

# EXISTENCE OF PERIODIC SOLUTIONS FOR A CLASS OF SECOND ORDER ORDINARY DIFFERENTIAL EQUATIONS

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**ABSTRACT.** We provide sufficient conditions for the existence of a periodic solution for a class of second order differential equations of the form  $\ddot{x} + g(x) = \varepsilon f(t, x, \dot{x}, \varepsilon)$ , where  $\varepsilon$  is a small parameter.

## 1. INTRODUCTION AND STATEMENT OF THE RESULTS

The second order differential equations of the form

$$\ddot{x} + g(x) = \varepsilon f(t, x, \dot{x}, \varepsilon),$$

have been studied by many authors because they have many applications, see for instance [2, 3, 7, 8, 9, 12, 15]. Two of the main families studied are the Duffing equations see [5, 6], ... or the forced pendulum see the nice survey [11] and the references quoted therein.

The aim of this work is to study periodic solutions of the second order differential equation

$$(1) \quad \ddot{x} + g(x) = \mu^{2n+1} p(t) + \mu^{4n+1} q(t, x, y, \mu),$$

where  $n$  is a positive integer,  $\mu$  is a small parameter, and the functions

$$g(x) = x + x^{2n+1} (b + xh(x)),$$

and  $h(x)$  are smooth,  $b \neq 0$ ,  $p(t)$  and  $q(t, x, y, \mu)$  are smooth and periodic with period  $2\pi$  in the variable  $t$ .

Let  $\Gamma(x)$  the Gamma function, see for more details [1], and let  $\alpha$  and  $\beta$  the first Fourier coefficients of the periodic function  $p(t)$ , i.e.

$$\alpha = \frac{1}{\pi} \int_0^{2\pi} p(t) \cos t \, dt, \quad \beta = \frac{1}{\pi} \int_0^{2\pi} p(t) \sin t \, dt.$$

Then our main result is the following.

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**Theorem 1.** *If  $\alpha\beta \neq 0$  then for  $\mu \neq 0$  sufficiently small the differential equation (1) has a  $2\pi$ -periodic solution  $\mathbf{x}(t, \mu)$  such that*

$$\mathbf{x}(0, \mu) = \pi^{\frac{1}{4n+2}} \left( \frac{\Gamma(n+2)}{2b\Gamma(n+\frac{3}{2})} \right)^{\frac{1}{2n+1}} \alpha \left( \frac{\beta^2}{\alpha^2} + 1 \right)^{-n} + O(\mu^{2n}).$$

Theorem 1 is proved in section 3. Its proof uses the averaging theory for computing periodic solutions, see section 2 for a summary of the results on this theory that we shall need.

## 2. THE AVERAGING THEORY

We want to study the  $T$ -periodic solutions of the periodic differential systems of the form

$$(2) \quad \mathbf{x}' = F_0(t, \mathbf{x}) + \varepsilon F_1(t, \mathbf{x}) + \varepsilon^2 F_2(t, \mathbf{x}, \varepsilon),$$

with  $\varepsilon > 0$  sufficiently small, where  $F_0, F_1 : \mathbb{R} \times \Omega \rightarrow \mathbb{R}^n$  and  $F_2 : \mathbb{R} \times \Omega \times (-\varepsilon_0, \varepsilon_0) \rightarrow \mathbb{R}^n$  are  $\mathcal{C}^2$  functions,  $T$ -periodic in the variable  $t$ , and  $\Omega$  is an open subset of  $\mathbb{R}^n$ . Let  $\mathbf{x}(t, \mathbf{z}, \varepsilon)$  be the solution of the differential system (2) such that  $\mathbf{x}(0, \mathbf{z}, \varepsilon) = \mathbf{z}$ . Suppose that the unperturbed system

$$(3) \quad \mathbf{x}' = F_0(t, \mathbf{x}),$$

has an open set  $V$  with  $\bar{V} \subset \Omega$  such that for each  $\mathbf{z} \in \bar{V}$ ,  $\mathbf{x}(t, \mathbf{z}, 0)$  is  $T$ -periodic.

Let  $\mathbf{y}$  be an  $n \times n$  matrix, and consider the first order variational equation

$$(4) \quad \mathbf{y}' = D_{\mathbf{x}} F_0(t, \mathbf{x}(t, \mathbf{z}, 0)) \mathbf{y},$$

of the unperturbed system (3) on the periodic solution  $\mathbf{x}(t, \mathbf{z}, 0)$ . Let  $M_{\mathbf{z}}(t)$  be the fundamental matrix of the linear differential system (4) with periodic coefficients such that  $M_{\mathbf{z}}(0)$  is the  $n \times n$  identity matrix.

**Theorem 2.** *Consider the function  $F : \bar{V} \rightarrow \mathbb{R}^n$*

$$(5) \quad f(\mathbf{z}) = \int_0^T M_{\mathbf{z}}^{-1}(t) F_1(t, \mathbf{x}(t, \mathbf{z}, 0)) dt.$$

*If there exists  $\alpha \in V$  with  $f(\alpha) = 0$  and*

$$(6) \quad \det((df/d\mathbf{z})(\alpha)) \neq 0,$$

*then there exists a  $T$ -periodic solution  $\mathbf{x}(t, \varepsilon)$  of system (2) such that  $\mathbf{x}(0, \varepsilon) = \alpha + O(\varepsilon)$ .*

The existence of the periodic solution of Theorem 2 is due to Malkin [10] and Roseau [13], for a shorter and easier proof see [4]. The proof for the stability follows in a similar way to the proof of Theorem 11.6 of [14].

### 3. PROOF OF THEOREM 1

The differential equation of second order (1) can be written as the first order differential system

$$(7) \quad \begin{aligned} \dot{x} &= y, \\ \dot{y} &= -x - x^{2n+1} (b + xh(x)) + \mu^{2n+1} p(t) + \mu^{4n+1} q(t, x, y, \mu). \end{aligned}$$

In order to apply the averaging theory described in section 2 to this differential system we do the scaling  $x \rightarrow \mu x$  and  $y \rightarrow \mu y$ . Hence the differential system (7) becomes

$$(8) \quad \begin{aligned} \dot{x} &= y, \\ \dot{y} &= -x + \mu^{2n} (-bx^{2n+1} + p(t)) + \mu^{4n} q^*(t, x, y, \mu). \end{aligned}$$

This system is written into the normal form (2) for applying the averaging theory described in section 2, where

$$(9) \quad \begin{aligned} \mathbf{x} &= (x, y), \\ \varepsilon &= \mu^{2n}, \\ \mathbf{F}_0(\mathbf{x}) &= (y, -x), \\ \mathbf{F}_1(\mathbf{x}, t) &= (0, -bx^{2n+1} + p(t)), \\ \mathbf{F}_2(\mathbf{x}, t, \varepsilon) &= (0, \bar{q}(t, x, y, \varepsilon)). \end{aligned}$$

From section 2 the solution  $\mathbf{x}(t, \mathbf{z}, 0) = (x(t, \mathbf{z}, 0), y(t, \mathbf{z}, 0))$  of system (8) with  $\varepsilon = 0$  satisfies  $\mathbf{x}(0, \mathbf{z}, 0) = \mathbf{z} = (x_0, y_0)$ , and consequently

$$\begin{aligned} x(t, \mathbf{z}, 0) &= x_0 \cos t + y_0 \sin t, \\ y(t, \mathbf{z}, 0) &= -x_0 \sin t + y_0 \cos t. \end{aligned}$$

The fundamental matrix  $M_{\mathbf{z}}(t) = M(t)$  of the the first order variational equation (4) satisfying (9) is

$$M(t) = \begin{pmatrix} \cos t & \sin t \\ -\sin t & \cos t \end{pmatrix}.$$

According with Theorem 2 in order to compute the  $2\pi$ -periodic solutions of the differential system (8) we must compute the integral

$$\begin{aligned} f(\mathbf{z}) &= \begin{pmatrix} f_1(x_0, y_0) \\ f_2(x_0, y_0) \end{pmatrix} \\ &= \int_0^{2\pi} M^{-1}(t) F_1(t, \mathbf{x}(t, \mathbf{z}, 0)) dt \\ &= \begin{pmatrix} b \int_0^{2\pi} \sin t (x_0 \cos t + y_0 \sin t)^{2n+1} dt - \int_0^{2\pi} p(t) \sin t dt \\ -b \int_0^{2\pi} \cos t (x_0 \cos t + y_0 \sin t)^{2n+1} dt + \int_0^{2\pi} p(t) \cos t dt \end{pmatrix}. \end{aligned}$$

Doing induction with respect to  $n$  it is not difficult to show that

$$\begin{aligned} \int_0^{2\pi} \sin t (x_0 \cos t + y_0 \sin t)^{2n+1} dt &= \frac{2\sqrt{\pi} \Gamma(\frac{3}{2} + n)}{\Gamma(2 + n)} y_0 (x_0^2 + y_0^2)^n, \\ \int_0^{2\pi} \cos t (x_0 \cos t + y_0 \sin t)^{2n+1} dt &= \frac{2\sqrt{\pi} \Gamma(\frac{3}{2} + n)}{\Gamma(2 + n)} x_0 (x_0^2 + y_0^2)^n. \end{aligned}$$

Therefore we must solve the system

$$\begin{pmatrix} f_1(x_0, y_0) \\ f_2(x_0, y_0) \end{pmatrix} = \begin{pmatrix} \frac{2\sqrt{\pi} b \Gamma(\frac{3}{2} + n)}{\Gamma(2 + n)} y_0 (x_0^2 + y_0^2)^n - \pi \beta_1 \\ -\frac{2\sqrt{\pi} b \Gamma(\frac{3}{2} + n)}{\Gamma(2 + n)} x_0 (x_0^2 + y_0^2)^n + \pi \alpha_1 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}.$$

This system has a unique solution

$$\begin{pmatrix} x_0^* \\ y_0^* \end{pmatrix} = \pi^{\frac{1}{4n+2}} \left( \frac{\Gamma(n+2)}{2b\Gamma(n+\frac{3}{2})} \right)^{\frac{1}{2n+1}} \begin{pmatrix} \alpha \left( \frac{\beta^2}{\alpha^2} + 1 \right)^{-n} \\ \beta \left( \frac{\alpha^2}{\beta^2} + 1 \right)^{-n} \end{pmatrix}.$$

The determinant (6) of the Jacobian matrix  $Df(x_0^*, y_0^*)$  is

$$\begin{aligned} \det(Df(x_0^*, y_0^*)) &= 4^{\frac{1}{2n+1}} (2n+1) \pi^{\frac{2n}{2n+1}+1} \left( \frac{\Gamma(n+2)}{b\Gamma(n+\frac{3}{2})} \right)^{-\frac{2}{2n+1}} \\ &\quad \left( \left( \beta \left( \frac{\alpha^2}{\beta^2} + 1 \right)^{-n} \right)^{\frac{2}{2n+1}} + \left( \alpha \left( \frac{\beta^2}{\alpha^2} + 1 \right)^{-n} \right)^{\frac{2}{2n+1}} \right)^{2n}, \end{aligned}$$

and by assumptions it is positive because  $\alpha\beta b \neq 0$ .

In summary all the assumptions of Theorem 2 hold and consequently from Theorem 2 it follows Theorem 1.

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#### REFERENCES

- [1] M. ABRAMOWITZ AND I.A. STEGUN, *Modified Bessel Functions I and K*, §9.6 in Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables, 9th printing. New York: Dover, 1972, pp 374–377.
- [2] D.R. ANDERSON, *Multiple periodic solutions for a second-order problem on periodic time scales*, Nonlinear Anal. **60** (2005), 101–115.
- [3] D.R. ANDERSON AND R.I. AVERY, *Existence of a periodic solution for continuous and discrete periodic second-order equations with variable potentials*, J. Appl. Math. Comput. **37** (2011), 297–312.
- [4] A. BUICA, J. P. FRANÇOISE AND J. LLIBRE, *Periodic solutions of nonlinear periodic differential systems with a small parameter*, Comm. Pure Appl. Anal. **6** (2006), 103–111.
- [5] H. CHEN AND Y. LI, *Stability and exact multiplicity of periodic solutions of Duffing equations with cubic nonlinearities*, Proc. Amer. Math. Soc. **135** (2007), 3925–3932.
- [6] G. DUFFING, *Erzwungen Schwingungen bei veränderlicher Eigenfrequenz und ihre technisch Bedeutung*, Sammlung Vieweg Heft 41/42, Vieweg, Braunschweig, 1918.
- [7] J.R. GRAEF, L. KONG AND H. WANG, *Existence, multiplicity, and dependence on a parameter for a periodic boundary value problem*, J. Differ. Equ. **245** (2008), 1185–1197.
- [8] J.K. HALE AND P.Z. TABOAS, *Interaction of damping and forcing in a second order equation*, Nonlinear Anal. **2** (1978), 77–84.
- [9] Y. LIU, W. GE AND Z. GUI, *Three positive periodic solutions of nonlinear differential equations with periodic coefficients*, Anal. Appl. **3**(2) (2005), 145–155.
- [10] I.G. MALKIN, *Some Problems of the theory of nonlinear oscillations*, Gosudarstv. Izdat. Tehn-Teor. Lit. Moscow, 1956 (in Russian).
- [11] J. MAWHIN, *Seventy-five years of global analysis around the forced pendulum equation*, Proceedings of Equadiff 9 CD rom, Masaryk University, Brno 1997, pp 115–145.
- [12] D. NUÑEZ, A. RIVERA AND G. ROSSODIVITA, *Stability of odd periodic solutions in a resonant oscillator*, Annali di Matematica (2017) **196**, 443–455.
- [13] M. ROSEAU, *Vibrations non linéaires et théorie de la stabilité*, Springer Tracts in Natural Philosophy, Vol. **8**, Springer, New York, 1985.
- [14] F. VEHRULST, *Nonlinear differential equations and dynamical systems*, Universitext, Springer, 1996.

- [15] X. WU, J. LI AND Y. ZHOU, *A priori bounds for periodic solutions of a Duffing equation*, J. Appl. Math. Comput. **26** (2008), 535-543.

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