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Insight into the degradation mechanisms of Atomic Layer Deposited TiO₂ as photoanode protective layer.

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KEYWORDS

TiO₂, water splitting, photoanode, protective layer, ALD, degradation.

ABSTRACT

Around 100 nm thick TiO_2 layers deposited by atomic layer deposition (ALD) have been investigated as anticorrosion protective films for silicon based photoanodes decorated with 5 nm NiFe catalyst in highly alkaline electrolyte. Completely amorphous layers presented high resistivity, meanwhile the ones synthesized at 300°C , having a fully anatase crystalline TiO_2 structure, introduced insignificant resistance, showing direct correlation between crystallization degree and electrical conductivity. The conductivity through crystalline TiO_2 layers has been found not to be homogeneous, presenting preferential conduction paths attributed to grain boundaries and defects within the crystalline structure. A correlation between the c-AFM measurements and grain interstitials can be seen, supported by HRTEM cross sectional images presenting defective regions in crystalline TiO_2 grains. It was found that the conduction mechanism goes through the injection of electrons coming from water oxidation from the electrocatalyst into the TiO_2 conduction band. Then, electrons are transported to the $\text{Si}/\text{SiO}_x/\text{TiO}_2$ interface where electrons recombine with holes given by the p^+n -Si junction. No evidences of intra band gap states in TiO_2 responsible of conductivity have been detected. Stability measurements of fully crystalline samples over 480 h in anodic polarization show a continuous current decay. Electrochemical impedance spectroscopy allows to identify that the main cause of deactivation is associated to the loss of TiO_2 electrical conductivity, corresponding to a self-passivation mechanism. This is proposed to reflect the effect of OH^- ions diffusing in the TiO_2 structure in anodic conditions by the electric field. This fact proves that a modification takes place in the defective zone of the layer, blocking the ability to transfer electrical charge through the layer. According to this mechanism, a regeneration of the degradation process is demonstrated possible based on ultraviolet illumination, which contributes to change the occupancy of TiO_2 electronic states and to recover the defective

zones conductivity. These findings confirm the connection between the structural properties of the ALD-deposited polycrystalline layer and the degradation mechanisms, and thus highlight main concerns towards fabricating long-lasting metal oxide protective layers for frontal illuminated photoelectrodes.

1. INTRODUCTION

Our society faces a significant challenge regarding energy production and climate change, where renewable energy sources are called to be crucial. However, the high variability on renewables such as wind and solar power is one of the main drawbacks nowadays. Storing solar energy into chemical bonds is an interesting candidate to face this variability, the so called “solar fuels”.¹⁻⁴ Among them, hydrogen produced from water electrolysis is a promising candidate as energy vector, substituting natural gas as heat source⁵ or fossil fuels used for transportation.⁶ Photoelectrochemical (PEC) water splitting was discovered by Fujishima and Honda in the late 70’s,⁷ but solar to hydrogen (STH) efficiencies have been far from industrial requirements for decades, although remarkable efforts have been put in understanding and enhancing photoactive metal oxide semiconductors and catalysts to perform the water splitting reaction.⁸⁻¹⁸

With the significant increase in productivity and cost reduction of the photovoltaic industry, which has led to its large scale implementation, many authors have put their attention into implementing these semiconductor materials and fabrication methods in PEC water splitting.¹⁹⁻²¹ Many materials such as monocrystalline²²⁻²⁴ and amorphous silicon,²⁵ chalcogenide-based compounds,^{26,27} III-V semiconductors²⁸⁻³² and even perovskites³³ have been proposed. However, all of these semiconductors face a common problem: the chemical stability of such materials when exposed to the highly acidic or alkaline electrolytes necessary for kinetically-efficient oxygen and hydrogen evolution reactions (OER and HER).²⁰ To overcome this, the use of protective layers was suggested, avoiding the contact between the electrolyte and the sensible photoabsorber materials, but their fabrication routes should assure that the photoabsorber suffers no harm or alteration. In addition, protective layers for frontal illuminated photoelectrodes need to be conductive, transparent, chemically stable. These requirements reduce the candidate

materials and methods. Among them, TiO_2 is found to be one of the best performing candidates,^{21,34–38} and among deposition techniques, atomic layer deposition (ALD) allows high layer thickness control, conformal layers almost pinhole free, and low crystallization temperatures thanks to intermediate-species mobility, avoiding alterations of inner layers of the photoelectrode.^{39–42} For instance, recent studies have reported interesting characteristics of ALD- TiO_2 such as increased adhesion on Si photoanodes⁴³ or enhanced plasmonic resonance in photovoltaic devices⁴⁴.

Photocathodes protected with polycrystalline TiO_2 have been widely reported as the conduction band alignment between the silicon and the TiO_2 metal oxide facilitates the electron transfer from the photocathode to the electrolyte through the TiO_2 layer that is n-type.^{22,26,45–49}

Regarding TiO_2 protecting silicon based photoanodes, there is a totally different casuistic. In this case, the photo-generated minority carriers, holes, must be transferred to the electrolyte through a protective n-type TiO_2 layer. Different works have studied TiO_2 as protective layer for photoanodes, mainly using silicon as the photoabsorbing material, either forming an heterojunction between the TiO_2 and the n-Si^{50–55} or with a buried p^+-n junction in silicon.^{28,56,57} Under these boundaries, as a n-type metal oxide layer was expected to be highly resistive when thicker than the tunnelling distance, first works with TiO_2 grown via ALD as protective layer were using thicknesses ranging of only few nanometers.^{50,51,55,57,58} However, these assumptions were put in evidence by the studies carried out by Hu *et al.*²⁸ These authors used layers with thicknesses over 100 nm which were found to have conductivity enough to validate the use of TiO_2 for photoanodes as feasible. They propose state-mediated hole transport through intra-band gap levels as the conduction path caused by the disorder of amorphous TiO_2 (“electronically

leaky layer”), in alignment with the silicon valence band, following the same mechanism postulated decades ago by Campet et al.⁵⁹

Other authors that claim not identifying this mid bandgap levels in their samples have proposed a mechanism based on a MIM-like (Metal-Insulator-Metal) structure defined by a TiO₂ layer between the photoabsorber and the electrocatalyst. Thus, Mei et al.⁵⁶ proposed that electrons were conducted through the conduction band for a p⁺-Si/Ti/TiO₂/Pt structure. So, under this model, electrons from the oxygen evolution reactions through the electrocatalyst are easily injected to the conduction band of the titanium oxide layer and transferred to the internal Si/TiO₂ interface at the photoanode where they recombine with the photogenerated holes. However, these works generally do not consider the polycrystalline nature of the TiO₂ layer and neither their degradation mechanisms, which they are often overlooked.

In this context, this work aims at revealing the role of the amorphous / polycrystalline nature of the atomically deposited TiO₂ layer, considering its defective zones, to determine the electrical conductivity pathways and their evolution over time as a consequence of the degradation mechanisms that take place, which determines the endurance to the use of the protective layers based on TiO₂ in silicon photoanodes for water splitting.

2. EXPERIMENTAL SECTION

TiO₂ coatings have been grown by ALD on p⁺-n silicon buried junctions and simultaneously on p⁺ degenerately doped silicon, to simulate direct injection of carriers in dark conditions.

p⁺-Si samples were created by cutting a degenerately doped silicon wafer (0.001 $\Omega\cdot\text{cm}$) in 1x1 cm² pieces, and 50 nm of Al were thermally evaporated as back contact.

For the p⁺n-Si samples, a 1 cm² active area was lithographically defined by SiO₂ passivation on a silicon n-type wafer (0.1-0.5 $\Omega\cdot\text{cm}$ resistivity). Boron was implanted in the defined front surface and activated by rapid thermal annealing, creating a 200 nm n⁺ region on top of the n-type substrate. As back contact, 1 μm of Al/0.5%Cu was sputtered on top of 30 nm of Ti to form a proper ohmic contact.

p⁺-Si and p⁺n-Si samples were sonicated for 5 min in a 1:1:1 isopropanol, acetone and DI water cleaning solution, followed by abundant rinsing and further 5 min sonication in DI water. Before ALD deposition, front surface was cleaned with 0.1M HF for 5 min, rinsed in DI water and immediately introduced in the ALD chamber.

A R200 Picosun Atomic Layer Deposition system was used to grow TiO₂ layers. TiCl₄ was selected as precursor because of its wide deposition window and lower cost compared with other metalorganic compounds. Thus, TiCl₄ and H₂O precursors at 19 °C were used in successive pulses at 8 mbar in N₂ flow atmosphere, with 0.1 s pulses and 10 s purges. Under these conditions, layers were grown at deposition temperatures of 100, 150 and 300 °C for 3300 cycles, corresponding to roughly 100 nm layers. Layer thickness was measured by evaluating the reflected spectra with a Sensofar interferometer device with ± 0.2 nm error. Finally, nickel-iron (NiFe) was deposited by thermal evaporation using an Oerlikon Univex 250 equipment, with a thickness of 5 nm measured by a quartz microbalance. Samples were then soldered to a Cu wire using Ag paint and epoxy protected leaving the front area exposed (Figure S1).

Surface morphology was observed with a Zeiss Series Auriga Field Emission Scanning Electron Microscope (FESEM). Structural characterization was carried out by X-ray diffraction (XRD) in a D8 Advance Bruker equipment with a Cu K α radiation source working at 40 kV and 40 mA with a 3 ° offset angle. Raman spectroscopy was performed using a Horiba LabRam HR 800 with 532 nm laser light with a 100x objective and an output power of 5 mW, 10 s acquisition time and 10 accumulations. High resolution transmission electron microscopy (HRTEM), high angle annular dark field (HAADF) scanning TEM (STEM) and electron energy loss spectroscopy (EELS) spectrum imaging (S.I.) were performed using a TECNAI F20 operated at 200 kV with a point to point resolution of 0.19 nm, with a coupled GATAN QUANTUM EELS detector. HRTEM samples were prepared using a focused ion beam (FIB) to select the desired region to image, previous protecting the surface with deposited Pt. AFM and conductivity AFM (c-AFM) measures were taken with a Park Systems XE-100 with platinum conductive cantilevers biasing the sample at +1 V. Due to the p⁺n-Si built-in voltage, only the samples on p⁺-Si substrates were measured by c-AFM. The photoelectrochemical measurements (cyclic voltammetries, chronoamperometries and electrochemical impedance spectroscopy) were obtained with a Biologic VMP-300 potentiostat using Ag/AgCl/KCl (3M) ($E^0 = 0.203 \text{ V}_{\text{RHE}}$) as reference electrode and platinum mesh as counter electrode. A quartz cell with flat faces was used with 100 ml of 1M KOH electrolyte and a 300 W xenon lamp with an AM 1.5G filter adjusted to 100 mW/cm², calibrated using a silicon diode (Gentec-EO, XLPF12-3S-H2-DO). 20 mW·cm⁻² 365 nm UV light is obtained from a 200 W Hg-Xe Hamamatsu LC8 light source lamp. I-V curves were obtained in two-electrode configuration depositing 50 nm of Au as top contact.

3. RESULTS AND DISCUSSION

3.1. Morphological and structural characterization

In this work, we have fabricated silicon photoanodes protected with TiO_2 layers prepared by Atomic Layer Deposition (ALD) in a temperature range from 100 to 300 °C from TiCl_4 precursor, which permits a wide temperature deposition window.⁶⁰

Figure 1 shows SEM surface images of the ALD TiO_2 coatings synthesized at the different temperatures. For the film obtained at 100°C (Figure 1a), a homogeneous coating is observed with no presence of grains, whereas using deposition temperatures of 150 °C some regions with different contrast appear, corresponding to grains embedded in an amorphous matrix (Figure 1b). Further increasing deposition temperature up to 300 °C (Figure 1c), yields to a highly crystalline surface with smaller grain sizes. These grains are related to the TiO_2 anatase phase, as can be seen in the XRD diffraction pattern for the sample synthesized at 300°C (Figure 1d). The TiO_2 anatase phase is more favourable to OH^- adsorption and thus causing higher ALD growth rate.^{22,61} No anatase XRD peaks or other crystalline phases are detected at 100 or 150 °C, as previously observed in literature.^{53,62}

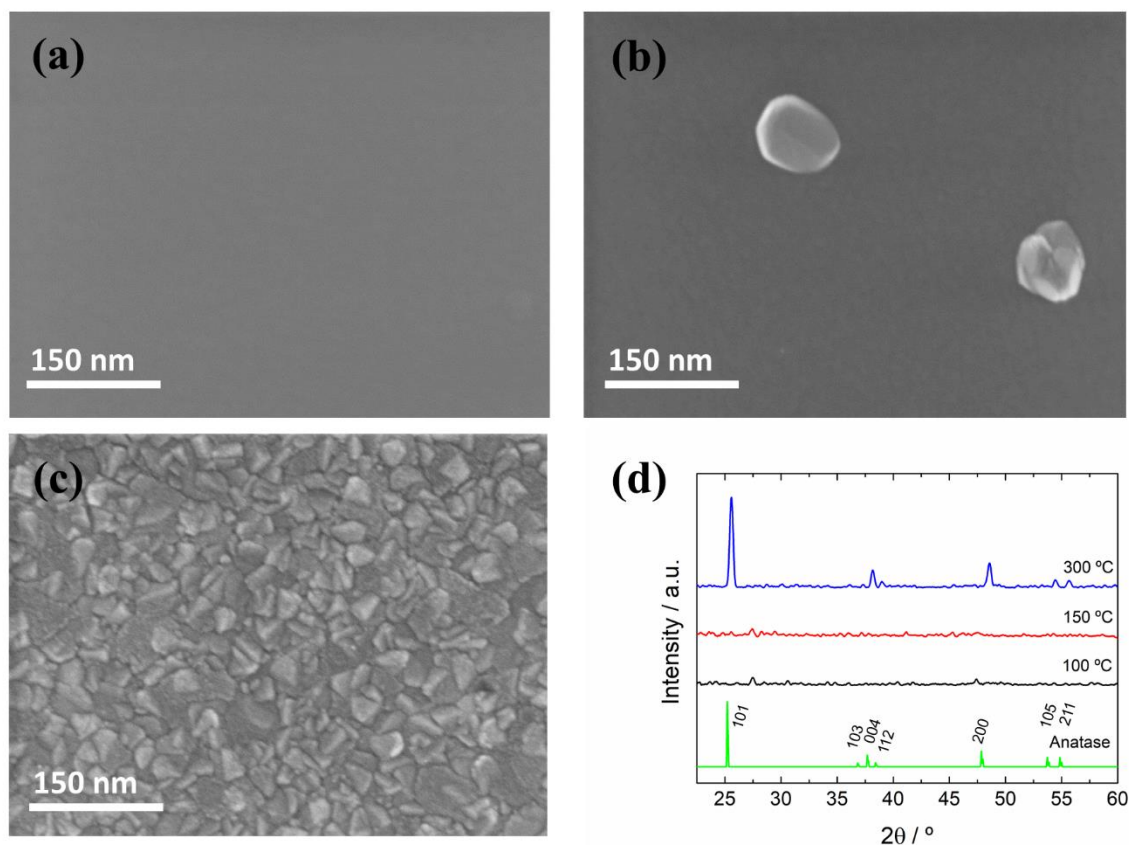


Figure 1. Top view SEM images of n⁺p-Si photoanodes protected with 100 nm TiO₂ grown at 100 °C (a), 150 °C (b) and 300 °C (c). (d) XRD spectra of TiO₂ layers grown at 100 °C (black), 150 °C (red) and 300 °C (blue) compared to the reference anatase pattern RRUFF R060277.9 (green).

As it can be seen in the SEM image for the sample synthesized at 150 °C (Figure 1b) the density of embedded crystals is rather low, and therefore XRD measurement does not allow the identification of the crystalline structure, being below the resolution. We used Raman spectroscopy equipped with an optical microscope, in order to carry out a more localised analysis. Optical view of the Si chip with an ALD-TiO₂ layer grown at 150 °C is shown in Figure 2a, where it is seen the presence of dark spots embedded in a homogenous grey film, which may be attributed to the submicron size crystals identified by SEM.⁶³ Figure 2b presents the Raman

spectra of two different spots, with a clear appearance of an intense peak in the darker spot, corresponding to the Eg mode of anatase located at 141.4 cm^{-1} . The peak is red-shifted with respect to the expected 144 cm^{-1} value,⁶⁴ being a sign of tensile stress, probably caused by partial amorphousness of the interface and reduced crystalline size.⁶⁵ A Raman mapping scan of the Eg main peak is shown in the inset of Figure 2a, which allows us to correlate the dark spots seen in the optical image with anatase crystallographic order presence.

Complete analysis by micro Raman spectroscopy of TiO₂ layers grown at 100, 150 and 300 °C is shown in Figure 2c. None of the spots analyzed of the sample synthesized at 100°C present the anatase peak, whereas is present at any spot at 300 °C samples, confirming the existence of a fully polycrystalline layer, in accordance with the XRD results and SEM. A deeper analysis of the Raman spectra shows a more intense peak for 300 °C layers compared with 150°C crystalline zones, and a blue shift of the Eg mode (of 10.3 cm^{-1}), that can be attributed to either surface pressure or phonon confinement effect that usually exist in nanometer-sized materials.⁶⁶ As seen in SEM images (Fig. 1c), nucleation is enhanced with increased deposition temperature and smaller grains are obtained by competitive growth.

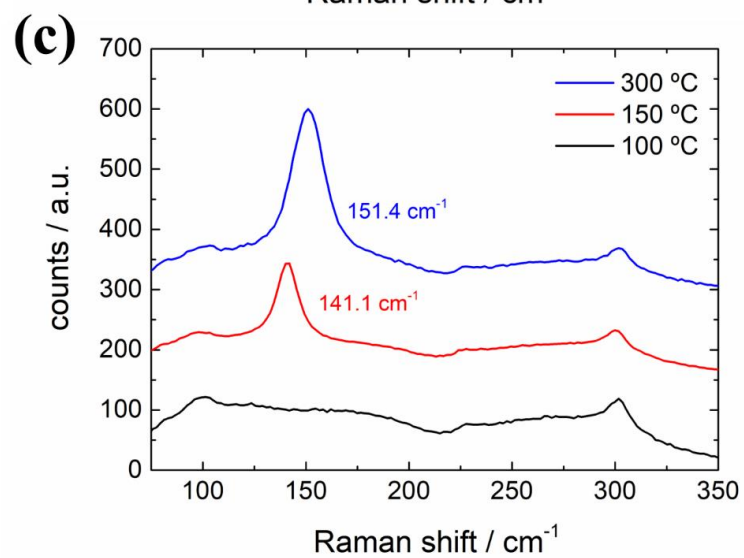
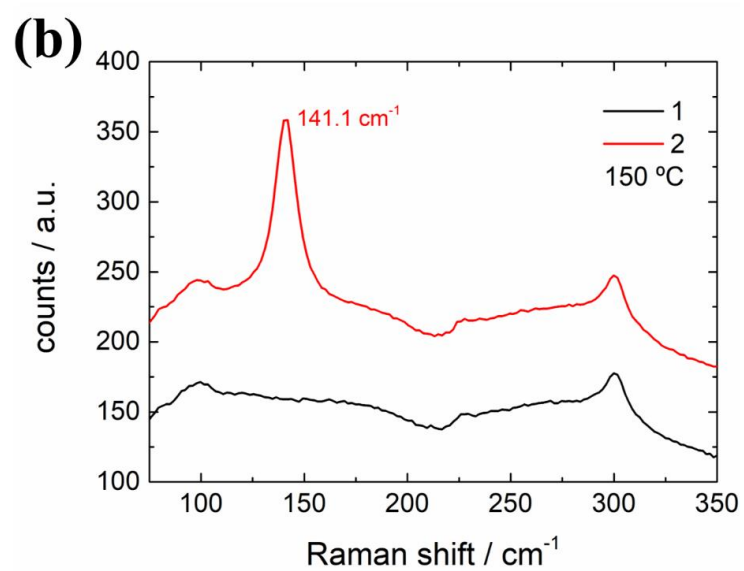
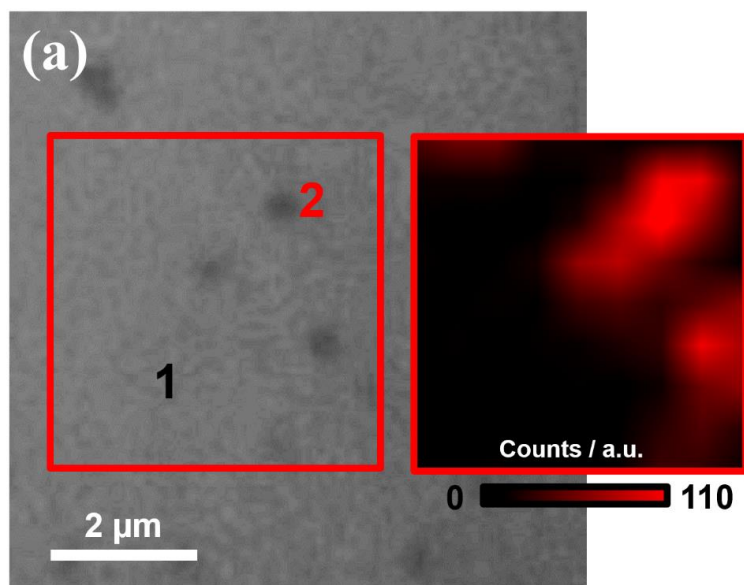


Figure 2. (a) Optical image of a protected silicon with TiO₂ grown by ALD at 150 °C. Right inset, counts map of the Eg Raman shift of the indicated region. (b) Raman shift spectra of two different positions of image a) corresponding to an amorphous region (1, black) or a crystalline grain (2, red). (c) Raman shift spectra of ALD-TiO₂ layers grown at 100 °C (black), 150 °C (red) and 300 °C (blue). The presented Raman shift spectra for the sample synthesized at 150 °C corresponds to a position where a TiO₂ grain can be seen. Peak maximum position is indicated for 150 and 300 °C.

3.2. *Electrochemical characterization*

In Figure 3a, cyclic voltammetries of three illuminated p⁺n-Si homojunctions are shown. The photoanodes, previously cleaned with HF (to eliminate the SiO_x) and immediately ALD-protected with TiO₂ grown at 100, 150 and 300 °C were then coated with a 5 nm film of thermally evaporated NiFe, to act as OER catalyst. For samples grown at 100 °C, no photocurrent can be observed. However, for samples deposited at 150°C and 300°C the same onset potential is obtained (around 0.96 V vs RHE) and whereas significant photocurrent is obtained for the former, highly increased currents are observed for the later. This fact, together with the saturation regime reached at more anodic potentials for the 150 °C sample, indicate a more resistive layer. The saturation current value, 22.5 mA/cm², is given by the photovoltaic quality of the prepared silicon-based homojunctions that were not optimized, and the lack of adjusted antireflective strategies.⁶⁷

In dark conditions, the same tendency observed for photoanodes is obtained in TiO₂-protected p⁺-Si anodes: increasing deposition temperature from 150 to 300 °C enhances OER current

(Figure 3b). The nickel-iron catalyst is necessary for efficient charge injection into the electrolyte^{28,68–71}, and if it is not present, no oxidative current is obtained in the studied range of potentials (dotted line in Figure 3b). In fact, Ni(Fe)OOH is the real catalytic phase obtained after cycling in alkaline electrolytes such as standard KOH solutions, one of the earth-abundant OER catalyst with lower overpotentials^{38,72–81}. The voltage difference between oxidative and reductive redox peaks of Ni hydroxide/oxyhydroxide is decreased when increasing TiO₂ deposition temperature (0.592 and 0.216 V for 150 and 300 °C). This fact is in accordance with the slopes of the OER current both in photoelectrodes (Figure 3a) and electrodes (Figure 3b), and it is related to higher series resistance in the system. Likewise, Ni²⁺/Ni³⁺ redox peaks present smaller area for the 150 °C sample than for the one synthesized at 300 °C, which indicates less active catalyst in the former sample, as the deposited NiFe quantity is the same in both samples. This fact may point out that charge transfer is boosted in the 300°C sample in comparison with the 150°C one.

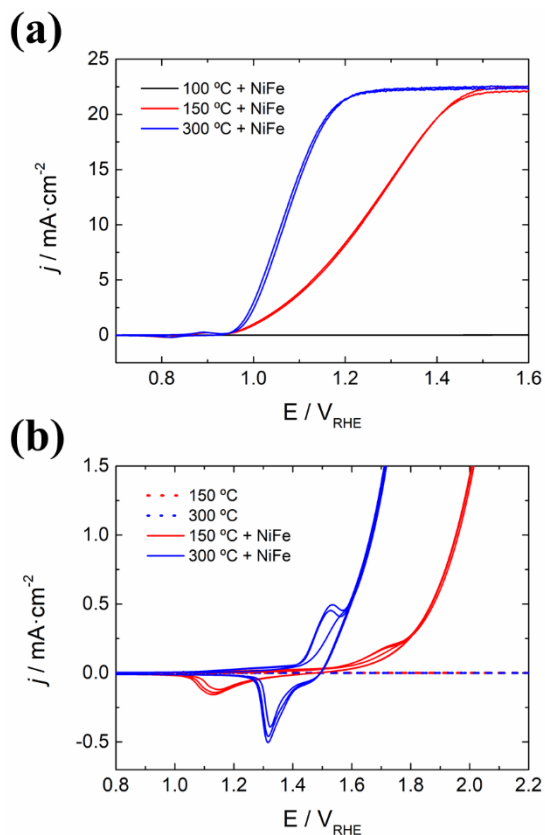


Figure 3. (a) Cyclic voltammeteries (1 sun AM 1.5G) of p^+n -Si photoanodes protected with 100 nm TiO_2 layers grown at 100 °C, 150°C and 300 °C, with 5 nm of nickel-iron thermally evaporated as OER catalyst. (b) Cyclic voltammeteries in dark of p^+ -Si substrates protected either with 150°C or 300 °C TiO_2 layers and decorated or not with 5 nm of thermally evaporated nickel-iron. 1 M KOH was used as electrolyte.

3.3. Conduction pathways in TiO_2 -ALD protected photoanodes

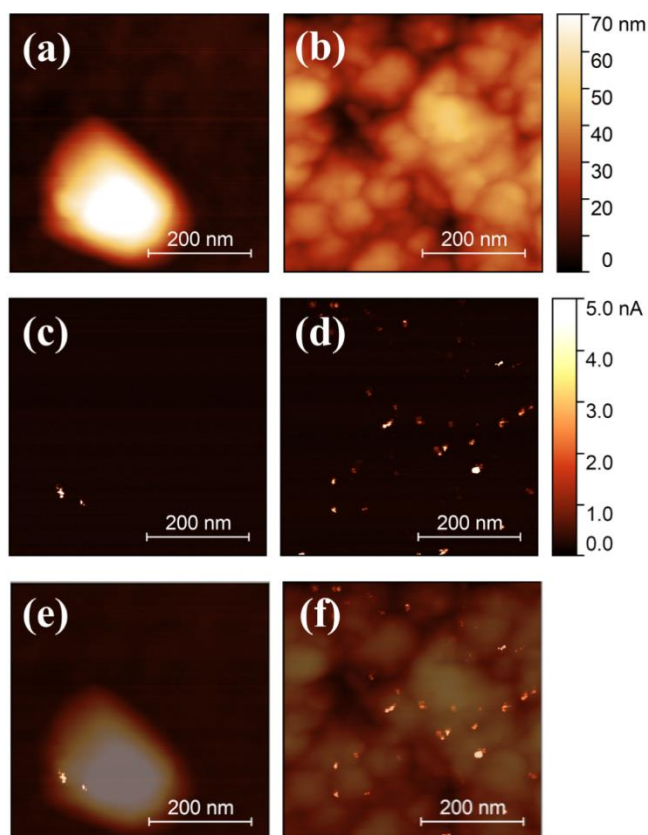


Figure 4. AFM height map of TiO₂ + NiFe layers grown at 150°C (a) and 300 °C (b), and its respective current intensity maps, (c) and (d) obtained at +1 V. (e) and (f) shows the overlays of the topography and current maps for each sample.

In Figure 4, we present the topography and conductivity atomic force microscopy maps (c-AFM) of ALD-TiO₂ samples grown on degenerately doped p⁺-Si substrates at 150 °C and 300 °C with 5 nm NiFe thermally evaporated on top. Measurements of TiO₂/NiFe layers were performed at +1 V, simulating the electron flow from the tip to the substrate across the layer, like in OER anodic conditions. Without NiFe catalyst, no current could be measured in the used range of potentials. As it can be seen in Figure 4 e-f, there is a correlation between topography maps and regions where current signal is detected: crystalline boundary regions, more abundant in the 300 °C sample than in the 150 °C, and no current in the amorphous regions. Additional images pointing

in the same direction are presented in supporting information Figure S2. Localized conduction of TiO₂-ALD layers observed by c-AFM has been also reported by other works. To explain the preferential conduction paths, Yu et. al.⁸² suggested the presence of crystallographically metastable regions, and Murakami et. al.⁸³ assumed an inhomogeneous distribution of oxygen vacancies (V_{ox}) known to be mobile under applied potentials and capable to agglomerate, defining preferential electron trap-assisted tunneling conductivity paths or conductive filaments, or Magneli phases in the grain boundaries such as it has been identified in the development of TiO₂ memristors.⁸⁴⁻⁸⁷

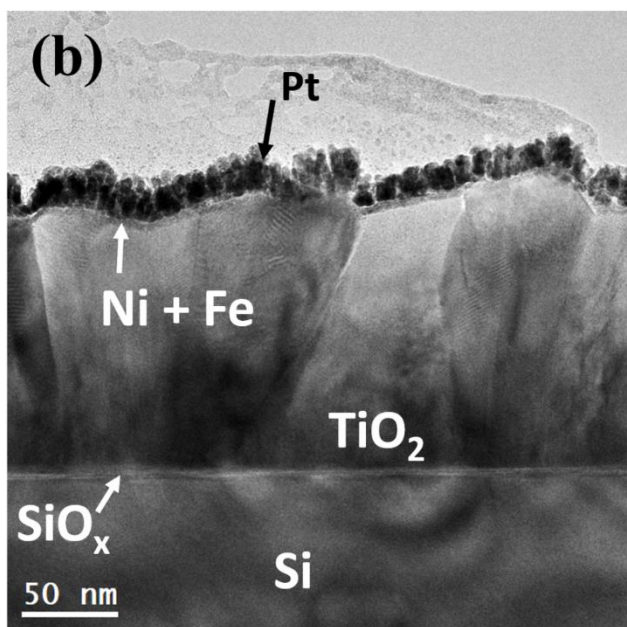
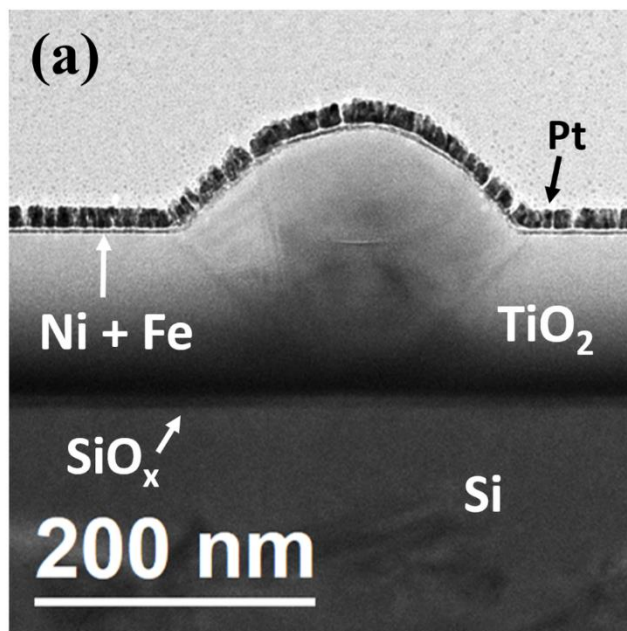


Figure 5. HRTEM cross sectional images of TiO₂ layers grown by ALD at 150 °C (a) and 300 °C (b). A Pt layer was deposited on top to protect the sample during the FIB sample preparation process.

Figure 5 shows the HRTEM images of the cross section obtained by FIB for the samples synthesized at 150 and 300°C, where multiple crystal orientations, grain boundaries and twin-like defects can be seen, embedded in an amorphous matrix for the 150 °C sample . Thus, unevenly distributed defects and oxygen vacancies should be expected, affecting the electrical states distribution. More HRTEM analysis is further depicted in Figure S3.

Columnar grains are present in both samples synthesized at 150 and 300 °C, but nucleation is significantly different. The HRTEM image of the sample grown at 150 °C shows that TiO₂ is crystalline only in a few-nanometer dome-like region close to the substrate, while the rest of the TiO₂ layer remains amorphous. However, the one prepared at 300 °C shows that the whole TiO₂ layer is crystalline. This can be attributed to the higher deposition temperature used, which enhances nucleation and lateral crystal propagation, together with crystal competition.⁴¹

The TiO₂/NiFe interface can be analysed by HRTEM EELS. In those as-prepared samples, the ~5nm apparently homogeneous NiFe catalyst layer can be seen (Figure S4a). After cycling in 1 M KOH (Figure S4b), a thicker layer can be observed, and the oxygen presence throughout the whole NiFe layer is increased, indicating oxidation of the metallic layer and OH⁻ incorporation, as reported in other works, where a restructuration of the NiFe towards a three dimensional nanoflakes porous layer takes place.⁸⁸ Moreover, there is no presence of Ni or Fe incorporation in the TiO₂ crystalline structure, or O or Ti concentration gradient can be detected above the EELS system resolution threshold, neither in the amorphous or crystalline structures nor the grain boundaries (Figures S5, S6).

To study the surface chemical composition of the layer, XPS analysis was performed, which is discussed in the supplementary information and Figure S7. To sum up, for as prepared samples, a

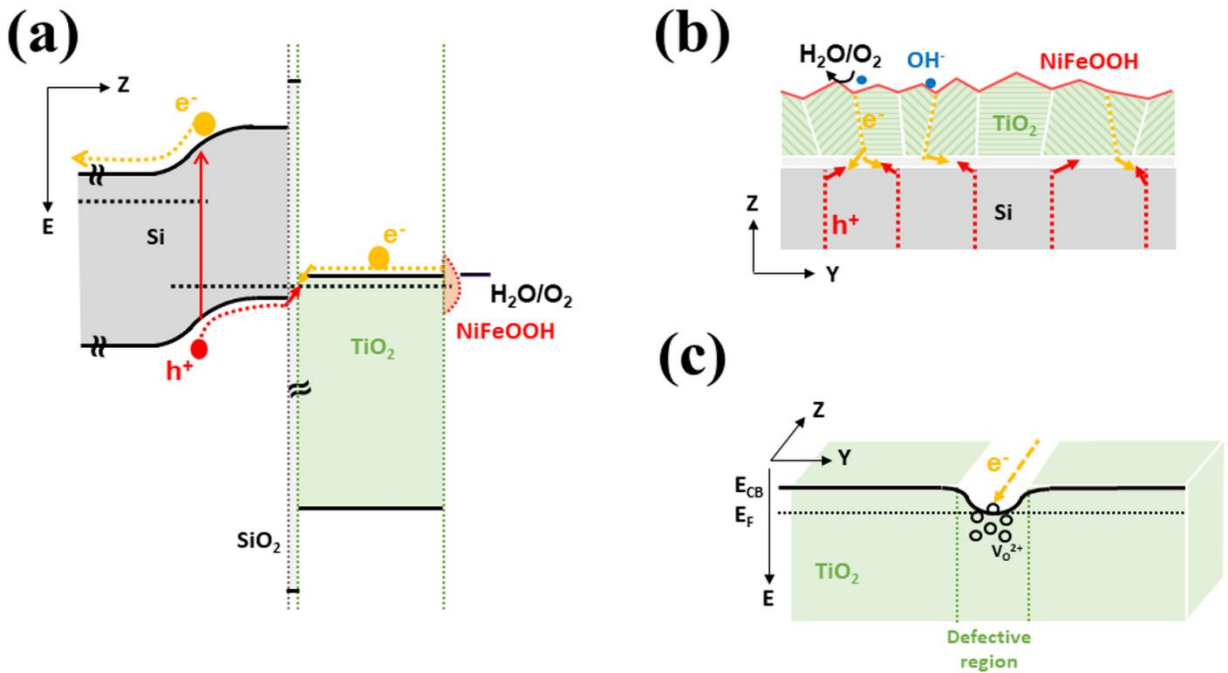
2.8:1 Ni:Fe ratio is observed, natively oxidized before anodic cycling and with higher oxygen content afterwards, caused by OH⁻ intercalation and oxidation of the coatings. Ni 2p spectra presents Ni(0) metallic peaks before but not after immersion in the electrolyte. For the cycled sample, multiple Ni(II) and Ni(III) peaks can be fitted (Figure 7a), probably by the presence of multiple chemical states such as NiO, Ni(OH)₂ and NiOOH, not easy to differentiate by XPS⁸⁹, and also correlated to Fe(III) forming Ni(Fe)OOH, as reported in other works^{38,69,88,90–92}.

XPS analysis also allows to obtain valence band (VB) positions of the superficial electrons, through the density and occupancy of electronic states of the surface (Figure S8). On one hand, as prepared TiO₂ layers at 150 and 300 °C present valence band potentials of ~2.75 eV under Fermi level, pointing at a higher n-type behaviour for the sample grown at 300 °C (larger Fermi level to valence band energies). On the other hand, no signal from intra-bandgap states is detected with a fine analysis of the 0-3 eV region in accordance with another works²⁸, although the small signal-to-noise ratio could be hindering the detection (Figure S9). Also, the oxidation of the metallic NiFe layer after immersion in the electrolyte at anodic potentials is revealed by the XPS spectra, still presenting slight free electrons behaviour. The most significant result is the excellent band alignment between TiO₂ conduction band and the Ni(Fe)OOH catalyst, that justify the injection of electrons from the oxygen evolution catalyst to the TiO₂ conduction band. The results obtained can be summarized in the mechanism proposed in Scheme 1, basing the conductivity in preferential charge transport pathways. The increased conductivity pathways density present for 300 °C deposited samples with respect to 150 °C ones (Figure 4 e-f) must be especially highlighted, in agreement with reported data in Figure 3. In addition, several features can be underlined:

1) The Si-TiO₂ interface has been found not to introduce significant electrical barrier to charge transport. A recombining contact for holes from the p⁺-Si valence band and electrons from the TiO₂ preferential paths conduction band is expected, facilitated by tunnelling through the ALD-caused SiO_x similar to the schemes proposed by other authors.⁹³

2) The transversal section to charge transport has been proven to be not homogeneous. Lower resistivity regions detected by c-AFM are attributed to more defective crystalline TiO₂ regions than bulk crystal grains, forming preferential conductivity paths through the grain boundaries and defects within the crystalline structure.⁸⁴

3) There is a favourable electron transfer from the Ni(Fe)OOH to the TiO₂ conduction band. This is facilitated by the permeable to electrolyte and energetically adaptive characteristics of the catalyst due to ionic diffusion, avoiding the energetic barrier formed between TiO₂ and the electrolyte and facilitating OER reaction.^{69,46}



Scheme 1. (a) Band diagram of a p⁺n-Si photoanode, TiO₂ protected with cycled Ni(Fe)OOH catalyst. (b) Cross section diagram with electron and holes flow paths depicted. c) Energy and cross section diagram corresponding to a preferential path or defective region, showing how extra states in interfacial defects generate a preferential path.

3.4. Endurance test: degradation mechanisms and stability

Crystalline TiO₂-protected p⁺-Si anodes and p⁺n-Si photoanodes were tested for stability in a 1 M KOH electrolyte at anodic potentials. The obtained current density values show significant decay for both electrodes (Figures 6 and 7). To further analyse these samples, Electrochemical Impedance Spectroscopy (EIS) was applied before and during stability test.

For the p⁺-Si anode in dark conditions, the current measured for more than 480 h at 1.8 V vs RHE, shows an exponential decay (Figure 6a), which indicates a degradation of the electrode. The EIS measurement shows two clear semicircles increasing over time (Figure 6b) revealing degradation correlated with two processes in series defined by different RC time constants, which allows us to propose an equivalent electrical circuit shown in Figure 6b. As seen in Figure S10, the first semicircle (R_2/C_2) is independent from the applied potential and the second one varies significantly between open circuit potential (OCP) and OER working conditions, proving a non-ohmic dependence of this second one. Our electrode has been modelled to an ohmic resistance, R_1 (corresponding to all resistive system contributions given by cables, back contact and substrate and electrolyte conductivity), a R_2/C_2 parallel circuit in series (corresponding to the depletion region within the TiO₂ protective layer) and a Randles circuit (R_3/CPE_3 , corresponding to the electrocatalytic OER charge transfer on Ni(Fe)OOH).

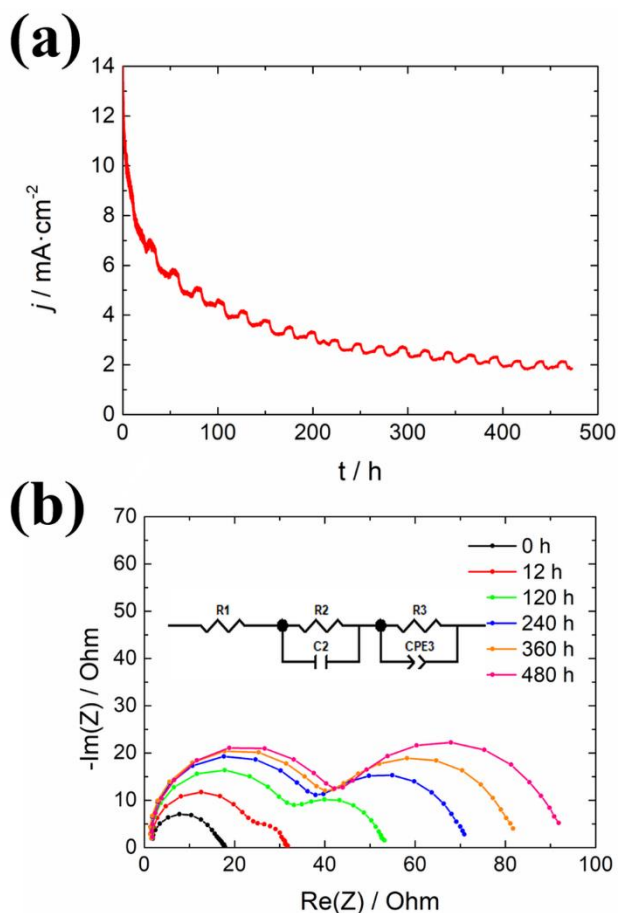


Figure 6. (a) Chronoamperometry of a $\text{p}^+\text{-Si}$ anode protected with ALD-deposited TiO_2 at 300 °C and NiFe decorated in 1 M KOH at a fixed potential of 1.8 V vs RHE. (fluctuations are due to daily temperature changes). (b) Nyquist plot of the electrochemical impedance spectroscopy (EIS) data from the same sample obtained every 12 h under potentiostatic 1.8 V vs RHE potential from 1 Hz to 50 kHz and an amplitude of 25mV. The inset shows the equivalent electric circuit used to fit the EIS data.

Fitting the proposed EIS circuit to our experiments (Table 1), a clear relation is observed between the current loss during stability tests and the semicircles increase. The current diminution rate is initially high, and it is reduced as current values decrease, pointing towards a relation between current passing and conductivity degradation. The resistance R_2 corresponding to the current flow

through the TiO₂ layer increases logarithmically over time at 1.8 V vs RHE, meanwhile R₃, the resistance attributed to charge transfer in the OER reaction, increases lineally (Figure S11).

Table 1. Values obtained after fitting the data presented in Figure 6b according to the circuit proposed.

Time (h)	R ₁ (Ω)	R ₂ (Ω)	C ₂ (F·cm ⁻²)	R ₃ (Ω)	CPE ₃ (F·cm ⁻² ·s ^{α₃-1})	α ₃
0	1.40	13.56	2.11·10 ⁻⁶	2.83	6.54·10 ⁻⁴	0.787
12	1.33	22.58	1.80·10 ⁻⁶	7.30	1.22·10 ⁻⁴	0.992
120	1.22	31.32	1.83·10 ⁻⁶	19.84	1.22·10 ⁻⁴	0.955
240	1.03	37.04	1.84·10 ⁻⁶	31.53	1.14·10 ⁻⁴	0.955
360	0.92	39.30	1.79·10 ⁻⁶	40.40	1.18·10 ⁻⁴	0.940
480	1.29	40.37	1.67·10 ⁻⁶	50.45	1.41·10 ⁻⁴	0.901

Owing to the fact that the interfacial SiO_x layer does not increase in thickness after stability test (Figure S12), the increase of R₂ over time can be related with the intrinsic self-passivating process of metal oxide surfaces, following a logarithmic rate law.⁹⁴ This is highlighting an interaction of OH⁻ ions from the electrolyte and the polycrystalline TiO₂ driven by the electric field due to anodic polarization. Some authors have previously reported chemical interaction of OH⁻ migration and defective regions, nulling preferential charge transport paths.⁹⁵

The significant decrease of crystalline TiO₂ conductivity has not previously been stated, but a deep review of the bibliography does not deny its plausibility. Other works have shown the possibility of OH⁻ to penetrate through a permeable Ni catalyst layer and diffuse through the TiO₂⁸² or ALD-TiO₂ dissolution on few hours scale⁹⁵. Oxygen mobility under polarization in TiO₂ has been known for years, forming preferential paths in the form of conductive filaments.^{96,97} The presence of such oxidant molecules, OH⁻, inside the TiO₂ structure can easily interact

with the states responsible of the n-type semiconductor properties and conductivity.⁹⁸ This fact has been depicted in Scheme S1a and b.

On the other hand, according to the Tafel equation describing the charge transfer at the electrode/electrolyte interface (eq. 1)⁹⁹, overpotential (η) depends logarithmically on the current density. Assuming both elements are in series (i_2 equals i_3 , i), and the resistance of the TiO₂ layer (R_2) increases while applied potential is kept constant, the current reaching the OER reaction (measured in EIS as R_3) will decrease logarithmically, giving an inverse logarithmic dependence of R_3 with R_2 (eq. 2), represented in Figure S13. A 5 times increase in the total resistance ($R_2 + R_3$) is in accordance with the equivalent current reduction observed after the 480 h stability test under anodic polarization.

$$\log\left(\frac{1}{R_{ct}}\right) = \frac{di}{d\eta} = a * \eta + b; \quad (1)$$

where $R_{ct} = R_3$ and $\eta = V - i * R_2$, giving:

$$\log\left(\frac{1}{R_3}\right) \propto (-R_2) \quad (2)$$

Following these equations, if $R_2 \propto \log(t)$ as it has been demonstrated in Figure S11, and substituting in eq. 2, we obtain $R_3 \propto t$.

In his turn, capacitance of the TiO₂ semicircle is not significantly modified as expected, as it corresponds to the thickness and dielectric constant of the TiO₂ and these are maintained.

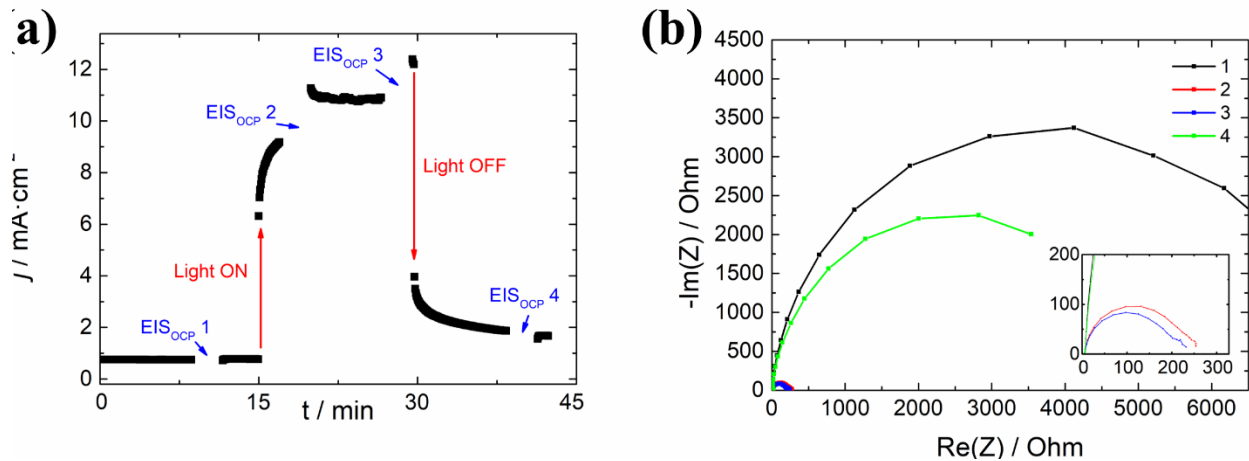


Figure 7. a) Chronoamperometry measurement at 1.3 V vs RHE of a p^+n -Si photoanode protected with TiO_2 ALD-deposited at 300 °C and NiFe under solar irradiation and superimposed UV light, after 333 h of stability test b) EIS measurement at OCP of a UV temporarily illuminated degraded photoanode, partially recovering the photocurrent via light-induced states. Measurements performed under 1 sun illumination, using 1 M KOH as electrolyte.

Similar stability test was performed with a p^+n -Si photoanode with TiO_2 deposited at 300 °C and NiFe decorated. Likewise, after 333 h at 1.3 V vs RHE photocurrent has been reduced significantly (Figure S14).

To validate the previously described degradation mechanism, superimposed UV illumination was used to modify the density of states within the TiO_2 electronic structure, achieving a significant recovery of the photocurrent (Figure 7a). This can be explained by the photoconductivity effect on TiO_2 ,¹⁰⁰ where the UV photons create electron-hole pairs changing the occupation of localized states, causing a significant reduction of R_2 (Figure 7b, Table S1).¹⁰¹ Switching off UV light gave a few-seconds decay followed by a several minutes persistent photoconductivity (PPC). These states created by oxygen desorption will eventually be

compensated again by OH^- diffusing inside the material, as illustrated in Scheme S1. This UV-light effect should only be expected in the near-surface region of TiO_2 due to small UV photon penetration depth, precisely where higher OH^- interaction would be.

4. CONCLUSIONS

TiO_2 protective layers grown by ALD have been deposited on silicon photoanodes to avoid corrosion, while allowing charge transfer to the electrolyte to perform the oxygen evolution reaction using a NiFe earth abundant catalyst. Layers consisting of approximately 100 nm have been grown at 100, 150 and 300 °C, corresponding to completely amorphous, partially crystalline and fully crystalline TiO_2 .

Completely amorphous layers presented high resistivity, meanwhile the ones synthesized at 300°C, having a fully anatase crystalline TiO_2 structure, introduced insignificant resistance, showing direct correlation between crystallization degree and electrical conductivity. The conductivity through crystalline TiO_2 layers has been found not to be homogeneous, presenting preferential conduction paths attributed to grain boundaries and defects within the crystalline structure. A correlation between the c-AFM measurements and grain interstitials can be seen, supported by HRTEM cross sectional images presenting defective regions in crystalline TiO_2 grains.

A conduction mechanism has been proposed assuming electrons coming from water oxidation are injected into the TiO_2 conduction band in preferential regions through the electrocatalyst. Then, electrons are transported to the $\text{Si}/\text{SiO}_x/\text{TiO}_2$ interface where electrons recombine with holes given by the p^+n -Si junction. No evidences of intra band gap states in TiO_2 responsible of conductivity have been detected.

On the other hand, current decay under continuous polarization was found, where OH^- ions are considered to diffuse into TiO_2 under anodic conditions and alter conduction paths conductivity. These changes were analysed under long-time (> 480 h) stability measurements at fixed voltage, which presented current decay over time. Deeper EIS analysis show two different time constant RC elements changing during endurance test, and allowed to track electrical variations in a series circuit structure. Associated TiO_2 resistance increases logarithmically over time, as expected for metal oxide surface passivation, corroborating a self-limited degradation. A 5 times increase on the EIS measured resistances associated to TiO_2 conductivity and OER catalysis is in accordance with the total current reduction observed after 480 h stability test in anodic polarization.

To corroborate this model, UV illumination was used to modify states occupancy, presenting persistent conductivity thanks to the partial recovery of the electronic states lost under working conditions. Further studies focused on higher resolution analysis than this work are needed to identify OH^- diffusion into TiO_2 structure and states loss, specifically in preferential regions. UV illumination constitutes a promising method to partially recover degradation. This confirms that the degradation mechanisms of TiO_2 still needs to be taken into consideration to enhance its long term use endurance as a transparent conductive and protective layer for frontal illumination water splitting photoanodes.

ASSOCIATED CONTENT

AFM, HRTEM, XPS, and EIS supplementary data, together with a photo of a finished device and a scheme of UV recovery, are supplied in an additional document

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