# INTEGRABILITY AND ZERO-HOPF BIFURCATION IN THE SPROTT A SYSTEM 

LUIS BARREIRA ${ }^{1}$, JAUME LLIBRE ${ }^{2}$ AND CLAUDIA VALLS ${ }^{1}$


#### Abstract

The first objective of this paper is to study the Darboux integrability of the polynomial differential system $$
\dot{x}=y, \quad \dot{y}=-x-y z, \quad \dot{z}=y^{2}-a
$$ and the second one is to show that for $a>0$ sufficiently small this model exhibits one small amplitude periodic solution that bifurcates from the origin of coordinates when $a=0$. This model was introduced by Hoover as the first example of a differential equation with a hidden attractor and it was used by Sprott to illustrate a differential equation having a chaotic behavior without equilibrium points, and now this system is known as the Sprott A system.


## 1. Introduction and statement of the main results

In differential equations it is very important to know whether a differential system presents chaos. One might think that it is not possible to have chaotic behavior in a system without equilibrium points, but now we know that this is not the case. In this paper we study one of these differential systems. More precisely we consider the well-known example of differential equation with a hidden chaotic attractor, which is a special case of the Nosé-Hoover oscillator [10] and as showed in [18] describes many natural phenomena. The Nosé-Hoover thermostat is a differential method used in molecular dynamics to keep the temperature around an average. It was first introduced by Nosé in 1985 [17] and developed further by Hoover [9]. In this paper we consider the Nosé-Hoover equation in the form

$$
\begin{equation*}
\dot{x}=y, \quad \dot{y}=-x-y z, \quad \dot{z}=y^{2}-a, \tag{1}
\end{equation*}
$$

where $a \in \mathbb{R}$ is a parameter and the dot denotes derivative with respect to the time $t$. In [21] Sprott listed this system as one of the differential systems without equilibrium points having chaotic behavior (when $a=1$ ), and from then system (1) is usually called Sprott A system, see for instance [16]. When $a=0$ the system has a continuum of equilibria, i.e. $(0,0, z)$ for

[^0]all $z \in \mathbb{R}$. For $a \neq 0$ it has no equilibria and for $a=1$ it generates a chaotic attractor, see for more details [21].

A number of dynamical aspects of system (1) have been analysed such as the formation of hidden chaotic attractors as well as the formation of nested invariant tori, and the dynamics at infinity (see [16] and the references therein). Here our first main objective is to further contribute to the understanding the complexity of the system, or more precisely the dynamics of system (1) by studying its integrability. We note that the existence of a first integral allows to reduce its study in one dimension. Moreover, inside the class of first integrals, the simpler ones are the so called Darboux first integrals, for more information about these kind of first integrals see for instance $[5,11,14,15]$ and subsection 2.1.

The second objective of this paper is to study the zero-Hopf bifurcation which exhibits the polynomial differential system (1) when $a=0$. The main tool up to now for studying a zero-Hopf bifurcation is to pass the system to the normal form of a zero-Hopf bifurcation, later on this introduction we provide the references about this. Our analysis of the zero-Hopf bifurcation is different, we study them directly using the averaging theory, see subsection 2.2 for a summary of the results of this theory that we need in this paper. We note that this paper is one of the few papers where we study a zero-Hopf bifurcation that is not isolated.

Let $U$ be an open and dense subset of $\mathbb{R}^{3}$, we say that a non-locally constant $C^{1}$ function $H: U \rightarrow \mathbb{R}$ is a first integral of system (1) on $U$ if $H(x(t), y(t), z(t))$ is constant for all of the values of $t$ such that for which $(x(t), y(t), z(t))$ is a solution of system (1) contained in $U$.

The first main result of this paper is the following. See section 2 for the definition of Darboux polynomial and Darboux first integral.

Theorem 1. The next statements hold for system (1):
(i) If $a \neq 0$ it has no Darboux polynomials.
(ii) If $a=0$ the unique Darboux polynomials have cofactor zero and are polynomials in the variable $H=x^{2}+y^{2}+z^{2}$.
(iii) If $a \neq 0$ it has no Darboux first integrals.
(iv) If $a=0$ the unique Darboux first integrals are Darboux functions in the variable $H$.

The second main objective of this paper is to show that system (1) exhibits one small amplitude periodic solution for $a>0$ sufficiently small that bifurcates from the non-isolated zero-Hopf equilibrium $(0,0,0)$ localized at the $z$-axis filled of equilibria when $a=0$.

We recall that a zero-Hopf equilibrium of a 3-dimensional autonomous differential system is an equilibrium such that it has a zero real eigenvalue and a pair of purely imaginary eigenvalues. We know that generically an
isolated zero-Hopf equilibrium exhibits a zero-Hopf bifurcation, which is a two-parameter unfolding of a 3-dimensional autonomous differential system with a zero-Hopf equilibria. The unfolding can exhibit different topological type of dynamics in a small neighborhood of this non-isolated equilibria and in some cases the zero-Hopf bifurcation can imply a local birth of chaos (see for instance $[3,7,8,12,20]$ ). This previous unfolding of a zero-Hopf bifurcation cannot be applied to our differential system (1) because first our system has only one parameter and second because the zero-Hopf equilibrium of system (1) is not isolated. For studying the zero-Hopf bifurcation of system (1), we shall use the averaging theory in a similar way at it was used in [2].

The second main result of this paper characterizes the zero-Hopf bifurcation of system (1).

Theorem 2. The differential system (1) has a zero-Hopf bifurcation at the non isolated equilibrium point $(0,0,0)$ when $a=0$ producing one small periodic solution for $a>0$ sufficiently small with shape

$$
x(t)=\sqrt{2 a} \cos t+O(a), \quad y(t)=-\sqrt{2 a} \sin t+O(a), \quad z(t)=O(a) .
$$

From the numerical work [16] on the differential system (1) we see that the periodic orbit found in Theorem 2 is surrounded by invariant tori.

The paper is organized as follows. In section 2 we present the basic definitions and necessary results to prove Theorems 1 and 2. In section 3 we prove Theorem 1 and in section 4 we prove Theorem 2.

## 2. Preliminaries

2.1. Darboux theory of integrability. We denote by $\mathbb{C}[x, y, z]$ the ring of polynomials in the variables $x, y, z$ and coefficients in $\mathbb{C}$. Given $f \in$ $\mathbb{C}[x, y, z] \backslash \mathbb{C}$ we say that the polynomial $f(x, y, z)$ is a Darboux polynomial of system (1) if there exists $K \in \mathbb{C}[x, y, z]$ called the cofactor such that

$$
\begin{equation*}
y \frac{\partial f}{\partial x}-(x+y z) \frac{\partial f}{\partial y}+\left(y^{2}-a\right) \frac{\partial f}{\partial z}=K f . \tag{2}
\end{equation*}
$$

The degree of $K$ is at most 1 . If $K=0$ then $f$ is a polynomial first integral. When a real polynomial system has a complex Darboux polynomial then it has also its conjugate. It is important to consider the complex Darboux polynoials of the real polynomial differential systems because sometimes they force the real integrability of the system. Note that we can write $f$ as

$$
\begin{equation*}
X f=K f, \quad X=y \frac{\partial}{\partial x}-(x+y z) \frac{\partial}{\partial y}+\left(y^{2}-a\right) \frac{\partial}{\partial z} . \tag{3}
\end{equation*}
$$

Let $f, g \in \mathbb{C}[x, y, z]$ and assume that $f$ and $g$ are relatively prime in the ring $\mathbb{C}[x, y, z]$, or that $g=1$. Then the function $\exp (f / g)$ is called an
exponential factor of system (1) if for some polynomial $L \in \mathbb{C}[x, y, z]$ of degree one, we have

$$
y \frac{\partial \exp (f / g)}{\partial x}-(x+y z) \frac{\partial \exp (f / g)}{\partial y}+\left(y^{2}-a\right) \frac{\partial \exp (f / g)}{\partial z}=L \exp (f / g)
$$

As before, we call $L$ the cofactor of $\exp (f / g)$. We note that in the definition of exponential factor if $f, g \in \mathbb{C}[x, y, z]$ then the exponential factor is a complex function. Again, when a real polynomial differential system has a complex exponential factor, then it has also its conjugate, and both are important for the existence of real first integrals of the system. The exponential factors are related with the multiplicity of the Darboux polynomials, see for more details [4], Chapter 8 of [6] and [14, 15].

A first integral is called a Darboux first integral if it is a first integral of the form

$$
\begin{equation*}
f_{1}^{\lambda_{1}} \cdots f_{p}^{\lambda_{p}} F_{1}^{\mu_{1}} \cdots F_{q}^{\mu_{q}} \tag{4}
\end{equation*}
$$

where $f_{i}$ are Darboux polynomials of system (1) for $i=1, \ldots, p$ and $F_{j}$ are exponential factors of system (1) for $j=1, \ldots, q$ and $\lambda_{i}, \mu_{j} \in \mathbb{C}$. The next result proved in [6] indicates how to find Darboux first integrals.

Proposition 3. Assume that a polynomial differential system of degree $m$ admits $p$ Darboux polynomials $f_{i}$ with cofactors $k_{i}$ for $i=1, \ldots, p$ and $q$ exponential factors $F_{j}=\exp \left(g_{j} / h_{j}\right)$ with cofactors $L_{j}$ for $j=1, \ldots, q$. Then there exist $\lambda_{i}$ and $\mu_{j} \in \mathbb{C}$ not all zero such that

$$
\sum_{i=1}^{p} \lambda_{i} K_{i}+\sum_{j=1}^{q} \mu_{j} L_{j}=0
$$

if and only if the function (4) is a Darboux first integral of that polynomial system.

The following result is proved in [6].
Proposition 4. Let $f$ be a Darboux polynomial and $f=\prod_{j=1}^{s} f_{j}^{\alpha_{j}}$ its decomposition into irreducible factors in $\mathbb{C}[x, y, z]$. Then $f$ us a Darboux polynomial if and only if all $f_{j}$ are Darboux polynomials. Moreover, if $K$ and $K_{j}$ are the cofactors of $f$ and $f_{j}$, then $K=\sum_{j=1}^{s} \alpha_{j} K_{j}$.

The following proposition is proved in [14, 15].
Proposition 5. Let $f, g \in \mathbb{C}[x, y, z]$. The following statements hold.
(1) If $\exp (f / g)$ is an exponential factor for the polynomial differential system (1) and $g$ is not a constant function, then $g$ is a Darboux polynomial of (1).
(2) The existence of an exponential factor of the form $\exp (f)$ is due to the multiplicity of the invariant plane at infinity.

The following proposition is well known and its proof is similar to the proof of Proposition 4.

Proposition 6. Let $F_{j}=\exp \left(f_{j} / g_{j}\right)$ for $j=1, \ldots, s$ be exponential factors of system (1) with cofactors $L_{j}$. Then

$$
F=\prod_{j=1}^{s} F_{j}=\exp \left(\sum_{j=1}^{s} \frac{f_{j}}{g_{j}}\right)
$$

is an exponential factor of system (1) with cofactor $L=\sum_{j=1}^{s} L_{j}$.
2.2. Averaging theory. We present a result from the averaging theory that we shall need for proving Theorem 2, for a general introduction to the averaging theory of integrability see $[13,19]$.

We consider the initial value problems

$$
\begin{equation*}
\dot{x}=\varepsilon F_{1}(t, x)+\varepsilon^{2} F_{2}(t, x, \varepsilon), \quad x(0)=x_{0} \tag{5}
\end{equation*}
$$

and

$$
\begin{equation*}
\dot{y}=\varepsilon g(y), \quad y(0)=x_{0} \tag{6}
\end{equation*}
$$

with $x, y$ and $x_{0}$ in some open subset $\Omega$ of $\mathbb{R}^{n}, t \in[0, \infty), \varepsilon \in\left(0, \varepsilon_{0}\right]$. We assume that $F_{1}$ and $F_{2}$ are periodic of period $T$ in the variable $t$, and we set

$$
\begin{equation*}
g(y)=\frac{1}{T} \int_{0}^{T} F_{1}(t, y) d t \tag{7}
\end{equation*}
$$

We will also use the notation $D_{x} g$ for all the first derivatives of $g$, and $D_{x x} g$ for all the second derivatives of $g$. For a proof of the next result, see [22].

Theorem 7. Assume that $F_{1}, D_{x} F_{1}, D_{x x} F_{1}$ and $D_{x} F_{2}$ are continuous and bounded by a constant independent of $\varepsilon$ in $[0, \infty) \times \Omega \times\left(0, \varepsilon_{0}\right]$ and that $y(t) \in \Omega$ for $t \in[0,1 / \varepsilon]$. Then the following statements hold:
(1) For $t \in[0,1 / \varepsilon]$, we have $x(t)-y(t)=O(\varepsilon)$ as $\varepsilon \rightarrow 0$;
(2) If $p \neq 0$ is a singular point of system (6) such that

$$
\begin{equation*}
\operatorname{det} D_{y} g(p) \neq 0 \tag{8}
\end{equation*}
$$

then there exists a periodic solution $x(t, \varepsilon)$ of period $T$ os system (5) which is close to $p$ and such that $x(0, \varepsilon)-p=O(\varepsilon)$ as $\varepsilon \rightarrow 0$;
(3) The stability of the periodic solution $x(t, \varepsilon)$ is given by the stability of the singular point $p$.

## 3. Proof of Theorem 1

To prove Theorem 1 we state and prove several auxiliary results. Let $\mathbb{N}$ be the set of positive integers and $\tau: \mathbb{C}[x, y, z] \rightarrow \mathbb{C}[x, y, z]$ be the automorphism

$$
\tau(x)=-x, \quad \tau(y)=-y, \quad \tau(z)=z
$$

Lemma 8. If $f$ is a Darboux polynomial for system (1) with cofactor $K$, then $\tau(f)$ is a Darboux polynomial of system (1) with cofactor $\tau(K)$.

Proof. Since $f$ is a Darboux polynomial for system (1) with cofactor $K$ in view of (3) we can write it as $X f=K f$. Since (1) is invariant by $\tau$ we have that $\tau_{*}(X)=X$, where $\tau_{*}$ is the push-forward associated to automorphism $\tau$. Therefore

$$
\tau(X f)(x, y, z)=\tau(K f)(x, y, z)
$$

The left-hand side is

$$
\begin{aligned}
\tau(X f)(x, y, z) & =\tau_{*}(X) \tau(f)(x, y, z)=X f\left(\tau^{-1}(x, y, z)\right)=X f(\tau(x, y, z)) \\
& =X f(-x,-y, z)=X \tau(f)(x, y, z)
\end{aligned}
$$

where we have used that $\tau^{-1}=\tau$. The right-hand side is

$$
\begin{aligned}
\tau(K f)(x, y, z) & =\left((K f) \circ \tau^{-1}\right)(x, y, z)=((K f) \circ \tau)(x, y, z)=(K f)(-x,-y, z) \\
& =K(-x,-y, z) f(-x,-y, z)=\tau(K) \tau(f)(x, y, z)
\end{aligned}
$$

From these two previous equalities we get

$$
X \tau(f)=\tau(K) \tau(f)
$$

and the lemma follows.
Lemma 9. If $f$ is an irreducible Darboux polynomial of degree $n \geq 1$ for system (1) with cofactor $K=k_{0}+k_{1} x+k_{2} y+k_{3} z$ with $k_{i} \in \mathbb{C}$ for $i=0,1,2,3$, then $g=f \cdot \tau(f)$ is a Darboux polynomial invariant by $\tau$ of degree $2 n$ with cofactor $2\left(k_{0}+k_{3} z\right)$. If the cofactor of $f$ is zero, then the cofactor of $g$ is also zero.

Proof. It follows from Lemma 8 that $\tau(f)$ is a Darboux polynomial of system (1) with cofactor $\tau(K)$. Therefore, by Proposition 4 the cofactor of $g=$ $f \cdot \tau(f)$ is $K+\tau(K)=2\left(k_{0}+k_{3} z\right)$. This concludes the proof of the lemma.
Lemma 10. Let $g$ be a Darboux polynomial of system (1) of degree $n$. Then the polynomial $f=g \cdot \tau(g)$ is a Darboux polynomial invariant by $\tau$ with degree $2 n$ and with cofactor $K=2 k_{0}-2 m z$ for $k_{0} \in \mathbb{C}$ and $m \in \mathbb{N} \cup\{0\}$. Moreover,

$$
\begin{equation*}
f_{2 n}(x, y, z)=y^{2 m} \sum_{l=0}^{n-m} a_{l}\left(x^{2}\right)^{n-m-l}\left(y^{2}+z^{2}\right)^{l}, \quad a_{l} \in \mathbb{C} . \tag{9}
\end{equation*}
$$

Proof. We write $f$ in sum of its homogeneous parts as

$$
f=\sum_{i=0}^{2 n} f_{i}, \quad f_{i} \text { is a homogeneous polynomial of degree } i
$$

Without loss of generality we can assume that $f_{2 n} \neq 0$ and $n \geq 1$. Computing the terms of degree $2 n+1$ in (2) we get

$$
-y z \frac{\partial f_{2 n}}{\partial y}+y^{2} \frac{\partial f_{2 n}}{\partial z}=2 k_{3} z f_{2 n}
$$

Solving this linear partial differential equation we get

$$
f_{2 n}=C_{n}\left(x, y^{2}+z^{2}\right) y^{-2 k_{3}}
$$

where $C_{n}$ is a polynomial in the variables $x, y^{2}+z^{2}$. Since $f_{2 n}$ is a polynomial we must have $k_{3}=-m$ with $m \in \mathbb{N} \cup\{0\}$. Moreover, since $f_{2 n}$ is a polynomial of degree $2 n$ it can be written as in (9) (note that $f_{2 n}$ is invariant by $\tau)$. This concludes the proof of the lemma.

Lemma 11. Let $g$ be a Darboux polynomial of system (1) of degree $n$. Then the polynomial $f=g \cdot \tau(g)$ is a Darboux polynomial invariant by $\tau$ with degree $2 n$ and with cofactor $K=-2 m z$ for $m \in \mathbb{N} \cup\{0\}, f_{2 n}=$ $\alpha_{0} y^{2 m}\left(x^{2}+y^{2}+z^{2}\right)^{n-m}$ for some $\alpha_{0} \in \mathbb{C} \backslash\{0\}$ and

$$
f_{2 n-1}=2 \alpha_{0} m z y^{2 m-1} \sum_{l=0}^{n-m}\binom{n-m}{l} x^{2 n-2 m-2 l+1}\left(y^{2}+z^{2}\right)^{l-1}
$$

Proof. We compute the terms of degree $2 n-1$ using $f_{2 n}$ given in (9). Taking the terms of degree $2 n$ from (2) we get

$$
\begin{equation*}
\mathcal{L}\left[f_{2 n-1}\right]=-2 m z f_{2 n-1}-y \frac{\partial f_{2 n}}{\partial x}+x \frac{\partial f_{2 n}}{\partial y}+2 k_{0} f_{2 n} \tag{10}
\end{equation*}
$$

where

$$
\begin{equation*}
\mathcal{L}=-y z \frac{\partial}{\partial y}+y^{2} \frac{\partial}{\partial z} \tag{11}
\end{equation*}
$$

Note that in view of the definition of $f_{2 n}$ we have

$$
\begin{aligned}
R(x, y, z)= & -y \frac{\partial f_{2 n}}{\partial x}+x \frac{\partial f_{2 n}}{\partial y}+2 k_{0} f_{2 n} \\
= & -2 x y^{2 m+1} \sum_{l=0}^{n-m-1} a_{l}(n-m-l)\left(x^{2}\right)^{n-m-l-1}\left(y^{2}+z^{2}\right)^{l} \\
& +2 m x y^{2 m-1} \sum_{l=0}^{n-m} a_{l}\left(x^{2}\right)^{n-m-l}\left(y^{2}+z^{2}\right)^{l} \\
& +2 x y^{2 m+1} \sum_{l=1}^{n-m} a_{l} l\left(x^{2}\right)^{n-m-l}\left(y^{2}+z^{2}\right)^{l-1} \\
& +2 k_{0} y^{2 m} \sum_{l=0}^{n-m} a_{l}\left(x^{2}\right)^{n-m-l}\left(y^{2}+z^{2}\right)^{l}
\end{aligned}
$$

$$
\begin{aligned}
= & 2 x y^{2 m+1} \sum_{l=0}^{n-m-1} b_{l}\left(x^{2}\right)^{n-m-l-1}\left(y^{2}+z^{2}\right)^{l} \\
& +2 m x y^{2 m-1} \sum_{l=0}^{n-m} a_{l}\left(x^{2}\right)^{n-m-l}\left(y^{2}+z^{2}\right)^{l} \\
& +2 k_{0} y^{2 m} \sum_{l=0}^{n-m} a_{l}\left(x^{2}\right)^{n-m-l}\left(y^{2}+z^{2}\right)^{l}
\end{aligned}
$$

where

$$
b_{l}=-(n-m-l) a_{l}+(l+1) a_{l+1} .
$$

We apply now the method of characteristic curves to solve (10). For the method of characteristic curves for solving linear partial differential equations we refer the reader, for instance, to Bleecker and Csordas [1].

The characteristic equations associated with the first linear partial differential equation of system

$$
\dot{y}=-y z, \quad \dot{z}=y^{2}
$$

are of the form

$$
\frac{d y}{d z}=-\frac{z}{y}
$$

This system has the general solution $y^{2}+z^{2}=\kappa$, where $\kappa$ is a constant.
According with the method of characteristics we make the change of variables

$$
\begin{equation*}
u=x, \quad v=z, \quad w=y^{2}+z^{2} \tag{12}
\end{equation*}
$$

Its inverse transformation is

$$
\begin{equation*}
x=u, \quad y= \pm \sqrt{w-v^{2}}, \quad z=v \tag{13}
\end{equation*}
$$

In the following for simplicity we only consider the case $y=+\sqrt{w-z^{2}}$. In what follows we always write $\bar{\theta}$ to denote a function $\theta=\theta(x, y, z)$ written in the $(u, v, w)$ variables, that is, $\bar{\theta}=\bar{\theta}(u, v, w)$. Under changes (12) and (13), equation in (10) becomes the following ordinary differential equation (for fixed $u, w$ )

$$
\begin{equation*}
\left(w-v^{2}\right) \frac{d \bar{f}_{2 n-1}}{d v}=-2 m v \bar{f}_{2 n-1}+\bar{R}(u, w, v) \tag{14}
\end{equation*}
$$

where

$$
\begin{aligned}
& \bar{R}(u, w, v)= 2 u \sqrt{w-v^{2}} \\
& 2 m+1 \\
& \sum_{l=0}^{n-m-1} b_{l}\left(u^{2}\right)^{n-m-l-1} w^{l} \\
&+2 m u \sqrt{w-v^{2}} \\
& 2 m-1 \\
& \sum_{l=0}^{n-m} a_{l}\left(u^{2}\right)^{n-m-l} w^{l} \\
&+2 k_{0} \sqrt{w-v^{2}}{ }^{2 m} \sum_{l=0}^{n-m} a_{l}\left(u^{2}\right)^{n-m-l} w^{l} .
\end{aligned}
$$

Equation (14) has the general solution

$$
\begin{aligned}
\bar{f}_{2 n-1}= & \bar{F}_{2 n-1}(u, w){\sqrt{w-v^{2}}}^{2 m}+{\sqrt{w-v^{2}}}^{2 m} \int \frac{\bar{R}(u, w, v)}{{\sqrt{w-v^{2}}}^{2 m+2}} d v \\
= & \bar{F}_{n-1}(u, w){\sqrt{w-v^{2}}}^{2 m} \\
& +2{\sqrt{w-v^{2}}}^{2 m} u \sum_{l=0}^{n-m-1} b_{l}\left(u^{2}\right)^{n-m_{2}-l-1} w^{l} \int \frac{1}{\sqrt{w-v^{2}}} d v \\
& +2 m{\sqrt{w-v^{2}}}^{2 m} u \sum_{l=0}^{n-m} a_{l}\left(u^{2}\right)^{n-m-l} w^{l} \int \frac{1}{\sqrt{w-v^{2}}} d v \\
& +2 k_{0}{\sqrt{w-v^{2}}}^{2 m} \sum_{l=0}^{n-m} a_{l}\left(u^{2}\right)^{n-m-l} w^{l} \int \frac{1}{w-v^{2}} d v,
\end{aligned}
$$

where $F_{2 n-1}$ is a smooth function in the variables $u, w$. Using that

$$
\begin{aligned}
\int \frac{1}{\sqrt{w-v^{2}}} d v & =\arctan \left(\frac{v}{\sqrt{w-v^{2}}}\right) \\
\int \frac{1}{\left(w-v^{2}\right)^{3 / 2}} d v & =\frac{v}{w \sqrt{w-v^{2}}} \\
\int \frac{1}{w-v^{2}} d v & =\frac{1}{\sqrt{w}} \arctan \left(\frac{v}{\sqrt{w}}\right),
\end{aligned}
$$

and that $f_{2 n-1}$ must be a homogeneous polynomial of degree $2 n-1$ we must have

$$
k_{0}=0 \quad \text { and } \quad b_{l}=0, \quad l=0, \ldots, n-m-1
$$

Therefore, we have

$$
a_{l}=\frac{n-m-l+1}{l} a_{l-1} \quad \text { that is } \quad a_{l}=\binom{n-m}{l} \alpha_{0}, \quad \alpha_{0} \in \mathbb{C} .
$$

This implies that

$$
\begin{align*}
f_{2 n} & =y^{2 m} \sum_{l=0}^{n-m}\binom{n-m}{l} \alpha_{0}\left(x^{2}\right)^{n-m-l}\left(y^{2}+z^{2}\right)^{l}  \tag{15}\\
& =\alpha_{0} y^{2 m}\left(x^{2}+y^{2}+z^{2}\right)^{n-m}
\end{align*}
$$

Note that

$$
f_{2 n-1}=y^{2 m} F_{n-1}\left(x, y^{2}+z^{2}\right)+2 m z y^{2 m-1} \sum_{l=0}^{n-m} a_{l} x^{2 n-2 m-2 l+1}\left(y^{2}+z^{2}\right)^{l-1}
$$

Since $f_{2 n-1}$ is a homogeneous polynomial of degree $2 n-1$ and must be invariant by $\tau$ we must have that $F_{2 n-1}$ is even in the variable $x$, and so $F_{2 n-1}$ must be zero. Then

$$
\begin{align*}
f_{2 n-1} & =2 m z y^{2 m-1} \sum_{l=0}^{n-m} a_{l} x^{2 n-2 m-2 l+1}\left(y^{2}+z^{2}\right)^{l-1}  \tag{16}\\
& =2 \alpha_{0} m z y^{2 m-1} \sum_{l=0}^{n-m}\binom{n-m}{l} x^{2 n-2 m-2 l+1}\left(y^{2}+z^{2}\right)^{l-1} .
\end{align*}
$$

This completes the proof of the lemma.
Theorem 12. The following holds for system (1).
(i) If $a \neq 0$ it has no Darboux polynomials invariant by $\tau$.
(ii) If $a=0$ the unique Darboux polynomials invariant by $\tau$ have cofactor zero and are polynomials in the variable $H=x^{2}+y^{2}+z^{2}$.

Proof. Assume that system (1) has a Darboux polynomial under the assumptions of Lemma 10. Using the notation and the computations in the proof of Lemma 10 we shall compute the terms of degree $2 n-2$ of $f$ using the terms of degree $2 n-1$ of (2). Thus we have

$$
\mathcal{L}\left[f_{2 n-2}\right]=-2 m z f_{2 n-2}-y \frac{\partial f_{2 n-1}}{\partial x}+x \frac{\partial f_{2 n-1}}{\partial y}+a \frac{\partial f_{2 n}}{\partial z}
$$

where $\mathcal{L}$ is given in (11). So from (15) and (16) we have

$$
\begin{aligned}
S(x, y, z)= & -y \frac{\partial f_{2 n-1}}{\partial x}+x \frac{\partial f_{2 n-1}}{\partial y}+a \frac{\partial f_{2 n}}{\partial z} \\
= & -2 \alpha_{0} m z y^{2 m} \sum_{l=0}^{n-m}\binom{n-m}{l}(2 n-2 m-2 l+1) x^{2 n-2 m-2 l}\left(y^{2}+z^{2}\right)^{l-1} \\
& +2 \alpha_{0} m z(2 m-1) y^{2 m-2} \sum_{l=0}^{n-m}\binom{n-m}{l} x^{2 n-2 m-2 l+2}\left(y^{2}+z^{2}\right)^{l-1} \\
& +4 \alpha_{0} m z y^{2 m} \sum_{l=0}^{n-m}\binom{n-m}{l} x^{2 n-2 m-2 l+2}(l-1)\left(y^{2}+z^{2}\right)^{l-2} \\
& +2 a \alpha_{0}(n-m) z y^{2 m}\left(x^{2}+y^{2}+z^{2}\right)^{n-m-1} .
\end{aligned}
$$

Changing the variables from $(x, y, z)$ to $(u, v, w)$ we obtain

$$
\left(w-v^{2}\right) \frac{d \bar{f}_{2 n-2}}{d v}=-2 m v \bar{f}_{2 n-2}+\bar{S}(u, v, w)
$$

Solving this linear differential equation we get

$$
\begin{aligned}
\bar{f}_{2 n-2}= & \bar{F}_{2 n-2}(u, w){\sqrt{w-v^{2}}}^{2 m} \\
& -2 \alpha_{0} m{\sqrt{w-v^{2}}}^{2 m} \sum_{l=0}^{n-m}\binom{n-m}{l}(2 n-2 m-2 l+1) u^{2 n-2 m-2 l} w^{l-1} \int \frac{v}{w-v^{2}} d v \\
& +2 \alpha_{0} m(2 m-1){\sqrt{w-v^{2}}}^{2 m} \sum_{l=0}^{n-m}\binom{n-m}{l} u^{2 n-2 m-2 l+2} w^{l-1} \int \frac{v}{\left(w-v^{2}\right)^{2}} d v \\
& +4 \alpha_{0} m{\sqrt{w-v^{2}}}^{2 m} \sum_{l=0}^{n-m}\binom{n-m}{l} u^{2 n-2 m-2 l+2}(l-1) w^{l-2} \int \frac{v}{w-v^{2}} d v \\
& +2 a \alpha_{0}(n-m) \sqrt{w-v^{2}}
\end{aligned}
$$

where $F_{2 n-2}$ is a smooth function in the variables $u, w$. Using that

$$
\begin{aligned}
\int \frac{v}{w-v^{2}} d v & =-\frac{1}{2} \log \left(w-v^{2}\right) \\
\int \frac{v}{\left(w-v^{2}\right)^{2}} d v & =\frac{1}{2\left(w-v^{2}\right)}
\end{aligned}
$$

From the fact that $f_{2 n-2}$ must be a polynomial we conclude that

$$
\alpha_{0} m=0 \quad \text { and } \quad a \alpha_{0}(n-m)=0
$$

Taking into account that $\alpha_{0} n \neq 0$ we must have $m=0$ and $a=0$.
Thus if $a \neq 0$ there are no Darboux polynomials invariant by $\tau$ for system (1). This proves statement (i) of Theorem 12.

If $a=0$ then $m_{2}=0, f_{2 n}=\alpha_{0}\left(x^{2}+y^{2}+z^{2}\right)^{n}, f_{2 n-1}=0$ and $f_{2 n-2}$ as in (9) with $n$ replaced by $n-1$. Proceeding as we did for $f_{2 n}$ we get that

$$
f_{2 n-2 k}=\alpha_{k}\left(x^{2}+y^{2}+z^{2}\right)^{n-k}, \quad \text { for } k=0, \ldots, n
$$

and

$$
f_{2 n-2 k-1}=0, \quad \text { for } k=0, \ldots, n-1
$$

So if $a=0$ there are no Darboux polynomials invariant by $\tau$ with nonzero cofactor for system (1) and the polynomial first integrals are polynomials in the variable $H=x^{2}+y^{2}+z^{2}$. This completes the proof of Theorem 12 .

Proof of Theorem 1. Let $f$ be an irreducible Darboux polynomial of degree $n$. In view of Lemma 10, the polynomial $g=f \cdot \tau(f)$ is a Darboux polynomial of system (1) with degree $2 n$ invariant by $\tau$. By Theorem 12 (i) if $a \neq 0$ this is not possible and so such $f$ cannot exists proving statement (i) of Theorem 1. If $a=0$ then in view of Theorem 12 (ii), the cofactor must be zero and $g=g\left(x^{2}+y^{2}+z^{2}\right)$. Since $x^{2}+y^{2}+z^{2}$ is irreducible, we get that both $f$ and $\tau(f)$ must be polynomials in the variable $x^{2}+y^{2}+z^{2}$ proving statement (ii) of Theorem 1.

To prove statement (iii) we proceed by contradiction. Assume that $G$ is a first integral of Darboux type of system (1). Then in view of its definition in (4) and taking into account statement (i) of Theorem 1 as well as Propositions 5 and 6 , we conclude that $G$ must be of the form

$$
G=\exp (f), \quad \text { where } f \in \mathbb{C}[x, y, z] \backslash \mathbb{C}
$$

Since $G$ is a first integral, after simplifying by $\exp (f)$, it must satisfy

$$
y \frac{\partial f}{\partial x}-(x+y z) \frac{\partial f}{\partial y}+\left(y^{2}-a\right) \frac{\partial f}{\partial z}=0
$$

So $f$ must be a polynomial first integral, which in view of statement (i) it is not possible proving statement (iii).

Now we prove statement (iv). Let $a=0$ and $G$ be a first integral of Darboux type of system (1). Then in view of its definition in (4) and taking into account statement (ii) of Theorem 1 as well as Propositions 5 and 6, we conclude that $G$ must be of the form

$$
G=\exp \left(f / g\left(x^{2}+y^{2}+z^{2}\right)\right)
$$

with $f \in \mathbb{C}[x, y, z], g$ a polynomial coprime with $f$. Since $G$ and $g\left(x^{2}+y^{2}+\right.$ $z^{2}$ ) are first integrals, after simplifying by $\exp \left(f / g\left(x^{2}+y^{2}+z^{2}\right)\right)$, we have

$$
y \frac{\partial f}{\partial x}-(x+y z) \frac{\partial f}{\partial y}+\left(y^{2}-a\right) \frac{\partial f}{\partial z}=0 .
$$

So $f$ must be a polynomial first integral, which in view of statement (ii) of Theorem 1 must be a polynomial in the variable $H=x^{2}+y^{2}+z^{2}$. This concludes the proof of statement (iv) and completes the proof of Theorem 1.

## 4. Proof of Theorem 2

We introduce the scaling

$$
X=\frac{x}{\sqrt{a}}, \quad Y=\frac{y}{\sqrt{a}}, \quad Z=\frac{z}{\sqrt{a}},
$$

where $a>0$. In these new variables system (1) becomes

$$
\begin{equation*}
\dot{X}=Y, \quad \dot{Y}=-X-\sqrt{a} Y Z, \quad \dot{Z}=\sqrt{a}\left(Y^{2}-1\right) \tag{17}
\end{equation*}
$$

Doing the cylindrical change of variables

$$
\begin{equation*}
X=r \cos \theta, \quad Y=r \sin \theta, \quad Z=Z \tag{18}
\end{equation*}
$$

in the region $r>0$, system (17) becomes

$$
\begin{equation*}
\dot{r}=-\sqrt{a} r Z \sin ^{2} \theta, \quad \dot{\theta}=-1-\sqrt{a} Z \cos \theta \sin \theta, \quad \dot{Z}=\sqrt{a}\left(-1+r^{2} \sin ^{2} \theta\right) \tag{19}
\end{equation*}
$$

It is easy to show that in a suitable neighborhood of $(r, Z)=(0,0)$ with $r>0$ we have $\dot{\theta} \neq 0$. Then choosing $\theta$ as the new independent variable, system (19) in a neighborhood of $(r, Z)=(0,0)$ with $r>0$ becomes

$$
\begin{aligned}
\frac{d r}{d \theta} & =\frac{\sqrt{a} r Z \sin ^{2} \theta}{1+\sqrt{a} Z \cos \theta \sin \theta}, \\
\frac{d Z}{d \theta} & =\frac{\sqrt{a}\left(1-r^{2} \sin ^{2} \theta\right)}{1+\sqrt{a} Z \cos \theta \sin \theta},
\end{aligned}
$$

Note that this system is periodic or period $2 \pi$ in $\theta$ and system (17) becomes

$$
\begin{align*}
\frac{d r}{d \theta} & =\sqrt{a} F_{11}(\theta, r, Z)+O(a) \\
\frac{d Z}{d \theta} & =\sqrt{a} F_{12}(\theta, r, Z)+O(a) \tag{20}
\end{align*}
$$

where

$$
F_{11}(\theta, r, Z)=r Z \sin ^{2} \theta, \quad F_{21}(\theta, r, Z)=1-r^{2} \sin ^{2} \theta
$$

Using the notation of the averaging theory in Theorem 7, we have that if we take $t=\theta, T=2 \pi, \varepsilon=\sqrt{a}, x=(r, Z)^{T}$ and

$$
\begin{aligned}
F_{1}(t, x) & =F_{1}(\theta, r, Z)=\binom{F_{11}(\theta, r, Z)}{F_{12}(\theta, r, Z)} \\
\varepsilon^{2} F_{2}(t, x) & =O(a)
\end{aligned}
$$

it is immediate to check that the differential system (20) is written in the normal form (5) for applying the averaging theory and that it satisfies the assumptions of Theorem 7. Now we compute the integral in (7) with $y=$ $(r, Z)^{T}$, and denoting

$$
g(y)=g(r, Z)=\binom{g_{11}(r, Z)}{g_{12}(r, Z)}
$$

we obtain

$$
\begin{aligned}
g_{11}(r, Z) & =\frac{r Z}{2} \\
g_{12}(r, Z) & =\frac{1}{2}\left(2-r^{2}\right)
\end{aligned}
$$

System $g_{11}(r, z)=g_{12}(r, Z)=0$ has a unique solution with $r>0$, namely $Z=0, r=\sqrt{2}$. The Jacobian (8) is

$$
\left|\begin{array}{ll}
Z / 2 & r / 2 \\
-r & 0
\end{array}\right|=\frac{r^{2}}{2}
$$

and evaluated on the solution $(r, Z)=(\sqrt{2}, 0)$ takes the value 1 . Then by Theorem 7 it follows that for any $a>0$ sufficiently small, system (20) has a periodic solution $(r(\theta, a), Z(\theta, a))$ such that $(r(0, a), Z(0, a))$ tends to $(\sqrt{2}, 0)$ when $a$ tends to zero.

Going back to the differential system (19) we get that this system for $a>0$ sufficiently small has one periodic solution of period approximately $2 \pi$ of the form

$$
r(t)=\sqrt{2}+O(\sqrt{a}), \quad \theta(t)=-t+O(\sqrt{a}), \quad Z(t)=O(\sqrt{a})
$$

This periodic solution become for the differential system (18) into one periodic solution of period also close to $2 \pi$ of the form

$$
X(t)=\sqrt{2} \cos t+O(\sqrt{a}), \quad Y(t)=-\sqrt{2} \sin t+O(\sqrt{a}), \quad Z(t)=O(\sqrt{a})
$$

for $a>0$ sufficiently small. Finally, we get for the differential system (1) the periodic solution

$$
x(t)=\sqrt{2 a} \cos t+O(a), \quad y(t)=-\sqrt{2 a} \sin t+O(a), \quad z(t)=O(a)
$$

of period near $2 \pi$ when $a>0$ sufficiently small. Clearly this periodic orbit tends to the equilibrium point $(0,0,0)$ when $a$ tends to zero. This concludes the proof of the theorem.

## Acknowledgements

The first and third authors are partially supported by FCT/Portugal through UID/MAT/04459/2013. The second author is partially supported by a FEDER-MINECO grant MTM2016-77278-P, a MINECO grant MTM2013-40998-P, and an AGAUR grant 2014SGR-568.

## References

[1] D. Bleecker and G. Csordas, Basic Partial Differential Equations Van Nostrand Reinhold, New York, 1992.
[2] V. Castellanos, V., J. Llibre, and I. Quilantán, Simultaneous periodic orbits bifurcating from two Zero-Hopf equilibria in a tritrophic food chain model J. Appl. Math. Phys. 1 (2013), 31-38.
[3] A.R. Champneys and V. Kirk, The entwined wiggling of homoclinic curves emerging from saddle-node/Hopf instabilities, Physica D 195 (2004), 77-105.
[4] C. Christopher, J. Llibre and J.V. Pereira, Multiplicity of invariant algebraic curves in polynomial vector fields, Pacific J. Math. 229 (2007), 63-117.
[5] G. Darboux, Mémoire sur les équations différentielles algébriques du premier ordre et du premier degrée (Mélanges). Bull. Sci. Math. 2ème série. 2 (1878), 60-96; 123144; 151-200.
[6] F. Dumortier, J. Llibre and J.C. Artés, Qualitative Theory of Planar Differential Systems. Springer, New York (2006).
[7] J. Guckenheimer, On a codimension two bifurcation, Lecture Notes in Math. 898 (1980), 99-142.
[8] J. Guckenheimer and P. Holmes, Nonlinear Oscillations, Dynamical Systems and Bifurcations of Vector Fields, Springer, 1983.
[9] W.G. Hoover, Canonical dynamics: Equilibrium phase-space distributions, Phys. Rev. A 31 (1985), 1695-1697.
[10] W.G. Hoover, Remark on "Some simple chaotic flows", Phys. Rev. E. 51 (1995), 759-760.
[11] J.P. Jouanolou, Equations de Pfaff algébriques. Lectures Notes in Mathematics, vol. 708. Springer, Berlin (1979).
[12] Y.A. Kuznetsov, Elements of Applied Bifurcation Theory, Springer-Verlag, 3rd edition, 2004.
[13] J. Llibre, R. Moeckel and C. Simó, Central configurations, periodic orbits and Hamiltonian systems, Advances Courses in Math. CRM Barcelona, Birkhäuser, 2015.
[14] J. Llibre and X. Zhang, Darboux theory of integrability in $\mathbb{C}^{n}$ taking into account the multiplicity, J. Differential Equations 246 (2009), 541-551.
[15] J. Llibre and X. Zhang, Darboux theory of integrability for polynomial vector fields in $\mathbb{R}^{n}$ taking into account the multiplicity at infinity, Bull. Sci. Math. 133 (2009), 765-778.
[16] M. Messias and A. Reinol, On the formation of hidden chaotic attractors and nested invariant tori in the Sprott A system, to appear in Nonlinear Dynamics.
[17] S. NosÉ, A unified formulation of the constant temperature molecular dynamics methods, J. Chem. Phys. 81 (1984), 511-519.
[18] H.A. Posch, W.G. Hoover and F.J. Vesely, Canonical dynamics of the Nosé oscillator: stability, order and chaos. Phys. Rev. A 33 (1986), 4253-4265.
[19] J.A. Sanders, F. Verhulst and J. Murdock, Averaging methods in nonlinear dynamical systems, Second edition, Applied Mathematical Sciences 59, Springer, New York, 2007.
[20] J. Scheurle and J. Marsden, Bifurcation to quasi-periodic tori in the interaction of steady state and Hopf bifurcations, SIAM. J. Math. Anal. 15 (1984), 1055-1074.
[21] J.C. Sprott, Some simple chaotic flows. Phsy. Rev. E 50, (1994), R647-R650.
[22] F. Verhulst, Nonlinear Differential Equations and Dynamical Systems. Springer, Berlin (1991).
${ }^{1}$ Departamento de Matemática, Instituto Superior Técnico, Universidade Técnica de Lisboa, Av. Rovisco Pais 1049-001, Lisboa, Portugal

E-mail address: barreira@math.ist.utl.pt, cvalls@math.ist.utl.pt
${ }^{2}$ Departament de Matemàtiques, Universitat Autònoma de Barcelona, 08193
Bellaterra, Barcelona, Catalonia, Spain
E-mail address: jllibre@mat.uab.cat


[^0]:    2010 Mathematics Subject Classification. Primary 34C05, 34A34, 34C14.
    Key words and phrases. Darboux integrability, Sprott A system, zero-Hopf bifurcation, averaging theory.

