Bianchi types VI_0 and VII_0 has the form

$$\dot{x}_{1} = x_{1}(-x_{4} + x_{5} + x_{6}),
\dot{x}_{2} = x_{2}(x_{4} - x_{5} + x_{6}),
\dot{x}_{3} = x_{3}(x_{4} + x_{5} - x_{6}),
\dot{x}_{4} = 9x_{1}(x_{1} - n_{2}x_{2}),
\dot{x}_{5} = -9n_{2}x_{2}(x_{1} - n_{2}x_{2}),
\dot{x}_{6} = 0,$$
(5)

with $n_2 = \pm 1$ depending on the Bianchi type.

The following result is well known, see for instance [11].

Proposition 2. Let F be an analytic function and let $F = \sum_i F_i$ be its decomposition into homogeneous polynomials of degree i. Then F is an analytic first integral of the homogeneous differential system (3) if and only if for all i F_i is a homogeneous polynomial first integral of system (3).

Proposition 2 shows that for studying the analytic first integrals of the homogeneous differential system (3) it is sufficient to study its homogeneous polynomial first integrals.

We assume from now on that $n_1 = 1, n_2 = \pm 1, n_3 = 0$; this case corresponds to Bianchi types VI₀ and VII₀, see Table 1. Proposition 3 shows the existence of two independent polynomial first integrals and Theorem 4 assures that there are no more polynomial first integrals independent from the ones provided by Proposition 3.

Proposition 3. Two independent polynomial first integrals for the Bianchi types VI_0 and VII_0 models are x_6 and

$$H = -9(x_1 - n_2 x_2)^2 - (x_4^2 + x_5^2 + x_6^2) + 2(x_4 x_5 + x_4 x_6 + x_5 x_6).$$
(6)

Proof. It follows easily by direct computations.

Theorem 4. Let $h(x_1, \ldots, x_6)$ be a polynomial first integral of system (5). Then h is a polynomial in the variables x_6 and H, where H is given by (6).

We state a lemma before proving Theorem 4. See [9] for a proof.

Lemma 5. Let x_k be one dimensional variables for k = 1, ..., n and n > 1. Let $f = f(x_1, ..., x_n)$ be a polynomial and let $f_1 = f(x_1, ..., x_n)|_{x_l=c_0}$, for some $l \in \{1, ..., n\}$, where c_0 is a constant. Then there exists a polynomial $g = g(x_1, ..., x_n)$ such that $f = f_1 + (x_l - c_0)g$.

Proof of Theorem 4. Let $h_1 = h|_{x_3=0}$. We consider the involution σ defined by $\sigma(x_1) = n_2 x_2$, $\sigma(x_2) = n_2 x_1$, $\sigma(x_3) = -x_3$, $\sigma(x_4) = x_5$, $\sigma(x_5) = x_4$ and $\sigma(x_6) = x_6$. We note that h_1 is invariant under the involution σ , namely $\sigma(h_1) = h_1$. By Lemma

5 there exists a polynomial g such that $h = h_1 + x_3 g$. System (5) on $x_3 = 0$ writes

$$\dot{x}_1 = x_1(-x_4 + x_5 + x_6),
\dot{x}_2 = x_2(x_4 - x_5 + x_6),
\dot{x}_4 = 9x_1(x_1 - n_2x_2),
\dot{x}_5 = -9n_2x_2(x_1 - n_2x_2),
\dot{x}_6 = 0.$$
(7)

Moreover h_1 is a first integral of system (7). Now we restrict system (7) to $x_1 = 0$ and we obtain

$$\dot{x}_2 = x_2(x_4 - x_5 + x_6),
\dot{x}_4 = 0,
\dot{x}_5 = 9x_2^2,
\dot{x}_6 = 0.$$
(8)

Let $h_2 = h_1|_{x_1=0}$; h_2 is a polynomial first integral of system (8). Let $\tilde{H} = H|_{x_1=0} = -9x_2^2 - (x_4^2 + x_5^2 + x_6^2) + 2(x_4x_5 + x_4x_6 + x_5x_6)$. It is clear that \tilde{H} , x_4 and x_6 are first integrals of system (8), and that they are independent. Since system (8) is a four dimensional system with three independent first integrals \tilde{H} , x_4 and x_6 , we have $h_2 = h_2(x_4, x_6, \tilde{H})$. Applying Lemma 5 we get $h_1 = h_2(x_4, x_6, \tilde{H}) + x_1g_1$, where $g_1 = g_1(x_1, x_2, x_4, x_5, x_6)$ is a polynomial and h_2 eventually can be zero. Since $\tilde{H} = H + 9x_1(x_1 - 2n_2x_2)$ we have

$$h_1 = h_2(x_4, x_6, H) + x_1g_1 = h_2(x_4, x_6, H) + x_1g_2,$$

for some polynomial $g_2 = g_2(x_1, x_2, x_4, x_5, x_6)$.

We claim that h_2 does not depend on x_4 . We prove the claim: as x_6 and H are also invariant by σ we have

$$h_2(x_4, x_6, H) + x_1g_2 = h_1 = \sigma(h_1) = h_2(x_5, x_6, H) + n_2x_2\sigma(g_2).$$

Since the most left hand side and the most right hand side of the previous expression are equal, it follows that $h_2(x_4, x_6, H) = h_2(x_5, x_6, H)$. Therefore $h_2 = h_2(x_6, H)$. In short the claim is proved.

We have $h_1 = h_2(x_6, H) + x_1g_2$. Next we prove that $g_2 \equiv 0$ proceeding by contradiction. Suppose that $g_2 \neq 0$ and write $x_1g_2 = x_1^{\alpha}g_3$, with $\alpha \in \mathbb{N}$, $g_3 \neq 0$ and $x_1 \nmid g_3$. Let X be the vector field associated to system (7), i.e.

$$X = x_1(-x_4 + x_5 + x_6)\frac{\partial}{\partial x_1} + x_2(x_4 - x_5 + x_6)\frac{\partial}{\partial x_2} + 9x_1(x_1 - n_2x_2)\frac{\partial}{\partial x_4} - 9n_2x_2(x_1 - n_2x_2)\frac{\partial}{\partial x_5}.$$

As $h_1 = h_2(x_6, H) + x_1^{\alpha}g_3$, x_6 and H are first integrals of system (7) we have

$$0 = \frac{X(h_1)}{x_1^{\alpha}} = \frac{X(x_1^{\alpha}g_3)}{x_1^{\alpha}}$$

= $x_1(-x_4 + x_5 + x_6)\frac{\partial g_3}{\partial x_1} + x_2(x_4 - x_5 + x_6)\frac{\partial g_3}{\partial x_2}$ (9)
+ $9x_1(x_1 - n_2x_2)\frac{\partial g_3}{\partial x_4} - 9n_2x_2(x_1 - n_2x_2)\frac{\partial g_3}{\partial x_5} + \alpha(-x_4 + x_5 + x_6)g_3.$

Let $g_4 = g_3|_{x_1=0}$. Since g_4 satisfies equation (9) applied on $x_1 = 0$, we obtain

$$x_2(x_4 - x_5 + x_6)\frac{\partial g_4}{\partial x_2} + 9x_2^2\frac{\partial g_4}{\partial x_5} + \alpha(-x_4 + x_5 + x_6)g_4 = 0.$$

Thus $x_2|g_4$ and $g_4 = x_2^\beta g_5$, with $\beta \in \mathbb{N}$, $g_5 \not\equiv 0$ and $x_2 \nmid g_5$. Hence

$$x_2(x_4 - x_5 + x_6)\frac{\partial g_5}{\partial x_2} + 9x_2^2\frac{\partial g_5}{\partial x_5} + \left[\alpha(-x_4 + x_5 + x_6) + \beta(x_4 - x_5 + x_6)\right]g_5 = 0,$$

and therefore $x_2|g_5$, a contradiction. Hence $g_2 \equiv 0$ and $h_1 = h_2(x_6, H)$.

Back to h, we proved that $h = h_2(x_6, H) + x_3g$. It remains to prove that $g \equiv 0$. Assume that $g \not\equiv 0$. Let $x_3g = x_3^{\gamma}\bar{g}_1$, with $\gamma \in \mathbb{N}$, $\bar{g}_1 \not\equiv 0$ and $x_3 \nmid \bar{g}_1$. Let Y be the vector field associated to system (5), i.e.

$$Y = x_1(-x_4 + x_5 + x_6)\frac{\partial}{\partial x_1} + x_2(x_4 - x_5 + x_6)\frac{\partial}{\partial x_2} + x_3(x_4 + x_5 - x_6)\frac{\partial}{\partial x_3} + 9x_1(x_1 - n_2x_2)\frac{\partial}{\partial x_4} - 9n_2x_2(x_1 - n_2x_2)\frac{\partial}{\partial x_5}.$$

As h, x_6 and H are first integrals of system (5), we have

$$0 = \frac{Y(h)}{x_3^{\gamma}} = \frac{Y(x_3^{\gamma}\bar{g}_1)}{x_3^{\gamma}}$$

= $x_1(-x_4 + x_5 + x_6)\frac{\partial\bar{g}_1}{\partial x_1} + x_2(x_4 - x_5 + x_6)\frac{\partial\bar{g}_1}{\partial x_2} + x_3(x_4 + x_5 - x_6)\frac{\partial\bar{g}_1}{\partial x_3}$ (10)
+ $9x_1(x_1 - n_2x_2)\frac{\partial\bar{g}_1}{\partial x_4} - 9n_2x_2(x_1 - n_2x_2)\frac{\partial\bar{g}_1}{\partial x_5} + \gamma(x_4 + x_5 - x_6)\bar{g}_1.$

Let $\bar{g}_2 = \bar{g}_1|_{x_3=0}$. We have that $\bar{g}_2 \not\equiv 0$ and that it satisfies equation (10) with $x_3 = 0$:

$$x_{1}(-x_{4}+x_{5}+x_{6})\frac{\partial\bar{g}_{2}}{\partial x_{1}} + x_{2}(x_{4}-x_{5}+x_{6})\frac{\partial\bar{g}_{2}}{\partial x_{2}} + 9x_{1}(x_{1}-n_{2}x_{2})\frac{\partial\bar{g}_{2}}{\partial x_{4}} - 9n_{2}x_{2}(x_{1}-n_{2}x_{2})\frac{\partial\bar{g}_{2}}{\partial x_{5}} + \gamma(x_{4}+x_{5}-x_{6})\bar{g}_{2} = 0.$$
(11)

We write $\bar{g}_2 = x_1^{\delta} \bar{g}_3$, with $\delta \in \mathbb{N} \cup \{0\}$, $\bar{g}_3 \neq 0$ and $x_1 \nmid \bar{g}_3$. Then, after dividing by x_1^{δ} we get

$$x_{1}(-x_{4}+x_{5}+x_{6})\frac{\partial\bar{g}_{3}}{\partial x_{1}} + x_{2}(x_{4}-x_{5}+x_{6})\frac{\partial\bar{g}_{3}}{\partial x_{2}} + 9x_{1}(x_{1}-n_{2}x_{2})\frac{\partial\bar{g}_{3}}{\partial x_{4}} - 9n_{2}x_{2}(x_{1}-n_{2}x_{2})\frac{\partial\bar{g}_{3}}{\partial x_{5}} + \left[\gamma(x_{4}+x_{5}-x_{6})+\delta(-x_{4}+x_{5}+x_{6})\right]\bar{g}_{3} = 0.$$
(12)

Let $\bar{g}_4 = \bar{g}_3|_{x_1=0}$. It satisfies equation (12) with $x_1 = 0$:

$$x_2(x_4 - x_5 + x_6)\frac{\partial \bar{g}_4}{\partial x_2} + 9x_2^2\frac{\partial \bar{g}_4}{\partial x_5} + \left[\gamma(x_4 + x_5 - x_6) + \delta(-x_4 + x_5 + x_6)\right]\bar{g}_4 = 0.$$
(13)

From this equation we have $x_2|\bar{g}_4$, and hence $\bar{g}_4 = x_2^{\lambda}\bar{g}_5$, with $\lambda \in \mathbb{N}$, $\bar{g}_5 \neq 0$ and $x_2 \nmid \bar{g}_5$. After dividing by x_2^{λ} we obtain

$$x_{2}(x_{4} - x_{5} + x_{6})\frac{\partial \bar{g}_{5}}{\partial x_{2}} + 9x_{2}^{2}\frac{\partial \bar{g}_{5}}{\partial x_{5}} + \left[\gamma(x_{4} + x_{5} - x_{6}) + \delta(-x_{4} + x_{5} + x_{6}) + \lambda(x_{4} - x_{5} + x_{6})\right]\bar{g}_{5} = 0.$$
(14)

We note that the factor of \bar{g}_5 in (14) is

$$(-\delta + \lambda + \gamma)x_4 + (\delta - \lambda + \gamma)x_5 + (\delta + \lambda - \gamma)x_6.$$

The coefficients of x_4, x_5, x_6 in this expression are identically zero if and only if $\delta = \lambda = \gamma = 0$, which is not possible because γ and λ are different from zero. Hence we get from (14) that $x_2|\bar{g}_5$, a contradiction. Back to h, we have $g \equiv 0$ and then $h = h_2(x_6, H)$. The theorem follows.

Undoing the changes of variables, system (2) for Bianchi types VI₀ and VII₀ has only two independent polynomial first integrals which are $H_1 = p_3q_3$ and $H_2 = -(q_1 - n_2q_2)^2 - 4(p_1^2q_1^2 + p_2^2q_2^2 + p_3^2q_3^2) + 8(p_1p_2q_1q_2 + p_1p_3q_1q_3 + p_2p_3q_2q_3)$, with $n_2 = \pm 1$ depending on the Bianchi type. We note that we can obtain the Hamiltonian (1) from H_2 : $\mathcal{H} = H_2/4$.

Proposition 3 and Theorem 4 prove statement (c) of Theorem 1.

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