

LIMIT CYCLES FOR TWO CLASSES OF CONTROL PIECEWISE LINEAR DIFFERENTIAL SYSTEMS

JAUME LLIBRE¹, REGILENE D. OLIVEIRA² AND CAMILA AP. B. RODRIGUES³

ABSTRACT. We study the bifurcation of limit cycles from the periodic orbits of $2n$ -dimensional linear centers when they are perturbed inside classes of continuous and classes of discontinuous piecewise linear differential systems of control theory of the form $\dot{x} = Ax + \phi(x_1)b$, where ϕ is a continuous or discontinuous piecewise linear function respectively.

1. INTRODUCTION AND STATEMENT OF THE MAIN RESULTS

In the control theory are relevant the *continuous piecewise linear differential systems* of the form

$$(1) \quad \dot{x} = Ax + \varphi(x_1)b,$$

with A a $m \times m$ matrix, $x, b \in \mathbb{R}^m$, $\varphi : \mathbb{R} \rightarrow \mathbb{R}$ is the continuous piecewise linear function

$$(2) \quad \varphi(x_1) = \begin{cases} -1 & \text{if } x_1 \in (-\infty, -1), \\ x_1 & \text{if } x_1 \in [-1, 1], \\ 1 & \text{if } x_1 \in (1, \infty), \end{cases}$$

where $x = (x_1, \dots, x_m)^T$, and the dot denotes the derivative with respect to the independent variable t , the time.

Also in control theory are important the *discontinuous piecewise linear differential systems* of the form (1) where $\psi : \mathbb{R} \rightarrow \mathbb{R}$ is the discontinuous piecewise linear function

$$(3) \quad \psi(x_1) = \begin{cases} -1 & \text{if } x_1 \in (-\infty, 0), \\ 1 & \text{if } x_1 \in (0, \infty). \end{cases}$$

For more details on these continuous and discontinuous piecewise linear differential systems see for instance the books [1, 3, 10, 11, 15, 17].

The analysis of discontinuous piecewise linear differential systems goes back mainly to Andronov and coworkers [2] and nowadays still continues to receive attention by many researchers. In particular, discontinuous piecewise linear differential systems appear in

a natural way in control theory and in the study of electrical circuits, see for instance the book [4] and the references quoted there. These systems can present complicated dynamical phenomena such as those exhibited by general nonlinear differential systems.

One of the main ingredients in the qualitative description of the dynamical behavior of a differential system is the number and the distribution of its limit cycles. The goal of this paper is to study the existence of limit cycles for a class of continuous and a class of discontinuous piecewise linear differential of the form (1).

More precisely, first we consider the class of continuous piecewise linear differential systems

$$(4) \quad \dot{x} = A_0 x + \varepsilon (Ax + \varphi(x_1)b),$$

with $|\varepsilon| \neq 0$ a sufficiently small real parameter, where A_0 is the $2n \times 2n$ matrix having on its principal diagonal the following 2×2 matrices

$$\begin{pmatrix} 0 & -(2k-1) \\ 2k-1 & 0 \end{pmatrix} \quad \text{for } k = 1, \dots, n,$$

in the order that they are defined and zeros in the complement, A is an arbitrary $2n \times 2n$ matrix and $b \in \mathbb{R}^{2n} \setminus \{0\}$. Note that for $\varepsilon = 0$ system (4) becomes

$$(5) \quad \dot{x}_1 = -x_2, \quad \dot{x}_2 = x_1, \quad \dots, \quad \dot{x}_{2n-1} = -(2n-1)x_{2n}, \quad \dot{x}_{2n} = (2n-1)x_{2n-1}.$$

Moreover, the origin of (5) is a *global isochronous center* in \mathbb{R}^{2n} , i.e. all its orbits different from the origin are periodic with period 2π .

In a similar way we consider the discontinuous piecewise linear differential systems

$$(6) \quad \dot{x} = A_0 x + \varepsilon (Ax + \psi(x_1)b).$$

Our main results on the limit cycles of the continuous and discontinuous piecewise linear differential systems (4) are the following ones.

Theorem 1. *For $|\varepsilon| > 0$ sufficiently small and using averaging theory of first order at most one limit cycle of the continuous piecewise linear differential system (4) bifurcates from the periodic orbits of system (5). Moreover there are systems (4) with $|\varepsilon| > 0$ sufficiently small having a such limit cycle.*

Theorem 1 is proved in section 3.

Theorem 2. *For $|\varepsilon| > 0$ sufficiently small and using averaging theory of first order at most one limit cycle of the discontinuous piecewise linear differential system (6) bifurcates from the periodic orbits of system (5). Moreover there are systems (4) with $|\varepsilon| > 0$ sufficiently small having a such limit cycle.*

Theorem 2 is proved in section 4.

If instead of the matrix A_0 we consider the matrix A_1 where A_1 is the $2n \times 2n$ matrix having on its principal diagonal the following 2×2 matrices

$$\begin{pmatrix} 0 & -k \\ k & 0 \end{pmatrix} \quad \text{for } k = 1, \dots, n,$$

in the order that they are defined and zeros in the complement, then we have the following results.

Theorem 3. *For $|\varepsilon| > 0$ sufficiently small the averaging theory of first order does not provide any information about the limit cycles of the continuous piecewise linear differential system*

$$(7) \quad \dot{x} = A_1 x + \varepsilon(Ax + \varphi(x_1)b),$$

which can bifurcate from the periodic orbits of system (5).

Theorem 4. *For $|\varepsilon| > 0$ sufficiently small the averaging theory of first order does not provide any information about the limit cycles of the discontinuous piecewise linear differential system*

$$(8) \quad \dot{x} = A_1 x + \varepsilon(Ax + \psi(x_1)b),$$

which can bifurcate from the periodic orbits of system (5).

Theorems 3 and (4) are proved in section 5.

Note the difference between the matrices A_0 and A_1 , in the matrix A_0 the non-zero entries are only the odd numbers $1, 3, \dots, 2n - 1$, while in the matrix A_1 the non-zero entries are the numbers $1, 2, \dots, n$. This difference provides that the continuous and discontinuous piecewise linear differential systems (4) and (6) can have limit cycles detected by the averaging theory, while for the continuous and discontinuous piecewise linear differential systems (7) and (8) the averaging theory cannot detect limit cycles.

Really for the control differential systems here studied the limit cycles that we obtain bifurcate from some periodic orbit of the $2n$ -dimensional linear differential center (5). This technique of finding limit cycles bifurcating from centers has been intensively studied in dimension 2, see for instance the book of Christopher and Li [9] and the hundreds of references quoted therein.

Other results different to the ones presented here, but which also study the limit cycles of control systems of the form (1) using averaging theory, can be found in [6, 7, 8, 12].

The main tools for proving the previous theorems are the extensions of the classical averaging theory for computing periodic solutions of \mathcal{C}^2 differential systems to continuous and discontinuous differential systems. In section 2 we summarize the extensions of the averaging theory that we shall use here for proving our results.

2. FIRST ORDER AVERAGING THEORY

For the classical averaging theory for finding periodic orbits of differential systems of class \mathcal{C}^2 see for instance the chapter 11 of the book of Verhulst [16].

In this section we present first the result on the continuous averaging theory that we will use for proving our Theorems 1 and 3. This theory uses the Brouwer degree of a continuous function and its proof can be found in [5].

Theorem 5. *We consider the following differential system*

$$(9) \quad \dot{x} = \varepsilon H(t, x) + \varepsilon^2 R(t, x, \varepsilon),$$

where $H : \mathbb{R} \times D \rightarrow \mathbb{R}^n$, $R : \mathbb{R} \times D \times (-\varepsilon_f, \varepsilon_f) \rightarrow \mathbb{R}^n$ are continuous functions, T -periodic in the first variable, and D is an open bounded subset of \mathbb{R}^n . We define $h : D \rightarrow \mathbb{R}^n$ as

$$(10) \quad h(z) = \int_0^T H(s, z) ds,$$

and assume that

- (i) H and R are locally Lipschitz with respect to x ;
- (ii) for $p \in D$ with $h(p) = 0$, there exists a neighborhood V of p such that $h(z) \neq 0$ for all $z \in \bar{V} \setminus \{p\}$ and the Brouwer degree $d_B(h, V, 0) \neq 0$.

Then, for $|\varepsilon| \neq 0$ sufficiently small, there exists an isolated T -periodic solution $x(t, \varepsilon)$ of system (9) such that $x(0, \varepsilon) \rightarrow p$ as $\varepsilon \rightarrow 0$.

Remark 6. Let $h : D \rightarrow \mathbb{R}^n$ be a C^1 function with $h(p) = 0$, where D is an open bounded subset of \mathbb{R}^n and $p \in D$. If the Jacobian of h at p is not zero, then there exists a neighborhood V of p such that $h(z) \neq 0$ for all $z \in \bar{V} \setminus \{p\}$, and the Brouwer degree $d_B(h, V, p) \in \{-1, 1\}$.

For a proof of Remark 6 see for instance [14].

For proving Theorems 2 and 4 we shall need the following extension of the averaging theory for computing periodic solutions to discontinuous differential systems done in [13].

Theorem 7. *We consider the following discontinuous differential system*

$$(11) \quad x'(t) = \varepsilon H(t, x) + \varepsilon^2 R(t, x, \varepsilon),$$

with

$$\begin{aligned} H(t, x) &= H_1(t, x) + \text{sign}(g(t, x))H_2(t, x), \\ R(t, x, \varepsilon) &= R_1(t, x, \varepsilon) + \text{sign}(g(t, x))R_2(t, x, \varepsilon), \end{aligned}$$

where $H_1, H_2 : \mathbb{R} \times D \rightarrow \mathbb{R}^n$, $R_1, R_2 : \mathbb{R} \times D \times (-\varepsilon_0, \varepsilon_0) \rightarrow \mathbb{R}^n$ and $g : \mathbb{R} \times D \rightarrow \mathbb{R}$ are continuous functions, T -periodic in the variable t and D is an open subset of \mathbb{R}^n . We also suppose that g is a C^1 function having 0 as a regular value.

Define the average function $h : D \rightarrow \mathbb{R}^n$ as

$$(12) \quad h(x) = \int_0^T H(t, x) dt.$$

We assume the following conditions.

- (i) H_1, H_2, R_1, R_2 are locally Lipschitz with respect to x ;
- (ii) there exists an open bounded subset $C \subset D$ such that, for $|\varepsilon| > 0$ sufficiently small, every orbit starting in \overline{C} reaches the set of discontinuity only at its crossing regions.
- (iii) for $a \in C$ with $h(a) = 0$, there exists a neighbourhood $U \subset C$ of a such that $h(z) \neq 0$ for all $z \in \overline{U} \setminus \{a\}$ and $d_B(h, U, 0) \neq 0$.

Then, for $|\varepsilon| > 0$ sufficiently small, there exists a T -periodic solution $x(t, \varepsilon)$ of system (11) such that $x(0, \varepsilon) \rightarrow a$ as $\varepsilon \rightarrow 0$.

3. PROOF OF THEOREM 1

The main tool for proving Theorem 1 is the averaging theory of first order for continuous differential systems presented in Theorem 5. In order to use this theorem we need to write the differential system (4) in the normal form (9), and for obtaining this we need to some changes of variables.

Lemma 8. *Doing the change of variables $(x_1, x_2, \dots, x_{2n}) \mapsto (\theta, r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1})$ defined by*

$$\begin{aligned} x_1 &= r \cos \theta, \\ x_2 &= r \sin \theta, \\ x_{2j-1} &= r_{j-1} \cos((2j-1)\theta + \theta_{j-1}), \\ x_{2j} &= r_{j-1} \sin((2j-1)\theta + \theta_{j-1}), \end{aligned}$$

for $j = 2, \dots, n$ system (4) is transformed into the system

$$(13) \quad \begin{aligned} \frac{dr}{d\theta} &= \varepsilon H_1(\theta, r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}) + \mathcal{O}(\varepsilon^2), \\ \frac{dr_{j-1}}{d\theta} &= \varepsilon H_{2(j-1)}(\theta, r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}) + \mathcal{O}(\varepsilon^2), \\ \frac{d\theta_{j-1}}{d\theta} &= \varepsilon H_{2j-1}(\theta, r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}) + \mathcal{O}(\varepsilon^2), \end{aligned}$$

where

$$H_1 = \sum_{l=1}^n r_{l-1} \left(F_{1,l} \cos \theta + F_{2,l} \sin \theta \right) + \varphi(r \cos \theta)(b_1 \cos \theta + b_2 \sin \theta),$$

and for $j = 2, 3, \dots, n$ we have

$$\begin{aligned}
H_{2(j-1)} &= \sum_{l=1}^n r_{l-1} \left(F_{2j-1,l} \cos((2j-1)\theta + \theta_{j-1}) + F_{2j,l} \sin((2j-1)\theta + \theta_{j-1}) \right) \\
&\quad + \varphi(r \cos \theta) [b_{2j-1} \cos((2j-1)\theta + \theta_{j-1}) + b_{2j} \sin((2j-1)\theta + \theta_{j-1})], \\
H_{2j-1} &= \sum_{l=1}^n \frac{r_{l-1}}{r_{j-1}} \left(F_{2j,l} \cos((2j-1)\theta + \theta_{j-1}) - F_{2j-1,l} \sin((2j-1)\theta + \theta_{j-1}) \right) \\
&\quad + (2j-1) \sum_{l=1}^n \frac{r_{l-1}}{r} \left(F_{1,l} \sin \theta - F_{2,l} \cos \theta \right) \\
&\quad + \varphi(r \cos \theta) \left(\frac{b_{2j}}{r_{j-1}} \cos((2j-1)\theta + \theta_{j-1}) - \frac{b_{2j-1}}{r_{j-1}} \sin((2j-1)\theta + \theta_{j-1}) \right) \\
&\quad - (2j-1) \varphi(r \cos \theta) \left(\frac{b_2}{r} \cos \theta - \frac{b_1}{r} \sin \theta \right),
\end{aligned}$$

with

$$F_{i,l} = F_{i,l}(r, \theta, \theta_{l-1}) = a_{i(2l-1)} \cos((2l-1)\theta + \theta_{l-1}) + a_{i(2l)} \sin((2l-1)\theta + \theta_{l-1}).$$

We take ε_0 sufficiently small, m arbitrarily large and

$$D_m = \left\{ (r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}) \in \left(\frac{1}{m}, m \right) \times \left[\mathbb{S}^1 \times \left(\frac{1}{m}, m \right) \right]^{n-1} \right\}.$$

Then the vector field of system (13) is well defined and continuous on $\mathbb{S}^1 \times D_m \times (-\varepsilon_0, \varepsilon_0)$. Moreover the system is 2π -periodic with respect to variable θ and locally Lipschitz with respect to variables $(r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1})$.

Proof. In the variables $(\theta, r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1})$ the differential system (4) becomes

$$\begin{aligned}
\dot{\theta} &= 1 + \frac{\varepsilon}{r} \left[\sum_{l=1}^n r_{l-1} \left(F_{2,l} \cos \theta - F_{1,l} \sin \theta \right) + \varphi(r \cos \theta) (b_2 \cos \theta - b_1 \sin \theta) \right], \\
\dot{r} &= \varepsilon H_1(\theta, r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}), \\
\dot{r}_{j-1} &= \varepsilon H_{2(j-1)}(\theta, r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}), \\
\dot{\theta}_{j-1} &= \varepsilon H_{2j-1}(\theta, r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}),
\end{aligned}$$

for $j = 2, 3, \dots, n$. Note that for $\varepsilon = 0$, $\dot{\theta}(t) > 0$ and hence for $|\varepsilon| \neq 0$ sufficiently small this property remains valid for each t when $(\theta, r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}) \in \mathbb{S}^1 \times D_m$. Now we take θ as the new independent variable. The right-hand side of the new system is well defined and continuous in $\mathbb{S}^1 \times D_m \times (-\varepsilon_0, \varepsilon_0)$ and it is 2π -periodic with respect to the new variable θ and locally Lipschitz with respect to $(r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1})$. Now system (8) can be obtained doing a Taylor series expansion in the parameter ε around $\varepsilon = 0$. \square

The next step is to find the corresponding average function (10) of system (8) that we denoted by $h = (h_1, h_2, \dots, h_{2(n-1)}, h_{2n-1}) : D_m \rightarrow \mathbb{R}^{n-1}$ and it is defined by

$$\begin{aligned} h_1 &= h_1(r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}) = \int_0^{2\pi} H_1(r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}) d\theta, \\ h_{2(j-1)} &= h_{2(j-1)}(r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}) = \int_0^{2\pi} H_{2(j-1)}(r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}) d\theta, \\ h_{2j-1} &= h_{2j-1}(r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}) = \int_0^{2\pi} H_{2j-1}(r, \theta_1, r_1, \dots, \theta_{n-1}, r_{n-1}) d\theta, \end{aligned}$$

for $j = 1, 2, \dots, n$. To calculate these integrals we will use the following equalities

$$\begin{aligned} \int_0^{2\pi} \cos((2j-1)\theta + \theta_{j-1}) \sin((2l-1)\theta + \theta_{l-1}) d\theta &= 0 \quad \text{for all integers } l, j > 1, \\ \int_0^{2\pi} \cos((2j-1)\theta + \theta_{j-1}) \cos((2l-1)\theta + \theta_{l-1}) d\theta &= \begin{cases} \pi & \text{if } l = j, \\ 0 & \text{if } l \neq j, \end{cases} \\ \int_0^{2\pi} \sin((2j-1)\theta + \theta_{j-1}) \sin((2l-1)\theta + \theta_{l-1}) d\theta &= \begin{cases} \pi & \text{if } l = j, \\ 0 & \text{if } l \neq j, \end{cases} \end{aligned}$$

and the next lemma.

For $r > 0$ and $j = 1, 2, \dots, n$ we denote

$$\begin{aligned} I_j(r) &= \int_0^{2\pi} \varphi(r \cos \theta) \cos((2j-1)\theta) d\theta, \\ J_j(r) &= \int_0^{2\pi} \varphi(r \cos \theta) \sin((2j-1)\theta) d\theta, \end{aligned}$$

where φ is the piecewise linear function (2).

Lemma 9. *The integrals I_j and $J_j(r)$ satisfy*

$$\begin{aligned} I_j(r) &= \begin{cases} \pi r & \text{if } j = 1 \text{ and } 0 < r \leq 1, \\ 0 & \text{if } j > 1 \text{ and } 0 < r \leq 1, \\ K(r) & \text{if } j = 1 \text{ and } r > 1, \\ L_j(r) & \text{if } j > 1 \text{ and } r > 1; \end{cases} \\ J_j(r) &= 0 \quad \text{for all } j = 1, 2, \dots, n \text{ and } r > 0. \end{aligned}$$

where

$$\begin{aligned} L_j(r) &= \frac{2}{j(2j-1)^2} \left((2j-1) \sqrt{-1+r^2} \cos((2j-1) \arctan \sqrt{-1+r^2}) \right. \\ &\quad \left. - \sin((2j-1) \arctan \sqrt{-1+r^2}) \right), \\ K(r) &= \pi r + \frac{2}{r} \sqrt{r^2-1} - 2r \arctan(\sqrt{r^2-1}). \end{aligned}$$

Proof. We consider two cases: $0 < r \leq 1$ and $r > 1$.

Case 1: $0 < r \leq 1$ In this case $|r \cos \theta| \leq 1$ and hence $\varphi(r \cos \theta) = r \cos \theta$ for all $\theta \in [0, 2\pi]$. Then if $j = 1$

$$\int_0^{2\pi} \varphi(r \cos \theta) \cos \theta d\theta = r \int_0^{2\pi} \cos^2 \theta d\theta = \pi r,$$

and

$$\int_0^{2\pi} \varphi(r \cos \theta) \sin \theta d\theta = r \int_0^{2\pi} \cos \theta \sin \theta d\theta = 0.$$

And if $j > 1$ then

$$\int_0^{2\pi} \varphi(r \cos \theta) \cos((2j-1)\theta) d\theta = r \int_0^{2\pi} \cos \theta \cos((2j-1)\theta) d\theta = 0,$$

$$\int_0^{2\pi} \varphi(r \cos \theta) \sin((2j-1)\theta) d\theta = r \int_0^{2\pi} \cos \theta \sin((2j-1)\theta) d\theta = 0.$$

Case 2: $r > 1$ In this case choose $\theta_c \in (0, \pi/2)$ such that $\cos \theta_c = 1/r$. If $j = 1$ we have

$$\begin{aligned} I_1(r) &= \int_0^{\theta_c} \cos \theta d\theta + r \int_{\theta_c}^{\pi-\theta_c} \cos^2 \theta d\theta - \int_{\pi-\theta_c}^{\pi+\theta_c} \cos \theta d\theta \\ &\quad + r \int_{\pi+\theta_c}^{2\pi-\theta_c} \cos^2 \theta d\theta + \int_{2\pi-\theta_c}^{2\pi} \cos \theta d\theta \\ &= \pi r + \frac{2}{r} \sqrt{r^2 - 1} - 2r \arctan(\sqrt{r^2 - 1}). \end{aligned}$$

The same reasoning can be applied to see that $J_1(r) = 0$. If $j > 1$ then

$$\begin{aligned} I_j(r) &= \int_0^{\theta_c} \cos((2j-1)\theta) d\theta + r \int_{\theta_c}^{\pi-\theta_c} \cos \theta \cos((2j-1)\theta) d\theta - \int_{\pi-\theta_c}^{\pi+\theta_c} \cos((2j-1)\theta) d\theta \\ &\quad + r \int_{\pi+\theta_c}^{2\pi-\theta_c} \cos \theta \cos((2j-1)\theta) d\theta + \int_{2\pi-\theta_c}^{2\pi} \cos((2j-1)\theta) d\theta \\ &= \frac{2}{j(2j-1)^2} \left((2j-1) \sqrt{-1+r^2} \cos((2j-1) \arctan \sqrt{-1+r^2}) \right. \\ &\quad \left. - \sin((2j-1) \arctan \sqrt{-1+r^2}) \right), \end{aligned}$$

and $J_j(r) = 0$. □

With the results presented previously we are able to prove Theorem 1. Since we can choose m sufficiently large to find the zeroes of the average function h in D_m it is sufficient to look for them in $(0, \infty) \times [\mathbb{S}^1 \times (0, \infty)]^{n-1}$. To calculate the expression of the average function we consider again two cases.

Case 1: $0 < r \leq 1$. In this case the system whose zeros can provide limit cycles of system (4) is

$$\begin{aligned}
 h_1 &= (a_{11} + a_{22} + b_1)\pi r, \\
 h_2 &= (a_{33} + a_{44})\pi r_1, \\
 h_3 &= (a_{43} - a_{34} + 3(a_{12} - a_{21} - b_2))\pi, \\
 &\vdots \\
 h_{2(n-1)} &= (a_{(2n-1)(2n-1)} + a_{(2n)(2n)})\pi r_{n-1}, \\
 h_{2n-1} &= (a_{(2n)(2n-1)} - a_{(2n-1)(2n)} + (2n-1)(a_{12} - a_{21} - b_2))\pi.
 \end{aligned}
 \tag{14}$$

Note that the variables $\theta_1, \theta_2, \dots, \theta_{n-1}$ does not appear explicitly into system (14). Hence, if this system has zeros, it has a continuum of zeros. Therefore the assumption (ii) of the averaging theory, presented in Theorem 5, is not satisfied and this theorem does not provide any information about the limit cycles of system (13).

Case 2: $r > 1$. Now the system whose zeros can provide limit cycles of system (13) is

$$\begin{aligned}
 h_1 &= (a_{11} + a_{22})\pi r + b_1 K(r), \\
 h_2 &= (a_{33} + a_{44})\pi r_1 + (b_3 \cos \theta_1 + b_4 \sin \theta_2)L_2(r), \\
 h_3 &= (a_{43} - a_{34} + 3(a_{12} - a_{21}))\pi - \\
 &\quad \frac{3b_2 r_1 K(r) - r(b_4 \cos \theta_1 - b_3 \sin \theta_1)L_2(r)}{rr_1}, \\
 &\vdots \\
 h_{2(n-1)} &= (a_{(2n-1)(2n-1)} + a_{(2n)(2n)})\pi r_{n-1} + \\
 &\quad (b_{2n-1} \cos \theta_{n-1} + b_{2n} \sin \theta_{n-1})L_n(r), \\
 h_{2n-1} &= (a_{(2n)(2n-1)} - a_{(2n-1)(2n)} + (2n-1)(a_{12} - a_{21}))\pi - \\
 &\quad \frac{(2n-1)b_2 r_{n-1} K(r) - r(b_{2n} \cos \theta_{n-1} - b_{2n-1} \sin \theta_{n-1})L_n(r)}{rr_{n-1}},
 \end{aligned}
 \tag{15}$$

For each $j \in \{2, 3, \dots, n\}$ we will study the zeros of the system

$$\begin{aligned} h_1 &= (a_{11} + a_{22})\pi r + b_1 K(r), \\ h_{2(j-1)} &= (a_{(2j-1)(2j-1)} + a_{(2j)(2j)})\pi r_{j-1} + \\ &\quad (b_{2j-1} \cos \theta_{j-1} + b_{2j} \sin \theta_{j-1})L_j(r), \\ h_{2j-1} &= \frac{(a_{(2j)(2j-1)} - a_{(2j-1)(2j)} + (2j-1)(a_{12} - a_{21}))\pi - \\ &\quad (2j-1)b_{2j}r_{j-1}K(r) - r(b_{2j} \cos \theta_{j-1} - b_{2j-1} \sin \theta_{j-1})L_j(r)}{rr_{j-1}}, \end{aligned}$$

Claim: *The function $K : (1, \infty) \rightarrow (\pi, 4)$ is a diffeomorphism.* Indeed note that K is twice differentiable with

$$K'(r) = \pi - 2 \frac{\sqrt{r^2 - 1}}{r^2} - 2 \arctan \sqrt{r^2 - 1},$$

and

$$K''(r) = -\frac{4}{r^3 \sqrt{r^2 - 1}} < 0$$

which implies that K' is a strictly decreasing function. Moreover $\lim_{r \rightarrow \infty} K'(r) = 0$ what means that $K'(r)$ has a horizontal asymptote given by the axis r and then $K'(r) \geq 0$. Suppose that there exists an $r_0 \in (1, \infty)$ such that $K'(r_0) = 0$. Then for all $r > r_0$ we have $K'(r) < K'(r_0) = 0$, contradiction. Therefore it follows that $K'(r) \neq 0$ for all $r \in (1, \infty)$ and the Inverse Function Theorem guarantees that K is a local diffeomorphism and since that K is a injective function we obtain the global diffeomorphism, ending the proof of this claim.

First we note that in order that the equation $h_1 = 0$ has solutions with $r > 1$ it is necessary that $b_1(a_{11} + a_{22}) < 0$. Moreover $K''(r) < 0$ implies that the graph of K is convex. In the plane of the graph of $K(r)$ the graph of $(a_{11} + a_{22})\pi r$ is a straight line passing through the origin and then both graphs can intersect at most in two points.

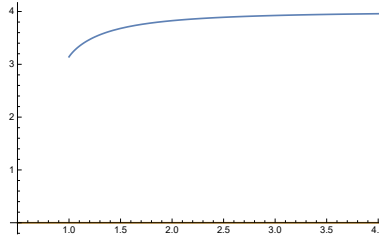


FIGURE 1. The graphic of the function $K(r)$.

But if some straight line intercept the graph of $K(r)$ in two points then it cannot pass through the origin, as we can see in Figure 1. Then the equation $h_1 = 0$ has at

most one solution if $r > 1$, and since that $K(r)$ is a diffeomorphism we can choose the coefficients a_{11}, a_{22} and b_1 so that this solution exists. We denote this solution by r_0 and we substitute it into the equations $h_{2(j-1)} = 0$ and $h_{2j-1} = 0$. Defining

$$\begin{aligned} A_j &= (a_{(2j-1)(2j-1)} + a_{(2j)(2j)})\pi, & B_j &= b_{2j-1}L_j(r_0), & C_j &= b_{2j}L_j(r_0), \\ D_j &= (a_{(2j)(2j-1)} - a_{(2j-1)(2j)} + (2j-1)(a_{12} - a_{21}))\pi - \frac{1}{r_0}(2j-1)b_2K(r_0), \\ u_j &= \cos \theta_{j-1}, & v_j &= \sin \theta_{j-1}, \end{aligned}$$

the system $h_{2(j-1)} = h_{2j-1} = 0$ is equivalent to the system

$$\begin{aligned} A_j r_{j-1} + B_j u_j + C_j v_j &= 0, \\ D_j r_{j-1} + C_j u_j - B_j v_j &= 0, \\ u_j^2 + v_j^2 - 1 &= 0. \end{aligned}$$

Using the two first equations we obtain

$$u_j = -\frac{(A_j B_j + C_j D_j) r_{j-1}}{B_j^2 + C_j^2}, \quad v_j = \frac{(B_j D_j - A_j C_j) r_{j-1}}{B_j^2 + C_j^2}.$$

Substituting these two expressions in the third equation we get

$$(A_j^2 + D_j^2) r_{j-1}^2 - B_j^2 - C_j^2 = 0.$$

Therefore at most there is one solution $r_{j-1} > 0$, which provide a unique u_j and v_j . Since we fixed an arbitrarily j to solve this system, the same reasoning can be applied to each pair of equations $h_{2(j-1)} = 0$ and $h_{2j-1} = 0$, concluding that system (15) has at most one solution. Moreover taking conveniently the parameters of the initial system (4) this solution exists and its Jacobian is not zero. Hence at most one limit cycle can bifurcate from the periodic orbits of the center of system (5) when we perturb it as in system (4), and there are systems for which a such limit cycles exist. This completes the proof of Theorem 1.

Now we present an explicit example of a continuous piecewise linear differential system (4) in \mathbb{R}^4 , and repeating for it the proof of Theorem 1 we will show it has one limit cycle. Consider the following differential system

$$(16) \quad \dot{x} = A_0 x + \varepsilon(Ax + \varphi(x_1)b),$$

where

$$A_0 = \begin{pmatrix} 0 & -1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -3 \\ 0 & 0 & 3 & 0 \end{pmatrix}, A = \begin{pmatrix} 2 & 1 & 0 & 0 \\ 0 & 2 & 0 & 0 \\ 0 & 0 & -2 & -1 \\ 0 & 0 & 0 & \frac{18\pi - \sqrt{3}}{9\pi} \end{pmatrix}, b = \begin{pmatrix} -\frac{24\pi}{3\sqrt{3} + 2\pi} \\ 1 \\ \frac{9(3 - 2\sqrt{3}\pi)}{2} \\ -1 \end{pmatrix}.$$

Doing the change of variables $x_1 = r \cos \theta$, $x_2 = r \sin \theta$, $x_3 = r_1 \cos(3\theta + \theta_1)$, $x_4 = r_1 \sin(3\theta + \theta_1)$ and taking θ as the new independent variable we obtain the system (17)

$$\begin{aligned} r'(\theta) &= \frac{dr}{d\theta} = \varepsilon H_1(\theta, r, \theta_1, r_1) + \mathcal{O}(\varepsilon^2), \\ r_1'(\theta) &= \frac{dr_1}{d\theta} = \varepsilon H_2(\theta, r, \theta_1, r_1) + \mathcal{O}(\varepsilon^2), \\ \theta_1'(\theta) &= \frac{d\theta_1}{d\theta} = \varepsilon H_3(\theta, r, \theta_1, r_1) + \mathcal{O}(\varepsilon^2), \\ H_1(\theta, r, \theta_1, r_1) &= 2r + \varphi(r \cos \theta) \sin \theta + \cos \theta \left(r \sin \theta - \frac{24\pi \varphi(r \cos \theta)}{3\sqrt{3} + 2\pi} \right), \\ H_2(\theta, r, \theta_1, r_1) &= -\frac{1}{18\pi} \left(18\pi \varphi(r \cos \theta) \sin(3\theta + \theta_1) + 9\pi r_1 \sin(2(3\theta + \theta_1)) \right. \\ &\quad \left. + \sqrt{3}r_1 + 81\pi(2\sqrt{3}\pi - 3)\varphi(r \cos \theta) \cos(3\theta + \theta_1) - \right. \\ &\quad \left. (\sqrt{3} - 36\pi)r_1 \cos(2(3\theta + \theta_1)) \right), \\ H_3(\theta, r, \theta_1, r_1) &= \sin^2(3\theta + \theta_1) + 2 \sin(2(3\theta + \theta_1)) - \frac{\sin(2(3\theta + \theta_1))}{6\sqrt{3}\pi} + 3 \sin^2 \theta \\ &\quad - \frac{72\pi \varphi(r \cos \theta) \sin \theta}{3\sqrt{3}r + 2\pi r} - \frac{\varphi(r \cos \theta) \cos(3\theta + \theta_1)}{r_1} - \frac{3\varphi(r \cos \theta) \cos \theta}{r} \\ &\quad + \frac{9\pi\sqrt{3}\varphi(r \cos \theta) \sin(3\theta + \theta_1)}{r_1} - \frac{27\varphi(r \cos \theta) \sin(3\theta + \theta_1)}{2r_1}. \end{aligned}$$

After some computations the average function $h = (h_1, h_2, h_3)$ defined in (10) is

$$\begin{aligned} h_1(r, \theta_1, r_1) &= 4\pi r - \frac{24\pi}{3\sqrt{3} + 2\pi} \left(\pi r + \frac{2\sqrt{r^2 - 1}}{r} - 2r \arctan(\sqrt{r^2 - 1}) \right), \\ h_2(r, \theta_1, r_1) &= \frac{\sqrt{3}}{3} \sin \theta_1 + \frac{3}{2}(2\sqrt{3}\pi - 3)\sqrt{3} \cos \theta_1 - \frac{\sqrt{3}}{9}r_1, \\ h_3(r, \theta_1, r_1) &= \frac{9\sqrt{3} \sin \theta_1}{2r_1} - \frac{9\pi \sin \theta_1}{r_1} + \frac{\sqrt{3} \cos \theta_1}{3r_1} - \frac{3}{2} \left(\sqrt{3} + \frac{2\pi}{3} \right) + 4\pi. \end{aligned}$$

In order to solve the system $h_1 = h_2 = h_3 = 0$ we can use the same reasoning applied in the proof of Theorem 1 obtaining that $(r^*, \theta_1^*, r_1^*) = (2, \pi/2, 3)$ is a zero of the average function. Moreover if $J = J(r, \theta_1, r_1)$ is the Jacobian matrix of h , then $\det J(2, \pi/2, 3) \neq 0$ which implies that we have a simple zero. By Theorem 5 system (17) and consequently system (16) has one limit cycle for $|\varepsilon| > 0$ sufficiently small.

4. PROOF OF THEOREM 2

This section is devoted to prove Theorem 2. According with Theorem 7 the same kind of arguments used for proving Theorem 1 can be applied to the discontinuous system (6), obtaining that the average function $h = (h_1, h_2, \dots, h_{2(n-1)}, h_{2n-1}) : D_m \rightarrow \mathbb{R}^{n-1}$ defined in (12) is

$$\begin{aligned}
 h_1 &= (a_{11} + a_{22})\pi r + b_1 \tilde{I}_1, \\
 h_{2(j-1)} &= (a_{(2j-1)(2j-1)} + a_{(2j)(2j)})\pi r_{j-1} + (b_{2j-1} \cos \theta_{j-1} + b_{2j} \sin \theta_{j-1})\tilde{I}_j, \\
 h_{2j-1} &= \frac{(a_{(2j)(2j-1)} - a_{(2j-1)(2j)} + (2j-1)(a_{12} - a_{21}))\pi - (2j-1)b_2 r_{j-1} \tilde{I}_1 - r(b_{2j} \cos \theta_{j-1} - b_{2j-1} \sin \theta_{j-1})\tilde{I}_j}{rr_{j-1}},
 \end{aligned} \tag{18}$$

for $j = 2, 3, \dots, n$, where

$$\tilde{I}_j = \begin{cases} -\frac{4}{(2j-1)} & \text{if } j \text{ is even,} \\ \frac{4}{(2j-1)} & \text{if } j \text{ is odd.} \end{cases}$$

In fact if we define

$$\begin{aligned}
 \tilde{I}_j &= \int_0^{2\pi} \psi(r \cos \theta) \cos((2j-1)\theta) d\theta, \\
 \tilde{J}_j &= \int_0^{2\pi} \psi(r \cos \theta) \sin((2j-1)\theta) d\theta,
 \end{aligned}$$

where ψ is the piecewise linear function given by (3). Then we have that

$$\begin{aligned}
 \tilde{I}_j &= \int_0^{2\pi} \psi(r \cos \theta) \cos((2j-1)\theta) d\theta \\
 &= \int_0^{\pi/2} \cos((2j-1)\theta) d\theta - \int_{\pi/2}^{3\pi/2} \cos((2j-1)\theta) d\theta + \int_{3\pi/2}^{2\pi} \cos((2j-1)\theta) d\theta \\
 &= -\frac{4}{(2j-1)} \cos(j\pi),
 \end{aligned}$$

and

$$\begin{aligned}
\tilde{J}_j &= \int_0^{2\pi} \psi(r \cos \theta) \sin((2j-1)\theta) d\theta \\
&= \int_0^{\pi/2} \sin((2j-1)\theta) d\theta - \int_{\pi/2}^{3\pi/2} \sin((2j-1)\theta) d\theta + \int_{3\pi/2}^{2\pi} \sin((2j-1)\theta) d\theta \\
&= -\frac{4}{(2j-1)} \sin(2j\pi) \cos(j\pi) = 0.
\end{aligned}$$

Note that \tilde{I}_j is a constant real number different from zero, and hence h_1 is a straight line, and then system (18) has at most one positive zero. Moreover if we choose conveniently the coefficients b_1 , a_{11} and a_{22} we can find a simple positive zero of system (18). This completes the proof of Theorem 2.

5. PROOF OF THEOREMS 3 AND 4

Doing the change of coordinates

$$\begin{aligned}
x_1 &= r \cos \theta, & x_2 &= r \sin \theta, \\
x_{2j-1} &= r_{j-1} \cos(j\theta + \theta_{j-1}), & x_{2j} &= r_{j-1} \sin(j\theta + \theta_{j-1}) \quad j \in \{2, 3, \dots, n\},
\end{aligned}$$

for $j = 2, 3, \dots, n$, to the continuous piecewise linear differential system (7), and working as in the proof of Theorem 1 we obtain that the average function $h = (h_1, h_2, \dots, h_{2n-1})$ defined in (10) now is given by

$$\begin{aligned}
(19) \quad h_1 &= (a_{11} + a_{22})\pi r + b_1 I_1(r), \\
h_{2(j-1)} &= (a_{(2j-1)(2j-1)} + a_{(2j)(2j)})\pi r_{j-1} + (b_{2j-1} \cos \theta_{j-1} + b_{2j} \sin \theta_{j-1}) I_j(r), \\
h_{2j-1} &= (a_{(2j)(2j-1)} - a_{(2j-1)(2j)} + j(a_{12} - a_{21}))\pi - \\
&\quad \frac{j b_2 r_{j-1} I_1(r) - r(b_{2j} \cos \theta_{j-1} - b_{2j-1} \sin \theta_{j-1}) I_j(r)}{r r_{j-1}},
\end{aligned}$$

where

$$I_j(r) = \int_0^{2\pi} \varphi(r \cos \theta) \cos(j\theta) d\theta.$$

Using exactly the same arguments than in the proof of Lemma 9 is possible to prove that

$$I_j(r) = \begin{cases} \pi r & \text{if } j = 1 \text{ and } 0 < r \leq 1, \\ 0 & \text{if } j \text{ is even and } 0 < r \leq 1, \\ L_j(r) & \text{if } j \text{ is odd and } r > 1, \end{cases}$$

where

$$L_j(r) = \frac{4}{j(j^2 - 1)} \left(j \sqrt{r^2 - 1} \cos(j \arctan(\sqrt{r^2 - 1})) - \sin(j \arctan(\sqrt{r^2 - 1})) \right).$$

The simple zeros of system (19) provide the existence of limit cycles for system (7) but since $I_j(r) = 0$ if j is even and $r > 1$, the variables θ_{j-1} , for $j = 2, 4, 6, \dots$ do not appear in the system $h_1 = h_2 = \dots = h_{2n-1} = 0$, so either this system has no zeros, or if it has zeros, then it has a continuum of zeros, and therefore the assumption (ii) of the averaging Theorem 9 does not hold, and consequently the averaging theory cannot say anything about the limit cycles of system (7). The same occurs for the case $0 < r \leq 1$. So we conclude that, using the averaging theory of first order, we can say nothing about the number of the limit cycles of system (7). This completes the proof of Theorems 3.

Now if we consider the discontinuous piecewise linear differential system (8), then its average function $h = (h_1, h_2, \dots, h_{2n-1})$ defined in (12) is

$$\begin{aligned}
 h_1 &= (a_{11} + a_{22})\pi r + b_1 \tilde{I}_1, \\
 h_{2(j-1)} &= (a_{(2j-1)(2j-1)} + a_{(2j)(2j)})\pi r_{j-1} + (b_{2j-1} \cos \theta_{j-1} + b_{2j} \sin \theta_{j-1})\tilde{I}_j, \\
 h_{2j-1} &= \frac{(a_{(2j)(2j-1)} - a_{(2j-1)(2j)} + j(a_{12} - a_{21}))\pi - j b_{2j} r_{j-1} \tilde{I}_1 - r(b_{2j} \cos \theta_{j-1} - b_{2j-1} \sin \theta_{j-1})\tilde{I}_j}{r r_{j-1}},
 \end{aligned}
 \tag{20}$$

where

$$\tilde{I}_j = \int_0^{2\pi} \psi(r \cos \theta) \cos(j\theta) d\theta.$$

Again we have that

$$\tilde{I}_j = \int_0^{2\pi} \psi(r \cos \theta) \cos(j\theta) d\theta = \begin{cases} 0 & \text{if } j \text{ is even,} \\ \pm \frac{4}{(2j-1)} & \text{if } j \text{ is odd,} \end{cases}$$

and we can see that again either no zeros of the function h , or a continuum of zeros, concluding that the averaging theory of first order given by Theorem 7 does not say anything about the limit cycles of system (8). This completes the proof of Theorems 4.

ACKNOWLEDGEMENTS

The first author is partially supported by a MINECO grant number MTM2013-40998-P, an AGAUR grant number 2014SGR568 and a FP7-PEOPLE-2012-IRSES grant number 318999. The second author is supported by a Projeto Temático FAPESP number 2014/00304-2. The first and the second authors are also supported by a FP7-PEOPLE-2012-IRSES grant number 316338 and a CAPES CSF-PVE grant 88881.030454/2013-01 from the program CSF-PVE. The third author has a PhD fellowship from CAPES-Brazil.

REFERENCES

- [1] M.A. AIZERMAN, *Theory of automatic control*, Pergamon Press, 1963.
- [2] A. ANDRONOV, A. VITT AND S. KHAIKIN, *Theory of Oscillations*, Pergamon Press, Oxford, 1966.
- [3] S. BARNETT AND R.G. CAMERON, *Introduction to mathematical control theory*, 2nd edition, Clarendon Press, Oxford, 1985.
- [4] M. DI BERNARDO, C.J. BUDD, A.R. CHAMPNEYS AND P. KOWALCZYK, *Piecewise-smooth Dynamical Systems*, Springer-Verlag London, Applied Mathematical Sciences, vol. **163**, 2008.
- [5] A. BUICA AND J. LLIBRE, *Averaging methods for finding periodic orbits via Brouwer degree*, Bull. Sci. Math. **128** (2004), 7–22.
- [6] A. BUICA AND J. LLIBRE, *Bifurcation of limit cycles from a four-dimensional center in control systems*, Int. J. of Bifurcation and Chaos, **15** (2005), 2653–2662.
- [7] C.A. BUZZI, J. LLIBRE, J.C. MEDRADO AND J. TORREGROSA, *Bifurcation of limit cycles from a center in \mathbb{R}^4 in resonance $1 : N$* , Dynamical Systems **24** (2009), 123–137.
- [8] P.T. CARDIN, T. DE CARVALHO AND J. LLIBRE, *Bifurcation of limit cycle from a n -dimensional linear center inside a class of piecewise linear differential systems*, Nonlinear Analysis: Theory, Methods and Appl. **75** (2012), 143–152.
- [9] C. CHRISTOPHER AND C. LI, *Limit cycles in differential equations*, Birkhauser, Boston, 2007.
- [10] J.E. GIBSON, *Nonlinear automatic control*, McGraw-Hill, New York, 1963.
- [11] H.K. KHALIL, *Nonlinear systems*, Macmillan, New York, 1992.
- [12] J. LLIBRE AND A. MAKHLOUF, *Bifurcation of limit cycles from a 4-dimensional center in $1 : n$ resonance*, Applied Mathematics and Computation **215** (2009), 140–149.
- [13] J. LLIBRE, D.D. NOVAES AND M.A. TEIXEIRA, *On the birth of limit cycles for non-smooth dynamical systems*, Bull. Sci. Math. **139** (2015), 229–244.
- [14] N.G. LLOYD, *Degree theory*, Cambridge Tracts in Math. vol. **73**, Cambridge Univ. Press, Cambridge, Great Britain, 1978.
- [15] A.I. MEES, *Dynamics of feedback systems*, John Wiley & Sons, 1981.
- [16] F. VERHULST, *Nonlinear Differential Equations and Dynamical Systems*, 2nd edition, Universitext, Springer 1996.
- [17] M. VIDYASAGAR, *Nonlinear systems analysis*, 2nd ed., Prentice-Hall, 1993.

¹ DEPARTAMENT DE MATEMÀTIQUES, UNIVERSITAT AUTÒNOMA DE BARCELONA, 08193 BELLATERRA, BARCELONA, CATALONIA, SPAIN

E-mail address: jllibre@mat.uab.cat

^{2,3} DEPARTMENT OF MATHEMATIC, ICMC-UNIVERSIDADE DE SÃO PAULO, C.P. 668, 13.560-970, SÃO CARLOS, SP, BRAZIL

E-mail address: regilene@icmc.usp.br, camilaap@icmc.usp.br