

HOPF BIFURCATION IN THE FULL REPRESSILATOR EQUATIONS

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ABSTRACT. In this paper we prove that the full repressilator equations, in dimension six undergo a supercritical Hopf bifurcation.

1. INTRODUCTION

Oscillatory networks are a particular kind of regulatory molecular networks, i.e., collections of interacting molecules in a cell. The regulatory oscillators can be used to study abnormalities of a process in the cell, from sleep disorders to cancer. So, they attract significant attention among biologists and biophysicists. There are many implementations of artificial oscillatory networks (see, e.g., [1, 2, 6, 7, 8, 10, 11]). One of them is the repressilator [5]. Its genetic implementation uses three proteins that cyclically repress the synthesis of one another. The following system of differential equations describes the behavior of the repressilator:

$$(1) \quad \begin{aligned} \dot{m}_1 &= -m_1 + \frac{\alpha}{1+v^n} + \alpha_0, \\ \dot{m}_2 &= -m_2 + \frac{\alpha}{1+w^n} + \alpha_0, \\ \dot{m}_3 &= -m_3 + \frac{\alpha}{1+u^n} + \alpha_0, \\ \dot{u} &= -\beta(u - m_1), \\ \dot{v} &= -\beta(v - m_2), \\ \dot{w} &= -\beta(w - m_3). \end{aligned}$$

Here u , v and w are proportional to the protein concentration, while m_i are proportional to the concentration of mRNA corresponding to those proteins. The nonlinear function $f(x) = \frac{\alpha}{1+x^n}$ reflects synthesis of the mRNAs from the DNA controlled by regulatory elements. The parameter α_0 represents uncontrolled part of the mRNA synthesis, and it is usually small. The

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explicit inclusion of the mRNA concentration variables into the model is given by β . Given that in general $\beta \ll 1$ and α_0 is very small, we consider $\alpha_0 = \varepsilon a$ and $\beta = \varepsilon b$, where a and b are positive constants and $\varepsilon > 0$ is sufficiently small. So, system (1) becomes

$$(2) \quad \begin{aligned} \dot{m}_1 &= -m_1 + \frac{\alpha}{1 + v^n} + \varepsilon a, \\ \dot{m}_2 &= -m_2 + \frac{\alpha}{1 + w^n} + \varepsilon a, \\ \dot{m}_3 &= -m_3 + \frac{\alpha}{1 + u^n} + \varepsilon a, \\ \dot{u} &= -\varepsilon b(u - m_1), \\ \dot{v} &= -\varepsilon b(v - m_2), \\ \dot{w} &= -\varepsilon b(w - m_3). \end{aligned}$$

In the papers [3, 4] the authors consider a reduced system of dimension three. The reduction assumes that the three excluded variables, i.e. m_i , evolve an order of magnitude faster than the other three. In [3] the authors prove that in the reduced system exhibits a supercritical Hopf bifurcation. The existence of a Hopf bifurcation in the reduced system does not imply that the full system, in dimension six, also has a supercritical Hopf bifurcation. It just gives an indication about its existence. Here in this work we consider the full system in dimension six and we extend the results of [3, 4] on the supercritical Hopf bifurcation to the 6-dimensional differential system (2). Our main result is the following one.

Theorem 1. *Let $n > 2$ be an integer. The point $(r_0, r_0, r_0, r_0, r_0, r_0)$ with*

$$r_0 = \sqrt[n]{\frac{2}{n-2}} + \varepsilon \frac{a - 3b \sqrt[n]{\frac{2}{n-2}}}{n-2} + \mathcal{O}(\varepsilon^2),$$

is an equilibrium of the differential system (2) with

$$\alpha = \alpha_{bif} = \frac{n}{n-2} \sqrt[n]{\frac{2}{n-2}} + \varepsilon \frac{n}{(n-2)^2} \left(-a(n-5) - 9b \sqrt[n]{\frac{2}{n-2}} \right) + \mathcal{O}(\varepsilon^2).$$

The eigenvalues of the linear part of system (2) at this equilibrium are $\{\pm \varepsilon \sqrt{3}bi + \mathcal{O}(\varepsilon^2), -1 + \varepsilon(-1 \pm \sqrt{3}i)b + \mathcal{O}(\varepsilon^2), -\varepsilon 3b + \mathcal{O}(\varepsilon^2), -1 + \varepsilon 2b + \mathcal{O}(\varepsilon^2)\}$. Moreover, there is a single supercritical Hopf bifurcation at $\alpha = \alpha_{bif}$ and there exists a small $\varepsilon_0 > 0$ such that for $\alpha_{bif} < \alpha < \alpha_{bif} + \varepsilon_0$ the system (2) has a stable limit cycle.

2. PROOF OF THEOREM 1

It is clear that system (2) has the equilibrium $p_0 = (r_0, r_0, r_0, r_0, r_0, r_0)$ where r_0 is solution of the equation

$$(3) \quad \frac{\alpha}{1+r^n} = r - \varepsilon a.$$

From (3) we have that $\alpha = (r_0 - \varepsilon a)(1 + r_0^n)$. So, substituting α in the linear part of system (2) at the equilibrium p_0 we get

$$A = \begin{pmatrix} -1 & 0 & 0 & 0 & \Delta & 0 \\ 0 & -1 & 0 & 0 & 0 & \Delta \\ 0 & 0 & -1 & \Delta & 0 & 0 \\ \varepsilon b & 0 & 0 & -\varepsilon b & 0 & 0 \\ 0 & \varepsilon b & 0 & 0 & -\varepsilon b & 0 \\ 0 & 0 & \varepsilon b & 0 & 0 & -\varepsilon b \end{pmatrix},$$

where $\Delta = -\frac{nr_0^{-1+n}(r_0 - \varepsilon a)}{1 + r_0^n}$. The eigenvalues of M are

$$\begin{aligned} & \varepsilon \frac{(-2 + (n-2 \pm i\sqrt{3}n)r_0^n)b}{2(1+r_0^n)} + \mathcal{O}(\varepsilon^2), & -1 + \varepsilon \frac{(-1 \pm i\sqrt{3})nr_0^n b}{2(1+r_0^n)} + \mathcal{O}(\varepsilon^2), \\ & -1 + \varepsilon \frac{nr_0^n b}{1+r_0^n} + \mathcal{O}(\varepsilon^2), & -\varepsilon \frac{(1 + (1+n)r_0^n) b}{1+r_0^n} + \mathcal{O}(\varepsilon^2). \end{aligned}$$

We impose that the real part of the eigenvalues $\varepsilon \frac{(-2 + (n-2 \pm i\sqrt{3}n)r_0^n)b}{2(1+r_0^n)} + \mathcal{O}(\varepsilon^2)$ is zero and we obtain

$$(4) \quad r_0 = \sqrt[n]{\frac{2}{n-2}} + \varepsilon \frac{a - 3b \sqrt[n]{\frac{2}{n-2}}}{n-2} + \mathcal{O}(\varepsilon^2).$$

Substituting (4) in (3) we get

$$(5) \quad \alpha_{bif} = \frac{n}{n-2} \sqrt[n]{\frac{2}{n-2}} - \varepsilon \frac{n}{(n-2)^2} \left(a(n-5) + 9b \sqrt[n]{\frac{2}{n-2}} \right) + \mathcal{O}(\varepsilon^2).$$

Substituting (4) in M and computing the eigenvalues we obtain $\pm \varepsilon \sqrt{3}bi + \mathcal{O}(\varepsilon^2)$, $-1 + \varepsilon(-1 \pm \sqrt{3}i)b + \mathcal{O}(\varepsilon^2)$, $-\varepsilon 3b + \mathcal{O}(\varepsilon^2)$, $-1 + \varepsilon 2b + \mathcal{O}(\varepsilon^2)$.

The linearization of (2) at p_0 has a pair of conjugate purely imaginary eigenvalues and the other four eigenvalues have negative real part. This is the setting for a Hopf bifurcation. We can expect to see a small-amplitude

limit cycle branching from the fixed point p_0 . It remains to compute the first Lyapunov coefficient $\ell_1(p_0)$ of (2) near p_0 . When $\ell_1(p_0) < 0$ the point p_0 is a weak focus of system (2) restricted to the central manifold of p_0 and the limit cycle that emerges from p_0 is stable. In this case we say that the Hopf bifurcation is supercritical. When $\ell_1(p_0) > 0$ the point p_0 is also a weak focus of system (2) restricted to the central manifold of p_0 but the limit cycle that borns from p_0 is unstable. In this second case we say that the Hopf bifurcation is subcritical.

Here we use the following result presented on page 180 of the book [9] for computing $\ell_1(p_0)$.

Lemma 2. *Let $\dot{x} = F(x)$ be a differential system having p_0 as an equilibrium point. Consider the third order Taylor approximation of F around p_0 given by $F(x) = Ax + \frac{1}{2!}B(x, x) + \frac{1}{3!}C(x, x, x) + \mathcal{O}(|x|^4)$. Assume that A has a pair of purely imaginary eigenvalues $\pm\lambda i$. Let q be the eigenvector of A corresponding to the eigenvalue λi , normalized so that $\bar{q} \cdot q = 1$, where \bar{q} is the conjugate vector of q . Let p be the adjoint eigenvector such that $A^T p = -\lambda i p$ and $\bar{p} \cdot q = 1$. If I denotes the 6×6 identity matrix, then*

$$\ell_1(p_0) = \frac{1}{2\lambda} \text{Re}(\bar{p} \cdot C(q, q, \bar{q}) - 2\bar{p} \cdot B(q, A^{-1}B(q, \bar{q})) + \bar{p} \cdot B(\bar{q}, (2\lambda i I - A)^{-1}B(q, q))).$$

In our case the linear part of system (2) at the equilibrium p_0 is

$$A = \begin{pmatrix} -1 & 0 & 0 & 0 & -2 + \varepsilon\sigma & 0 \\ 0 & -1 & 0 & 0 & 0 & -2 + \varepsilon\sigma \\ 0 & 0 & -1 & -2 + \varepsilon\sigma & 0 & 0 \\ b\varepsilon & 0 & 0 & -b\varepsilon & 0 & 0 \\ 0 & b\varepsilon & 0 & 0 & -b\varepsilon & 0 \\ 0 & 0 & b\varepsilon & 0 & 0 & -b\varepsilon \end{pmatrix} + \mathcal{O}(\varepsilon^2),$$

where

$$\sigma = 6b + \frac{1}{(n-2)^2} \left(b^2(51 - 39n) + ba12 \sqrt[n]{\frac{n-2}{2}} - a^2 \sqrt[n]{\left(\frac{n-2}{2}\right)^2 (n-1)} \right).$$

We have that A has an eigenvalue $\varepsilon\sqrt{3}bi + \mathcal{O}(\varepsilon^2)$. Now we compute the bi- and tri-linear functions B and C . Considering the vector field $(f_1, f_2, f_3, f_4, f_5, f_6)$ associated to the differential system (2) we observe that all second and third

derivatives vanishes except $\frac{\partial^2 f_1}{\partial v^2}$, $\frac{\partial^2 f_2}{\partial w^2}$, $\frac{\partial^2 f_3}{\partial u^2}$, $\frac{\partial^3 f_1}{\partial v^3}$, $\frac{\partial^3 f_2}{\partial w^3}$ and $\frac{\partial^3 f_3}{\partial u^3}$. Computing these derivatives, taking into account (4) and (5) we get that

$$\frac{\partial^2 f_1}{\partial v^2}(p_0) = \frac{\partial^2 f_2}{\partial w^2}(p_0) = \frac{\partial^2 f_3}{\partial u^2}(p_0) = \gamma,$$

$$\frac{\partial^3 f_1}{\partial v^3}(p_0) = \frac{\partial^3 f_2}{\partial w^3}(p_0) = \frac{\partial^3 f_3}{\partial u^3}(p_0) = \delta,$$

where

$$\begin{aligned} \gamma = & -2(n-5) \sqrt[n]{\frac{n-2}{2}} \\ & + \varepsilon \sqrt[n]{\frac{2^{n-2}}{(n-2)^{n-2}}} \left(a(5n-13) + 3b \sqrt[n]{\frac{2}{n-2}} (n^2 - 12n + 23) \right) + \mathcal{O}(\varepsilon^2), \end{aligned}$$

and

$$\begin{aligned} \delta = & -2 \sqrt[n]{\left(\frac{n-2}{2}\right)^2} (n^2 - 15n + 38) \\ & + \varepsilon \frac{2}{n-2} \sqrt[n]{\left(\frac{n-2}{2}\right)^3} (2a(7n^2 - 57n + 98) \\ & + 3b \sqrt[n]{\frac{2}{n-2}} (n^3 - 31n^2 + 182n - 272)) + \mathcal{O}(\varepsilon^2). \end{aligned}$$

So, the bilinear function B is given by

$$B((x_1, y_1, z_1, u_1, v_1, w_1), (x_2, y_2, z_2, u_2, v_2, w_2)) = (\gamma v_1 v_2, \gamma w_1 w_2, \gamma u_1 u_2, 0, 0, 0),$$

and the tri-linear function C is given by the expression

$$\begin{aligned} C((x_1, y_1, z_1, u_1, v_1, w_1), (x_2, y_2, z_2, u_2, v_2, w_2), (x_3, y_3, z_3, u_3, v_3, w_3)) = \\ (\delta v_1 v_2 v_3, \delta w_1 w_2 w_3, \delta u_1 u_2 u_3, 0, 0, 0). \end{aligned}$$

Computing the normalized eigenvector q of A , associated to the eigenvalue $\varepsilon\sqrt{3}bi + \mathcal{O}(\varepsilon^2)$, we obtain

$$\begin{aligned} q = & \left(\frac{1 - \sqrt{3}i}{\sqrt{15}}, -\frac{2}{\sqrt{15}}, \frac{1 + \sqrt{3}i}{\sqrt{15}}, \frac{-1 - \sqrt{3}i}{2\sqrt{15}}, \frac{-1 + \sqrt{3}i}{2\sqrt{15}}, \frac{1}{\sqrt{15}} \right) \\ & + \varepsilon \left(\frac{-30(\sqrt{3} + i)b - (\sqrt{3} - 3i)\sigma}{30\sqrt{5}}, \frac{60bi + 2\sqrt{3}\sigma}{30\sqrt{5}}, \frac{30(\sqrt{3} - i)b - (\sqrt{3} + 3i)\sigma}{30\sqrt{5}}, \right. \\ & \left. -\frac{(2\sqrt{3} + 6i)\sigma}{30\sqrt{5}}, -\frac{(2\sqrt{3} - 6i)\sigma}{30\sqrt{5}}, \frac{4\sqrt{3}\sigma}{30\sqrt{5}} \right) + \mathcal{O}(\varepsilon^2). \end{aligned}$$

The normalized adjoint eigenvector of the transpose matrix A with the eigenvalue $-\varepsilon\sqrt{3}bi$ is

$$\begin{aligned} p = & \left(0, 0, 0, \frac{-\sqrt{15} - 3\sqrt{5}i}{6}, \frac{-\sqrt{15} + 3\sqrt{5}i}{6}, \frac{\sqrt{15}}{3} \right) \\ & + \varepsilon \left(\frac{(-5 - 5\sqrt{3}i)b}{2\sqrt{15}}, \frac{(-5 + 5\sqrt{3}i)b}{2\sqrt{15}}, \frac{5b}{\sqrt{15}}, \frac{10b + (1 + \sqrt{3}i)\sigma}{\sqrt{15}}, \right. \\ & \left. \frac{-5(1 + \sqrt{3}i)b + (1 - \sqrt{3}i)\sigma}{\sqrt{15}}, \frac{-5(1 - \sqrt{3}i)b - 2\sigma}{\sqrt{15}} \right) + \mathcal{O}(\varepsilon^2). \end{aligned}$$

According to Lemma 2, in order to compute $\ell_1(p_0)$, we need to compute first A^{-1} and $(2\sqrt{3}b\varepsilon iI - A)^{-1}$. We have that $A^{-1} = \frac{1}{\varepsilon}A_{-1} + A_0 + \varepsilon A_1 + \mathcal{O}(\varepsilon^2)$,

$$\begin{aligned} A_{-1} = & \frac{1}{9b} \begin{pmatrix} 0 & 0 & 0 & 8 & 2 & -4 \\ 0 & 0 & 0 & -4 & 8 & 2 \\ 0 & 0 & 0 & 2 & -4 & 8 \\ 0 & 0 & 0 & -1 & 2 & -4 \\ 0 & 0 & 0 & -4 & -1 & 2 \\ 0 & 0 & 0 & 2 & -4 & -1 \end{pmatrix}, \\ A_0 = & \frac{1}{27b} \begin{pmatrix} -3b & 6b & -12b & -4\sigma & 5\sigma & -4\sigma \\ -12b & -3b & 6b & -4\sigma & -4\sigma & 5\sigma \\ 6b & -12b & -3b & 5\sigma & -4\sigma & -4\sigma \\ -3b & 6b & -12b & -4\sigma & 5\sigma & -4\sigma \\ -12b & -3b & 6b & -4\sigma & -4\sigma & 5\sigma \\ 6b & -12b & -3b & 5\sigma & -4\sigma & -4\sigma \end{pmatrix} \text{ and} \end{aligned}$$

$$A_1 = \frac{\sigma}{81b} \begin{pmatrix} -12b & 15b & -12b & -10\sigma & 8\sigma & -\sigma \\ -12b & -12b & 15b & -\sigma & -10\sigma & 8\sigma \\ 15b & -12b & -12b & 8\sigma & -\sigma & -10\sigma \\ -12b & 15b & -12b & -10\sigma & 8\sigma & -\sigma \\ -12b & -12b & 15b & -\sigma & -10\sigma & 8\sigma \\ 15b & -12b & -12b & 8\sigma & -\sigma & -10\sigma \end{pmatrix}.$$

and the expression of $(2\sqrt{3}b\epsilon iI - A)^{-1}$ is very large and we present it in the appendix.

The first, second and third terms of $\ell_1(p_0)$ given in Lemma 2 respectively are

$$Re(\bar{p} \cdot C(q, q, \bar{q})) = \frac{1}{15} \sqrt[n]{\left(\frac{n-2}{2}\right)^2} (n^2 - 15n + 38)b\epsilon + \mathcal{O}(\epsilon^2),$$

$$Re(-2\bar{p} \cdot B(q, A^{-1}B(q, \bar{q}))) = -\frac{4}{45} \sqrt[n]{\left(\frac{n-2}{2}\right)^2} (n^2 - 10n + 25)b\epsilon + \mathcal{O}(\epsilon^2)$$

and

$$Re(\bar{p} \cdot B(\bar{q}, (2\lambda iI - A)^{-1}B(q, q))) = 0 + \mathcal{O}(\epsilon^2).$$

Consequently we get

$$\ell_1(p_0) = -\frac{2}{45\sqrt{3}} \sqrt[n]{\left(\frac{n-2}{2}\right)^2} (n+7)(n-2) + \mathcal{O}(\epsilon).$$

As we said before $\ell_1(p_0) < 0$ implies that we have a supercritical Hopf bifurcation at $\alpha = \alpha_{bif}$, so there exists $\epsilon_0 > 0$ such that for $\alpha_{bif} < \alpha < \alpha_{bif} + \epsilon_0$ the system (2) has a stable limit cycle.

3. CONCLUSIONS

The repressilator model is an implementation of an artificial oscillatory network used for studying the collections of interacting molecules in a cell. The model is given by a 6-dimensional differential system. O. Buse, A. Kuznetsov and R. Pérez published two nice papers (see [3, 4]) analysing a reduced system of dimension three. In this reduced system they show the existence of a supercritical Hopf bifurcation. Since the reduction is reasonable we may expect that such supercritical Hopf bifurcation must also occurs in the actual 6-dimensional differential system. We prove that this is the case.

4. APPENDIX

The matrix $(2\sqrt{3}b\epsilon iI - A)^{-1}$ is given by

$$(2\sqrt{3}b\epsilon iI - A)^{-1} = \begin{pmatrix} C_{11} & C_{12} & C_{13} & C_{14} & C_{15} & C_{16} \\ C_{21} & C_{22} & C_{23} & C_{24} & C_{25} & C_{26} \\ C_{31} & C_{32} & C_{33} & C_{34} & C_{35} & C_{36} \\ C_{41} & C_{42} & C_{43} & C_{44} & C_{45} & C_{46} \\ C_{51} & C_{52} & C_{53} & C_{54} & C_{55} & C_{56} \\ C_{61} & C_{62} & C_{63} & C_{64} & C_{65} & C_{66} \end{pmatrix},$$

where $C_{ij} = A_{ij} + B_{ij}$, for $i, j = 1, 2, \dots, 6$.

$$\begin{aligned} A_{11} &= \frac{1}{189} (213 - 16i\sqrt{3}), \\ A_{12} &= \frac{2}{189} (-9 + 34i\sqrt{3}), \\ A_{13} &= \frac{4}{189} (-15 - 4i\sqrt{3}), \\ A_{14} &= \frac{4(12(\sqrt{3} - 124i)b\epsilon - (124\sqrt{3} + 3i)\sigma\epsilon + 72\sqrt{3} + 18i)}{81(10\sqrt{3} + 27i)b\epsilon}, \\ A_{15} &= \frac{3156\sqrt{3}b\epsilon + 6144ib\epsilon + 512\sqrt{3}\sigma\epsilon - 789i\sigma\epsilon - 864\sqrt{3} + 666i}{810\sqrt{3}b\epsilon + 2187ib\epsilon}, \\ A_{16} &= \frac{4(12(50 + 87i\sqrt{3})b\epsilon + (261 - 50i\sqrt{3})\sigma\epsilon + 18i\sqrt{3} - 225)}{81(27 - 10i\sqrt{3})b\epsilon}, \\ A_{21} &= \frac{4(12(50 + 87i\sqrt{3})b\epsilon + (261 - 50i\sqrt{3})\sigma\epsilon + 18i\sqrt{3} - 225)}{81(27 - 10i\sqrt{3})b\epsilon}, \\ A_{22} &= \frac{1}{189} (213 - 16i\sqrt{3}), \\ A_{23} &= \frac{1}{189} (213 - 16i\sqrt{3}), \\ A_{24} &= \frac{4(12(50 + 87i\sqrt{3})b\epsilon + (261 - 50i\sqrt{3})\sigma\epsilon + 18i\sqrt{3} - 225)}{81(27 - 10i\sqrt{3})b\epsilon}, \\ A_{25} &= \frac{4(12(50 + 87i\sqrt{3})b\epsilon + (261 - 50i\sqrt{3})\sigma\epsilon + 18i\sqrt{3} - 225)}{81(27 - 10i\sqrt{3})b\epsilon}, \\ A_{26} &= \frac{4(12(50 + 87i\sqrt{3})b\epsilon + (261 - 50i\sqrt{3})\sigma\epsilon + 18i\sqrt{3} - 225)}{81(27 - 10i\sqrt{3})b\epsilon}, \\ A_{31} &= \frac{4(12(50 + 87i\sqrt{3})b\epsilon + (261 - 50i\sqrt{3})\sigma\epsilon + 18i\sqrt{3} - 225)}{81(27 - 10i\sqrt{3})b\epsilon}, \end{aligned}$$

$$\begin{aligned}
 A_{32} &= \frac{4(12(50 + 87i\sqrt{3})b\epsilon + (261 - 50i\sqrt{3})\sigma\epsilon + 18i\sqrt{3} - 225)}{81(27 - 10i\sqrt{3})b\epsilon}, \\
 A_{33} &= \frac{4(12(50 + 87i\sqrt{3})b\epsilon + (261 - 50i\sqrt{3})\sigma\epsilon + 18i\sqrt{3} - 225)}{81(27 - 10i\sqrt{3})b\epsilon}, \\
 A_{34} &= \frac{4(12(50 + 87i\sqrt{3})b\epsilon + (261 - 50i\sqrt{3})\sigma\epsilon + 18i\sqrt{3} - 225)}{81(27 - 10i\sqrt{3})b\epsilon}, \\
 A_{35} &= \frac{4(12(50 + 87i\sqrt{3})b\epsilon + (261 - 50i\sqrt{3})\sigma\epsilon + 18i\sqrt{3} - 225)}{81(27 - 10i\sqrt{3})b\epsilon}, \\
 A_{36} &= \frac{4(12(50 + 87i\sqrt{3})b\epsilon + (261 - 50i\sqrt{3})\sigma\epsilon + 18i\sqrt{3} - 225)}{81(27 - 10i\sqrt{3})b\epsilon}, \\
 A_{41} &= \frac{1}{189}(9 - 34i\sqrt{3}), \\
 A_{42} &= \frac{2}{189}(15 + 4i\sqrt{3}), \\
 A_{43} &= \frac{4}{189}(-3 + 2i\sqrt{3}), \\
 A_{44} &= -\frac{912\sqrt{3}b\epsilon + 480ib\epsilon + 40\sqrt{3}\sigma\epsilon - 228i\sigma\epsilon - 432\sqrt{3} + 333i}{810\sqrt{3}b\epsilon + 2187ib\epsilon}, \\
 A_{45} &= \frac{1188\sqrt{3}b\epsilon - 984ib\epsilon - 82\sqrt{3}\sigma\epsilon - 297i\sigma\epsilon + 36\sqrt{3} + 450i}{810\sqrt{3}b\epsilon + 2187ib\epsilon}, \\
 A_{46} &= \frac{4(12(\sqrt{3} + 44i)b\epsilon + (44\sqrt{3} - 3i)\sigma\epsilon - 9(4\sqrt{3} + i))}{81(10\sqrt{3} + 27i)b\epsilon}, \\
 A_{51} &= \frac{4}{189}(-3 + 2i\sqrt{3}), \\
 A_{52} &= \frac{4}{189}(-3 + 2i\sqrt{3}), \\
 A_{53} &= \frac{2}{189}(15 + 4i\sqrt{3}), \\
 A_{54} &= \frac{4(12(\sqrt{3} + 44i)b\epsilon + (44\sqrt{3} - 3i)\sigma\epsilon - 9(4\sqrt{3} + i))}{81(10\sqrt{3} + 27i)b\epsilon}, \\
 A_{55} &= -\frac{912\sqrt{3}b\epsilon + 480ib\epsilon + 40\sqrt{3}\sigma\epsilon - 228i\sigma\epsilon - 432\sqrt{3} + 333i}{810\sqrt{3}b\epsilon + 2187ib\epsilon}, \\
 A_{56} &= \frac{1188\sqrt{3}b\epsilon - 984ib\epsilon - 82\sqrt{3}\sigma\epsilon - 297i\sigma\epsilon + 36\sqrt{3} + 450i}{810\sqrt{3}b\epsilon + 2187ib\epsilon}, \\
 A_{61} &= \frac{2}{189}(15 + 4i\sqrt{3}), \\
 A_{62} &= \frac{4}{189}(-3 + 2i\sqrt{3}), \\
 A_{63} &= \frac{1}{189}(9 - 34i\sqrt{3}), \\
 A_{64} &= \frac{1188\sqrt{3}b\epsilon - 984ib\epsilon - 82\sqrt{3}\sigma\epsilon - 297i\sigma\epsilon + 36\sqrt{3} + 450i}{810\sqrt{3}b\epsilon + 2187ib\epsilon},
 \end{aligned}$$

$$\begin{aligned}
A_{65} &= \frac{1188\sqrt{3}b\epsilon - 984ib\epsilon - 82\sqrt{3}\sigma\epsilon - 297i\sigma\epsilon + 36\sqrt{3} + 450i}{810\sqrt{3}b\epsilon + 2187ib\epsilon}, \\
A_{66} &= \frac{1188\sqrt{3}b\epsilon - 984ib\epsilon - 82\sqrt{3}\sigma\epsilon - 297i\sigma\epsilon + 36\sqrt{3} + 450i}{810\sqrt{3}b\epsilon + 2187ib\epsilon}, \\
B_{11} &= \frac{2\epsilon \left((-4800 - 5921i\sqrt{3})b + (-362 + 316i\sqrt{3})\sigma \right)}{3969}, \\
B_{12} &= \frac{\epsilon \left(8(215\sqrt{3} + 42i)b - (2\sqrt{3} + 259i)\sigma \right)}{81(4\sqrt{3} + i)}, \\
B_{13} &= \frac{\epsilon \left(8(215\sqrt{3} + 42i)b - (2\sqrt{3} + 259i)\sigma \right)}{81(4\sqrt{3} + i)}, \\
B_{14} &= \frac{2\epsilon}{81(2\sqrt{3} - 3i)^3 b} \left(96(444\sqrt{3} - 131i)b^2 - 52(41\sqrt{3} + 324i)b\sigma \right. \\
&\quad \left. + (-516\sqrt{3} + 271i)\sigma^2 \right), \\
B_{15} &= \frac{\epsilon}{81(2\sqrt{3} - 3i)^3 b} \left(24(-1760\sqrt{3} + 5807i)b^2 + 2(9247\sqrt{3} + 5952i)b\sigma \right. \\
&\quad \left. + 8(14\sqrt{3} - 215i)\sigma^2 \right), \\
B_{16} &= \frac{\epsilon}{81(2\sqrt{3} - 3i)^3 b} \left(-48(1058\sqrt{3} + 2825i)b^2 + 16(-1021\sqrt{3} \right. \\
&\quad \left. + 1362i)b\sigma + (758\sqrt{3} + 1259i)\sigma^2 \right), \\
B_{21} &= \frac{4\epsilon \left((-146\sqrt{3} + 528i)b + (32\sqrt{3} + 23i)\sigma \right)}{81(4\sqrt{3} + i)}, \\
B_{22} &= \frac{2\epsilon \left((-4800 - 5921i\sqrt{3})b + (-362 + 316i\sqrt{3})\sigma \right)}{3969}, \\
B_{23} &= \frac{\epsilon \left(8(215\sqrt{3} + 42i)b - (2\sqrt{3} + 259i)\sigma \right)}{81(4\sqrt{3} + i)}, \\
B_{24} &= \frac{\epsilon}{81(2\sqrt{3} - 3i)^3 b} \left(-48(1058\sqrt{3} + 2825i)b^2 + 16(-1021\sqrt{3} \right. \\
&\quad \left. + 1362i)b\sigma + (758\sqrt{3} + 1259i)\sigma^2 \right), \\
B_{25} &= \frac{2\epsilon}{81(2\sqrt{3} - 3i)^3 b} \left(96(444\sqrt{3} - 131i)b^2 \right. \\
&\quad \left. - 52(41\sqrt{3} + 324i)b\sigma + (-516\sqrt{3} + 271i)\sigma^2 \right), \\
B_{26} &= \frac{\epsilon}{81(2\sqrt{3} - 3i)^3 b} \left(24(-1760\sqrt{3} + 5807i)b^2 + 2(9247\sqrt{3} \right. \\
&\quad \left. + 5952i)b\sigma + 8(14\sqrt{3} - 215i)\sigma^2 \right), \\
B_{31} &= \frac{\epsilon \left(8(215\sqrt{3} + 42i)b - (2\sqrt{3} + 259i)\sigma \right)}{81(4\sqrt{3} + i)}, \\
B_{32} &= \frac{2\epsilon \left((-4800 - 5921i\sqrt{3})b + (-362 + 316i\sqrt{3})\sigma \right)}{3969}, \\
B_{33} &= \frac{2\epsilon \left((-4800 - 5921i\sqrt{3})b + (-362 + 316i\sqrt{3})\sigma \right)}{3969},
\end{aligned}$$

$$\begin{aligned}
 B_{34} &= \frac{\epsilon}{81(2\sqrt{3}-3i)^3 b} \left(24(-1760\sqrt{3}+5807i)b^2 + 2(9247\sqrt{3} \right. \\
 &\quad \left. + 5952i)b\sigma + 8(14\sqrt{3}-215i)\sigma^2 \right), \\
 B_{35} &= \frac{\epsilon}{81(2\sqrt{3}-3i)^3 b} \left(-48(1058\sqrt{3}+2825i)b^2 + 16(-1021\sqrt{3} \right. \\
 &\quad \left. + 1362i)b\sigma + (758\sqrt{3}+1259i)\sigma^2 \right), \\
 B_{36} &= \frac{2\epsilon}{81(2\sqrt{3}-3i)^3 b} \left(96(444\sqrt{3}-131i)b^2 - 52(41\sqrt{3}+324i)b\sigma \right. \\
 &\quad \left. + (-516\sqrt{3}+271i)\sigma^2 \right), \\
 B_{41} &= \frac{2\epsilon(-259\sqrt{3}b+6ib+8\sqrt{3}\sigma+22i\sigma)}{81(4\sqrt{3}+i)}, \\
 B_{42} &= -\frac{\epsilon(8(-23\sqrt{3}+96i)b+(40\sqrt{3}+19i)\sigma)}{81(4\sqrt{3}+i)}, \\
 B_{43} &= \frac{4\epsilon(2(35\sqrt{3}+54i)b+(6\sqrt{3}-13i)\sigma)}{81(4\sqrt{3}+i)}, \\
 B_{44} &= \frac{2\epsilon}{81(2\sqrt{3}-3i)^3 b} \left(96(14\sqrt{3}-215i)b^2 - 52(53\sqrt{3}+6i)b\sigma \right. \\
 &\quad \left. + (2\sqrt{3}+259i)\sigma^2 \right), \\
 B_{45} &= \frac{\epsilon}{81(2\sqrt{3}-3i)^3 b} \left(24(758\sqrt{3}+1259i)b^2 + 2(1627\sqrt{3} \right. \\
 &\quad \left. - 3810i)b\sigma - 8(32\sqrt{3}+23i)\sigma^2 \right), \\
 B_{46} &= \frac{\epsilon}{81(2\sqrt{3}-3i)^3 b} \left(48(-516\sqrt{3}+271i)b^2 + 16(131\sqrt{3} \right. \\
 &\quad \left. + 576i)b\sigma + (252\sqrt{3}-253i)\sigma^2 \right), \\
 B_{51} &= \frac{4\epsilon(2(35\sqrt{3}+54i)b+(6\sqrt{3}-13i)\sigma)}{81(4\sqrt{3}+i)}, \\
 B_{52} &= \frac{2\epsilon(-259\sqrt{3}b+6ib+8\sqrt{3}\sigma+22i\sigma)}{81(4\sqrt{3}+i)}, \\
 B_{53} &= -\frac{\epsilon(8(-23\sqrt{3}+96i)b+(40\sqrt{3}+19i)\sigma)}{81(4\sqrt{3}+i)}, \\
 B_{54} &= \frac{\epsilon}{81(2\sqrt{3}-3i)^3 b} \left(48(-516\sqrt{3}+271i)b^2 + 16(131\sqrt{3} \right. \\
 &\quad \left. + 576i)b\sigma + (252\sqrt{3}-253i)\sigma^2 \right), \\
 B_{55} &= \frac{2\epsilon}{81(2\sqrt{3}-3i)^3 b} \left(96(14\sqrt{3}-215i)b^2 - 52(53\sqrt{3}+6i)b\sigma \right. \\
 &\quad \left. + (2\sqrt{3}+259i)\sigma^2 \right), \\
 B_{56} &= \frac{\epsilon}{81(2\sqrt{3}-3i)^3 b} \left(24(758\sqrt{3}+1259i)b^2 + 2(1627\sqrt{3} \right. \\
 &\quad \left. - 3810i)b\sigma - 8(32\sqrt{3}+23i)\sigma^2 \right),
 \end{aligned}$$

$$\begin{aligned}
B_{61} &= -\frac{\epsilon (8(-23\sqrt{3} + 96i)b + (40\sqrt{3} + 19i)\sigma)}{81(4\sqrt{3} + i)}, \\
B_{62} &= \frac{4\epsilon (2(35\sqrt{3} + 54i)b + (6\sqrt{3} - 13i)\sigma)}{81(4\sqrt{3} + i)}, \\
B_{63} &= \frac{4\epsilon (2(35\sqrt{3} + 54i)b + (6\sqrt{3} - 13i)\sigma)}{81(4\sqrt{3} + i)}, \\
B_{64} &= \frac{\epsilon}{81(2\sqrt{3} - 3i)^3 b} \left(24(758\sqrt{3} + 1259i)b^2 + 2(1627\sqrt{3} \right. \\
&\quad \left. - 3810i)b\sigma - 8(32\sqrt{3} + 23i)\sigma^2 \right), \\
B_{65} &= \frac{\epsilon}{81(2\sqrt{3} - 3i)^3 b} \left(48(-516\sqrt{3} + 271i)b^2 + 16(131\sqrt{3} \right. \\
&\quad \left. + 576i)b\sigma + (252\sqrt{3} - 253i)\sigma^2 \right) \text{ and} \\
B_{66} &= \frac{2\epsilon}{81(2\sqrt{3} - 3i)^3 b} \left(96(14\sqrt{3} - 215i)b^2 - 52(53\sqrt{3} + 6i)b\sigma \right. \\
&\quad \left. + (2\sqrt{3} + 259i)\sigma^2 \right).
\end{aligned}$$

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