

## PHASE PORTRAITS OF THE HIGGINS–SELKOV SYSTEM

JAUME LLIBRE

Departament de Matemàtiques, Universitat Autònoma de Barcelona  
08193 Bellaterra, Barcelona, Catalonia, Spain

MARZIEH MOUSAVI

Department of Mathematical Sciences, Isfahan University of Technology,  
Isfahan 84156-83111, Iran

ABSTRACT. In this paper we study the dynamics of the Higgins–Selkov system

$$\dot{x} = 1 - xy^\gamma, \quad \dot{y} = \alpha y(xy^{\gamma-1} - 1),$$

where  $\alpha$  is a real parameter and  $\gamma > 1$  is an integer. We classify the phase portraits of this system for  $\gamma = 3, 4, 5, 6$ , in the Poincaré disc for all the values of the parameter  $\alpha$ . Moreover, we determine in function of the parameter  $\alpha$  the regions of the phase space with biological meaning.

### 1. INTRODUCTION AND STATEMENT OF THE MAIN RESULTS

The *glycolysis* is the first step in the transformation of glucose into energy for cellular metabolism. Higgins in [9] used a mathematical model for investigating sustained oscillations in the glycolysis process. But his model did not show any limit cycle corresponding to sustained oscillations observed in the experiments. For solving this problem Selkov in [12] introduced another mathematical model for studying the glycolysis now called the Higgins–Selkov system. This system is

$$\frac{dx}{dt} = 1 - xy^\gamma, \quad \frac{dy}{dt} = \alpha y(xy^{\gamma-1} - 1), \quad (1)$$

where  $x$  and  $y$  are state variables,  $t$  is the time variable,  $\alpha$  is a real parameter and  $\gamma > 1$  is an integer.

The point  $(1, 1)$  is the unique singular point of system (1) for all values of parameter  $\alpha \neq 0$  and  $\gamma > 1$ . It is shown that a supercritical Hopf bifurcation occurs at  $\alpha = 1/(\gamma - 1)$ . In 2018, Artés et al. [2] characterized the dynamics of system (1) with  $\gamma = 2$  and  $\alpha \in \mathbb{R} \setminus (1, 3)$ , and proposed a conjecture about existence and uniqueness of a limit cycle when  $\alpha \in (1, 3)$ . Chen and Tang in [6] have proved this conjecture which completes the phase portraits of system (1) with  $\gamma = 2$ . Brechmann and Rendall in [3] researched the uniqueness of limit cycles, and additionally they proved that there are no limit cycles when  $\alpha \in (0, 1/(\gamma - 1))$  for all  $\gamma > 1$ . Recently, it appears in arXiv the preprint of Brechmann and Rendall [4] where they show that for  $\alpha > 0$  the top right hand quarter of each phase portrait is topologically equivalent to that for  $\gamma = 2$  for any integer  $\gamma > 1$ .

---

2010 *Mathematics Subject Classification.* 34C05, 34C07, 34C08.

*Key words and phrases.* Higgins–Selkov system, limit cycle, Phase portrait, Poincaré compactification.

In this paper we provide a complete study of the dynamics of system (1) for  $\gamma = 3, 4, 5$  and 6, and characterize the existence or non–existence of limit cycles in function of the parameter  $\alpha$ . Moreover, we provide the qualitative behaviour of the system in a neighborhood of the infinity for all values of the integer  $\gamma > 1$  using the Poincaré compactification. Finally we classify the phase portraits of system (1) in the Poincaré disc for  $\gamma = 3, 4, 5$  and 6.

For a definition of the *Poincaré compactification* and the *Poincaré disc* see Chapter 5 of [7], for the notions of *separatrix*, *canonical regions* and the concept of *topologically equivalent differential systems* see section 1.9 of [7]. Here, the number of separatrices and canonical region of a phase portrait will be denoted by  $S$  and  $R$ , respectively.

Our main results are the following two theorems.

**Theorem 1.** *The phase portraits of the Higgins–Selkov system (1) in the Poincaré disc for  $\gamma = 3$  and 5 are topologically equivalent to one of the phase portraits in Figure 1 as follows:*

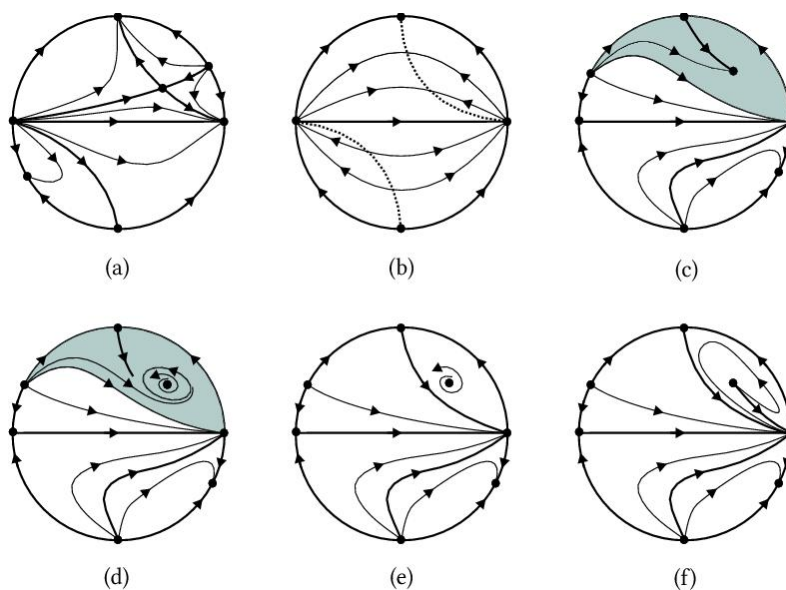
- (i) *Figure 1(a) corresponds to  $\alpha < 0$  with  $S = 19$  and  $R = 6$ ,*
- (ii) *Figure 1(b) corresponds to  $\alpha = 0$  with  $S = \infty$ ,*
- (iii) *Figure 1(c) corresponds to  $\alpha \in (0, 1/(\gamma - 1)]$  with  $S = 17$  and  $R = 4$ ,*
- (iv) *Figure 1(d) corresponds to  $\alpha \in (1/(\gamma - 1), \alpha_\gamma^*)$  where  $\alpha_\gamma^* \in (1/(\gamma - 1), \alpha_\gamma)$ ,  $\alpha_3 = (1261 + 57\sqrt{57})/686$  and  $\alpha_5 = 24(411 + 41\sqrt{41})/3125$  with  $S = 18$  and  $R = 5$ ,*
- (v) *Figure 1(e) corresponds to  $\alpha = \alpha_\gamma^*$  with  $S = 16$  and  $R = 4$ ,*
- (vi) *Figure 1(f) corresponds to  $\alpha > \alpha_\gamma^*$  with  $S = 17$  and  $R = 4$ .*

Moreover for all odd integer  $\gamma > 1$  the phase portraits of Figure 1(a) – (c) hold for  $\alpha < 0$ ,  $\alpha = 0$  and  $\alpha \in [0, 1/(\gamma - 1))$ , respectively; and the qualitative behaviour of system (1) at infinity is the same for any given  $\alpha \in \mathbb{R}$ .

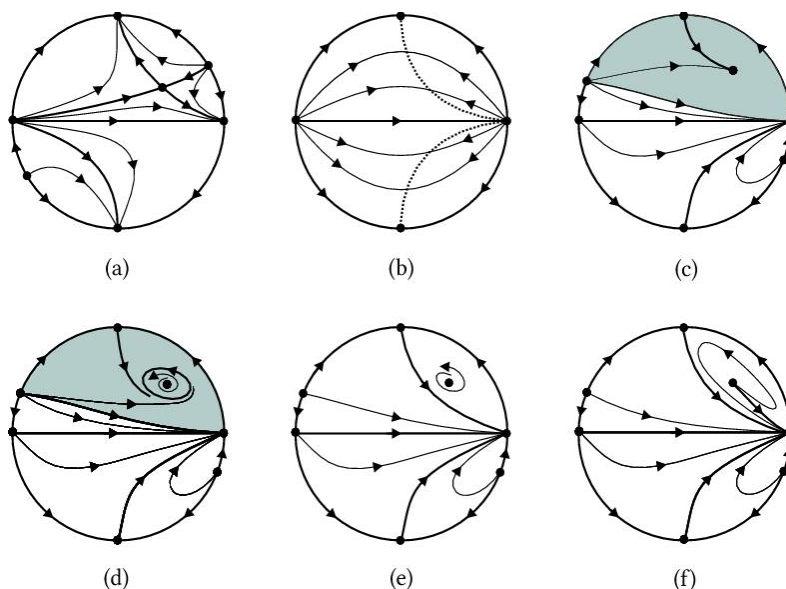
**Theorem 2.** *The phase portraits of the Higgins–Selkov system (1) in the Poincaré disc for  $\gamma = 4$  and 6 are topologically equivalent to one of the phase portraits in Figure 2 as follows:*

- (i) *Figure 2(a) corresponds to  $\alpha < 0$  with  $S = 19$  and  $R = 6$ ,*
- (ii) *Figure 2(b) corresponds to  $\alpha = 0$  with  $S = \infty$ ,*
- (iii) *Figure 2(c) corresponds to  $\alpha \in (0, 1/(\gamma - 1)]$  with  $S = 17$  and  $R = 4$ ,*
- (iv) *Figure 2(d) corresponds to  $\alpha \in (1/(\gamma - 1), \alpha_\gamma^*)$  where  $\alpha_\gamma^* \in (1/(\gamma - 1), \alpha_\gamma)$ ,  $\alpha_4 = 46875/14641$  and  $\alpha_6 = (189384302 + 10121703\sqrt{681})/47045881$  with  $S = 18$  and  $R = 5$ ,*
- (v) *Figure 2(e) corresponds to  $\alpha = \alpha_\gamma^*$  with  $S = 16$  and  $R = 4$ ,*
- (vi) *Figure 2(f) corresponds to  $\alpha > \alpha_\gamma^*$  with  $S = 17$  and  $R = 4$ .*

Moreover for all even integer  $\gamma > 1$  the phase portraits of Figure 1(a) – (c) hold for  $\alpha < 0$ ,  $\alpha = 0$  and  $\alpha \in [0, 1/(\gamma - 1))$ , respectively; and the qualitative behaviour of system (1) at infinity is the same for any given  $\alpha \in \mathbb{R}$ .



**Figure 1.** The phase portraits of system (1) for  $\gamma = 3$  and 5 in the Poincaré disc. The shaded areas correspond to the initial conditions of the orbits having a final finite evolution, so these are the initial conditions with biological meaning. In the phase portrait (c) the final behaviour is a stable singular point, and in the phase portrait (d) is a stable limit cycle.



**Figure 2.** The phase portraits of system (1) for  $\gamma = 4$  and 6 in the Poincaré disc. The shaded areas correspond to the initial conditions of the orbits having a final finite evolution, so these are the initial conditions with biological meaning. In the phase portrait (c) the final behaviour is a stable singular point, and in the phase portrait (d) is a stable limit cycle.

1

## 2. THE PROOF OF THEOREMS 1 AND 2

**2.1. The local phase portraits of the finite singular points.** The Higgins–Selkov system (1) has a unique singular point  $(1, 1)$  and the Jacobian matrix of the

system at this point is

$$J = \begin{pmatrix} -1 & -\gamma \\ \alpha & \alpha(\gamma - 1) \end{pmatrix}.$$

The determinant and trace of the matrix  $J$  are equal to  $\alpha$  and  $\alpha(\gamma - 1) - 1$ , respectively. The eigenvalues of the Jacobian matrix are

$$\frac{1}{2} \left[ \alpha(\gamma - 1) - 1 \pm \sqrt{[\alpha(\gamma - 1) - 1]^2 - 4\alpha} \right].$$

Therefore the singular point  $(1, 1)$  is

- a hyperbolic saddle if  $\alpha < 0$ ;
- a stable hyperbolic node if  $\alpha \in (0, (\gamma + 1 - 2\sqrt{\gamma})/(\gamma - 1)^2]$ ;
- a stable focus if  $\alpha \in ((\gamma + 1 - 2\sqrt{\gamma})/(\gamma - 1)^2, 1/(\gamma - 1)]$ ;
- a stable weak focus if  $\alpha = 1/(\gamma - 1)$ , in this case a Hopf bifurcation occurs;
- an unstable hyperbolic focus if  $\alpha \in (1/(\gamma - 1), (\gamma + 1 + 2\sqrt{\gamma})/(\gamma - 1)^2]$ ; and
- an unstable hyperbolic node if  $\alpha \geq (\gamma + 1 + 2\sqrt{\gamma})/(\gamma - 1)^2$ .

For  $\alpha = 1/(\gamma - 1)$  the Jacobian matrix  $J$  has a pair of imaginary eigenvalues. If it is considered the real part of the eigenvalues as a function of parameter  $\alpha$  then the derivative of this function with respect to  $\alpha$  at  $\alpha = 1/(\gamma - 1)$  is nonzero. Therefore, according to Theorem 3.4.2 of [8] a Hopf bifurcation occurs. Then, by using the formula (3.4.11) of [8] we obtain the first Lyapunov coefficient which is equal to  $-1/16$ . Hence, we have a supercritical bifurcation for  $\alpha$  slightly larger than  $1/(\gamma - 1)$ , and the unique bifurcated limit cycle is stable. Moreover, if  $\alpha = 0$ , then all the points on the curve  $x = y^{-\gamma}$  are singular points of system (1) and all the parallel lines to the  $x$ -axis are invariant, therefore the phase portraits of system (1) are given in Figure 1(b) and 2(b) corresponding to the cases that the integer  $\gamma$  is odd and even, respectively.

**2.2. The local phase portraits of the infinite singular points.** In order to study the qualitative behaviour of the solutions of 2-dimensional polynomial differential systems at infinity, we can use the Poincaré compactification technique. The base of this technique is projecting a vector field in  $\mathbb{R}^2$  onto a sphere which is a compact set. For more details about the Poincaré compactification, see chapter 5 of [7].

In the local chart  $U_1$  system (1) becomes

$$\dot{u} = -uv^{\gamma+1} + u^{\gamma+1} + \alpha u^\gamma - \alpha uv^\gamma, \quad \dot{v} = -v(v^{\gamma+1} - u^\gamma). \quad (2)$$

Let  $\alpha \neq 0$ , then the infinite singular points of system (1) in local chart  $U_1$  are  $(0, 0)$  and  $(-\alpha, 0)$ . The Jacobian matrix of system (2) at point  $(0, 0)$  is identically zero. The blow up technique is used for studying the dynamics of the system in a neighborhood of the linearly zero singular points. For more details about the blow up, see [1]. Here we shall apply the blow up technique separately when the integer  $\gamma$  is odd and even.

If  $\gamma$  is odd, for desingularizing the origin of system (2), we use two vertical blow ups. First we do a vertical blow up  $(u, v) \mapsto (w_1 v, v)$  where  $w_1$  is the new variable. Therefore system (2) in the new variables and after cancelling the common factor  $v^{\gamma-1}$  becomes

$$\dot{w}_1 = -\alpha w_1(v - w_1^{\gamma-1}), \quad \dot{v} = -v^2(v - w_1^\gamma), \quad (3)$$

where the origin is the unique singular point on the line  $v = 0$ , and the Jacobian matrix of the system at  $(0, 0)$  is again identically zero. Hence, we use another vertical

blow up  $(w_1, v) \mapsto (w_2v, v)$  where  $w_2$  is the new variable. Thus system (3) in the new variables and after eliminating the common factor  $v$  is

$$\dot{w}_2 = w_2(v - \alpha - v^\gamma w_2^\gamma + \alpha v^{\gamma-2} w_2^{\gamma-1}), \quad \dot{v} = v^2(-1 + v^{\gamma-1} w_2^\gamma). \quad (4)$$

The origin is the unique singular point of system (4) on the line  $v = 0$ , it is semi-hyperbolic and its local phase portrait is determined using Theorem 2.19 of [7]. By going back from  $(w_2, v)$  to  $(w_1, v)$  and from  $(w_1, v)$  to  $(u, v)$ , we get that the point  $(0, 0)$  of the original system (2) is the union of one parabolic and four hyperbolic sectors when  $\alpha > 0$ , see Figure 1(c)-(f); and the origin is a stable node when  $\alpha < 0$ , see Figure 1(a).

If  $\gamma$  is even we do a horizontal blow up  $(u, v) \mapsto (u, uw)$  where  $w$  is the new variable. Then system (2) in the new variables writes

$$\dot{u} = u(u - u^2 w^{\gamma+1} + \alpha - \alpha u w^\gamma), \quad \dot{w} = \alpha w(-1 + u w^\gamma). \quad (5)$$

where by rescaling of time, the common factor  $u^{\gamma-1}$  has been eliminated. The point  $(0, 0)$  is the unique singular point of system (5) on  $u = 0$ , it is a hyperbolic saddle when  $\alpha \neq 0$ . Now by performing a blowing down and returning from  $(u, w)$  to  $(u, v)$ , we obtain that the origin of local chart  $U_1$  is formed by two hyperbolic and two parabolic sectors, see Figure 2(a) and 2(c)-(f).

Now we consider the singular point  $(-\alpha, 0)$  of system (2) in the local chart  $U_1$ . If  $\gamma$  is odd then the Jacobian matrix at this point has two positive eigenvalues when  $\alpha < 0$ . Therefore  $(-\alpha, 0)$  is a hyperbolic unstable node, see Figure 1(a). Also, when  $\alpha > 0$ , the eigenvalues of the Jacobian matrix are negative and this implies that the point  $(-\alpha, 0)$  is a hyperbolic stable node, see Figure 1(c)-(f).

If  $\gamma$  is even, then the Jacobian matrix of system (2) at the point  $(-\alpha, 0)$  has two positive eigenvalues when  $\alpha \neq 0$ . It implies that  $(-\alpha, 0)$  is a hyperbolic unstable node, see Figure 2(a) and 2(c)-(f).

In the local chart  $U_2$  system (1) becomes

$$\dot{u} = v^{\gamma+1} - u - \alpha u^2 + \alpha u v^\gamma, \quad \dot{v} = -\alpha v(u - v^\gamma). \quad (6)$$

The origin is a singular point of system (6) and the Jacobian matrix at this point has a unique zero eigenvalue. Therefore the origin is a semi-hyperbolic singular point and for determining the local behaviour of system (6) in the neighborhood of the origin we use Theorem 2.19 of [7] and we obtain that if  $\gamma$  is odd, then the origin is union of one parabolic and two hyperbolic sectors when  $\alpha \neq 0$ , see Figure 1. In the case  $\gamma$  is even, the origin is a saddle when  $\alpha > 0$  and a stable node when  $\alpha < 0$ , see Figure 2.

**2.3. The periodic orbits.** According to Corollary 2 of section 3.12 of [11], we know that the region limited by a periodic orbit of a  $C^1$  planar differential system contains at least one singular point and if there are a finite number of singular points in this region, then the sum of the topological indices of these singular points is equal to one.

Since system (1) for  $\alpha < 0$  has only one finite singular point (a saddle point) and the index of this point is  $-1$  (see Theorem 7 of section 3.12 of [11]), then system (1) for  $\alpha < 0$  has no periodic orbits. Moreover if  $\alpha = 0$  then  $\dot{y} = 0$ , and in this case system (1) has no periodic solutions.

**Lemma 3.** *The Higgins–Selkov system (1) for  $\gamma = 3, 4, 5$  and  $6$ , has no periodic solutions if  $\alpha > \alpha_\gamma$ , where  $\alpha_3 = (1261 + 57\sqrt{57})/686$ ,  $\alpha_4 = 46875/14641$ ,  $\alpha_5 = 24(411 + 41\sqrt{41})/3125$  and  $\alpha_6 = (189384302 + 10121703\sqrt{681})/47045881$ .*

*Proof.* For system (1) we have

$$\dot{x}|_{x=0} = 1 > 0 \quad \text{and} \quad \dot{y}|_{y=0} = 0,$$

and also if there exists a periodic orbit, it must surround the singular point  $(1, 1)$ . Therefore if system (1) has a periodic solution, then it must be contained in the quadrant  $Q = \{(x, y) \in \mathbb{R}^2 : x > 0, y > 0\}$  for all  $\gamma > 1$ . The divergence of system (1) is

$$f(x, y) = -y^\gamma + \alpha\gamma xy^{\gamma-1} - \alpha.$$

The curve  $f(x, y) = 0$  separates the quadrant  $Q$  into two open components  $f(x, y) > 0$  and  $f(x, y) < 0$ . We prove that for  $\gamma = 3, 4, 5, 6$ , this curve  $\{(x, y) : f(x, y) = 0\} \cap Q$  is transversal to the flow of system (1) if  $\alpha > (1261 + 57\sqrt{57})/686$ ,  $46875/14641$ ,  $24(411 + 41\sqrt{41})/3125$  and  $(189384302 + 10121703\sqrt{681})/47045881$ , respectively. Therefore a periodic orbit cannot intersect the curve  $f(x, y) = 0$ , and it must be contained in a simple connected region either  $\{f(x, y) > 0\} \cap Q$ , or  $\{f(x, y) < 0\} \cap Q$ . Then the existence of a such periodic solution is impossible by the Bendixson criterion, see Theorem 7.10 of [7].

By solving  $f(x, y) = 0$  we obtain  $x = (y^\gamma + \alpha)/\alpha\gamma y^{\gamma-1}$ . Consider the function

$$\begin{aligned} T_\alpha(x, y) &= \left( \frac{\partial f}{\partial x} \dot{x} + \frac{\partial f}{\partial y} \dot{y} \right) \Big|_{x = \frac{y^\gamma + \alpha}{\alpha\gamma y^{\gamma-1}}} := \frac{1}{\gamma} P_\alpha(y) \\ &= \frac{-(1 + \gamma)y^{2\gamma} + \alpha(\gamma - 2)y^\gamma + \alpha\gamma^2 y^{\gamma-1} - \alpha^2(\gamma - 1)^2}{\gamma}. \end{aligned}$$

Hence for all positive integer  $\gamma$ ,  $P_\alpha$  is a polynomial of degree  $2\gamma$ . In order to prove that  $T_\alpha(x, y)$  in the first quadrant does not change sign, it is sufficient to verify that polynomial  $P_\alpha$  has no real roots. Here, we use the notions of *discrimination matrix*, *discriminant sequence* and Theorem 8 in the Appendix to prove that  $P_\alpha$  for  $\gamma = 3, 4, 5, 6$  does not have real roots.

First we consider  $\gamma = 3$  and prove that the polynomial  $P_\alpha$  of degree 6 for  $\alpha > (1261 + 57\sqrt{57})/686$  has no real roots. The elements of the discriminant sequence of polynomial  $P_\alpha$  are

$$\begin{aligned} D_1 &= d_2 = 96, \\ D_2 &= d_4 = 0, \\ D_3 &= d_6 = -13824\alpha^2, \\ D_4 &= d_8 = 186624\alpha^3(-1536 + 7\alpha), \\ D_5 &= d_{10} = 30233088\alpha^5(-1024 + 621\alpha), \\ D_6 &= d_{12} = -2176782336\alpha^8(1024 - 1261\alpha + 343\alpha^2). \end{aligned}$$

Therefore, according to the Appendix for  $\alpha > (1261 + 57\sqrt{57})/686$  the sign list of sequence  $\{D_1, D_2, \dots, D_6\}$  is  $[1, 0, -1, \text{sign}(D_4), 1, -1]$ , so we have the following *revised sign list*

$$[1, -1, -1, 1 \text{ or } -1, 1, -1].$$

Then for any sign of the fourth element in the above list, the number of the sign changes in the revised list is 3. Hence, according to Theorem 8,  $P_\alpha(y)$  has 3 pairs of distinct conjugate imaginary roots. Hence it has no real roots and this implies that  $T_\alpha(x, y) < 0$  in  $Q$ . This concludes the proof of the lemma for  $\gamma = 3$ .

Now we consider  $\gamma = 4$ . Again by using discrimination matrix and Theorem 8 we show that the polynomial  $P_\alpha$  of degree 8 for  $\alpha > 46875/14641$  has no real roots. The elements of the discriminant sequence of polynomial  $P_\alpha$  are

$$\begin{aligned} D_1 &= d_2 = 200, \\ D_2 &= d_4 = 0, \\ D_3 &= d_6 = 0, \\ D_4 &= d_8 = -12800000\alpha^3, \\ D_5 &= d_{10} = 409600000\alpha^4(12500 + 11\alpha), \\ D_6 &= d_{12} = 39321600000\alpha^6(62500 + 4387\alpha), \\ D_7 &= d_{14} = -11324620800000\alpha^9(-15625 + 48257\alpha), \\ D_8 &= d_{16} = 9784472371200000\alpha^{12}(1 + \alpha)(-46875 + 14641\alpha). \end{aligned}$$

For  $\alpha > 46875/14641$  the sign list of sequence  $\{D_1, D_2, \dots, D_8\}$  is  $[1, 0, 0, -1, 1, 1, -1, 1]$ , Therefore the corresponding revised sign list is

$$[1, -1, -1, -1, 1, 1, -1, 1].$$

The number of the sign changes in this revised list is 4. Thus  $P_\alpha(y)$  has 4 pairs of distinct conjugate imaginary roots and no real roots. It implies that  $T_\alpha(x, y) < 0$  in  $Q$ . This concludes the proof of the lemma for  $\gamma = 4$ .

If  $\gamma = 5$  in a similar way to the previous cases, we show that the polynomial  $P_\alpha$  of degree 10 for  $\alpha > 24(411 + 41\sqrt{41})/3125$  has only imaginary roots. The elements of the discriminant sequence of the polynomial  $P_\alpha$  are

$$\begin{aligned} D_1 &= d_2 = 360, \\ D_2 &= d_4 = 0, \\ D_3 &= d_6 = 0, \\ D_4 &= d_8 = 0, \\ D_5 &= d_{10} = 23619600000\alpha^4, \\ D_6 &= d_{12} = -22143375000000\alpha^5(-9600 + \alpha), \\ D_7 &= d_{14} = -6643012500000000\alpha^7(-48000 + 343\alpha), \\ D_8 &= d_{16} = -32651735040000000000000\alpha^{10}, \\ D_9 &= d_{18} = 26572050000000000000000\alpha^{13}(-20352 + 8575\alpha), \\ D_{10} &= d_{20} = -7085880000000000000000000\alpha^{16}(1 + \alpha)(18432 - 19728\alpha + 3125\alpha^2). \end{aligned}$$

For  $\alpha > 24(411 + 41\sqrt{41})/3125$  the sign list of sequence  $\{D_1, D_2, \dots, D_{10}\}$  is

$$[1, 0, 0, 0, 1, \text{sign}(D_6), \text{sign}(D_7), -1, 1, -1],$$

so the corresponding revised sign list is

$$[1, -1, -1, 1, 1, 1 \text{ or } -1, 1 \text{ or } -1, -1, 1, -1].$$

The case that  $\text{sign}(D_6) = -1$  and  $\text{sign}(D_7) = 1$  is impossible, in the other cases the number of the sign changes in this revised list is 5. Therefore  $P_\alpha(y)$  has 5 pairs of distinct conjugate imaginary roots and no real roots. It implies that  $T_\alpha(x, y) < 0$  in  $Q$ . This completes the proof of the lemma for  $\gamma = 5$ .

For  $\gamma = 6$  the elements of discriminant sequence of the polynomial  $P_\alpha$  of degree 12 are

$$\begin{aligned} D_1 &= d_2 = 588, \\ D_2 &= d_4 = 0, \\ D_3 &= d_6 = 0, \\ D_4 &= d_8 = 0, \\ D_5 &= d_{10} = 0, \\ D_6 &= d_{12} = 78690759081984\alpha^5, \\ D_7 &= d_{14} = -265581311901696\alpha^6(66706983 + 608\alpha), \\ D_8 &= d_{16} = -430241725280747520\alpha^8(155649627 + 93272\alpha), \\ D_9 &= d_{18} = -15944515252329719468851200000\alpha^{11}, \\ D_{10} &= d_{20} = 408976816222257304376033280000000\alpha^{14}, \\ D_{11} &= d_{22} = -753691861204384855818240000000\alpha^{17}(-1647086 + 898529\alpha), \\ D_{12} &= d_{24} = 38108015454154290462720000000000\alpha^{20}(-720600125 - 378768604\alpha \\ &\quad + 47045881\alpha^2). \end{aligned}$$

For  $\alpha > (189384302 + 10121703\sqrt{681})/47045881$  the sign list of sequence  $\{D_1, D_2, \dots, D_{12}\}$  is  $[1, 0, 0, 0, 0, 1, -1, -1, -1, 1, -1, 1]$ , then the corresponding revised sign list is

$$[1, -1, -1, 1, 1, 1, -1, -1, -1, 1, -1, 1].$$

The number of the sign changes in this revised list is 6. According to Theorem 8,  $P_\alpha(y)$  has 6 pairs of distinct conjugate imaginary roots. It implies that  $T_\alpha(x, y) < 0$  in  $Q$ . This completes the proof of the lemma for  $\gamma = 6$ .  $\square$

To simplify the investigation about the limit cycles of the Higgins–Selkov system (1), it is introduced the following change of variables

$$x = \frac{\alpha - X - Y}{\alpha}, \quad y = 1 + X, \quad \frac{dt}{d\tau} = \frac{1}{y^\gamma}.$$

By performing this change of variables and again denoting  $X, Y$  and  $\tau$  by  $x, y$  and  $t$ , respectively, system (1) transforms into the following Liénard differential system

$$\frac{dx}{dt} = -y + \alpha - x - \frac{\alpha}{(1+x)^{\gamma-1}}, \quad \frac{dy}{dt} = \frac{\alpha x}{(1+x)^\gamma}. \quad (7)$$

In [3] by considering Liénard system (7) for  $x > -1$ , it was proved that if the unique singular point of system (1) is unstable for a given value of parameter  $\alpha$ , then there is at most one periodic orbit and if such a solution exists, it must be asymptotically stable. Moreover the authors of [3], by means of Theorem 2.1 of [10], have shown if  $\alpha \in (0, 1/(\gamma - 1))$  then system (7) has no limit cycles. This implies that when the unique singular point of the Higgins–Selkov system (1) is stable, then there is no limit cycles (see Theorem 7 of [3]).



In the following result we show that system (7) and consequently the Higgins–Selkov (1) in the case  $\gamma = 3, 4, 5, 6$  has no limit cycle when  $\alpha = 1/(\gamma - 1)$ .

**Lemma 4.** *If  $\alpha = 1/(\gamma - 1)$  then system (7) for  $\gamma = 3, 4, 5, 6$  has no limit cycles.*

*Proof.* By doing the change of variable  $y \rightarrow -y$  on system (7), we can use Corollary 10 of [5] such that

$$\bar{F}(x) = \alpha - x - \frac{\alpha}{(1+x)^{\gamma-1}}, \quad \bar{g}(x) = \frac{\alpha x}{(1+x)^\gamma},$$

and  $x \in (-1, \infty)$ . If  $\alpha = 1/(\gamma - 1)$ , it can be easily checked that the conditions (i)-(iii) of Proposition 9 of [5] hold. To complete this proof, it suffices to show that there is no solution for the system

$$\bar{F}(x_1) = \bar{F}(x_2), \quad \bar{\lambda}(x_1) = \bar{\lambda}(x_2), \quad (8)$$

where  $\bar{F}'(x) = \bar{f}(x)$ ,  $\bar{\lambda}(x) = \bar{g}(x)/\bar{f}(x)$ , and  $-1 < x_2 < 0 < x_1$ . From  $\bar{\lambda}(x_1) = \bar{\lambda}(x_2)$  and  $\alpha = 1/(\gamma - 1)$ , the following equation is obtained

$$(x_2 - x_1) + x_1(1 + x_2)^\gamma - x_2(1 + x_1)^\gamma = 0,$$

by using the binomial expansion and after simplifying, we can rewrite the above equation as

$$x_1 x_2 \left[ \binom{\gamma}{2} (x_2 - x_1) + \binom{\gamma}{3} (x_2^2 - x_1^2) + \dots + \binom{\gamma}{\gamma} (x_2^{\gamma-1} - x_1^{\gamma-1}) \right] = 0, \quad (9)$$

where  $\binom{\gamma}{n} = \frac{\gamma!}{(\gamma-n)!n!}$  and  $n < \gamma$  is a positive integer. Now for  $\gamma = 3, 4, 5, 6$ , we show that there is no solution for equation (9) such that  $-1 < x_2 < 0 < x_1$ , thus there is no solution for system (8).

By substituting  $\gamma = 3, 4, 5, 6$  into (9), we obtain

- $x_1 x_2 (x_2 - x_1)(3 + x_1 + x_2) = 0$  for  $\gamma = 3$ ,
- $x_1 x_2 (x_2 - x_1) [(6 + 4x_2 + x_2^2) + x_1(4 + x_2) + x_1^2] = 0$  for  $\gamma = 4$ ,
- $x_1 x_2 (x_2 - x_1) [10 + x_1^3 + x_2(10 + 5x_2 + x_2^2) + x_1^2(5 + x_2) + x_1(10 + 5x_2 + x_2^2)] = 0$  for  $\gamma = 5$ ,
- $x_1 x_2 (x_2 - x_1) [(15 + 20x_2 + 15x_2^2 + 6x_2^3) + x_2^4 + x_1^3(6 + x_2) + x_1^2(15 + 6x_2 + x_2^2) + x_1(20 + 15x_2 + 6x_2^2 + x_2^3) + x_1^4] = 0$  for  $\gamma = 6$ .

In all four equations listed above, the condition  $-1 < x_2 < 0 < x_1$  implies that there is no solution for these equations. Therefore system (8) has no solutions and by using Corollary 10 of [5] we conclude that system (7) for  $\gamma = 3, 4, 5, 6$  and  $\alpha = 1/(\gamma - 1)$  has no limit cycles.  $\square$

Now we study the existence of limit cycle of system (1) when  $\alpha > 1/(\gamma - 1)$ . For this purpose we perform the linear change of variables  $(x, y, t) \rightarrow (x, \sqrt{\alpha}y, -1/\sqrt{\alpha}t)$  in system (7) when  $\alpha > 0$ , and we obtain the following system

$$\frac{dx}{dt} = y + \sqrt{\alpha} - \frac{1}{\sqrt{\alpha}}x - \frac{\sqrt{\alpha}}{(1+x)^{\gamma-1}}, \quad \frac{dy}{dt} = \frac{-x}{(1+x)^\gamma}. \quad (10)$$

Then we show that the vector field of system (10) is a generalized rotated vector field (see Definition 3.3 of [14]).

**Lemma 5.** *If  $\gamma > 1$  is a fixed integer, then the vector field of system (10) for all real  $\alpha > 0$  is a generalized rotated vector field.*

*Proof.* The unique singular point of system (10) at the origin for all values of the parameter  $\alpha$  remains unchanged and also the value of the following determinant

$$\begin{vmatrix} y + \sqrt{\alpha_1} - \frac{1}{\sqrt{\alpha_1}}x - \frac{\sqrt{\alpha_1}}{(1+x)^{\gamma-1}} & \frac{-x}{(1+x)^\gamma} \\ y + \sqrt{\alpha_2} - \frac{1}{\sqrt{\alpha_2}}x - \frac{\sqrt{\alpha_2}}{(1+x)^{\gamma-1}} & \frac{-x}{(1+x)^\gamma} \end{vmatrix},$$

for all  $x > -1$  and  $y \in \mathbb{R}$ , is equal to

$$(\sqrt{\alpha_2} - \sqrt{\alpha_1}) \left( \frac{x((x+1)^{\gamma-1} - 1)}{(1+x)^{2\gamma-1}} + \frac{1}{\sqrt{\alpha_1\alpha_2}} \frac{x^2}{(1+x)^\gamma} \right) \geq 0,$$

where  $\alpha_1 < \alpha_2$  and  $\alpha_1, \alpha_2 > 0$ . The value of this determinant can only be zero for  $x = 0$ . Therefore, according to Definition 3.3 of [14], the vector field of system (10) is a generalized rotated vector field.  $\square$

**Remark 6.** *Since system (10) has been obtained from system (7) by a linear change of variables, hence all facts about existence and number of limit cycles of the both systems are the same.*

**Lemma 7.** *For each  $\gamma = 3, 4, 5$  and  $6$ , there exists a real number  $\alpha_\gamma^* \in (1/(\gamma-1), \alpha_\gamma]$  such that system (10) has a unique limit cycle when  $\alpha \in (1/(\gamma-1), \alpha_\gamma^*)$  and no periodic solution when  $\alpha \in (\alpha_\gamma^*, \alpha_\gamma]$  where  $\alpha_\gamma$  for  $\gamma = 3, 4, 5, 6$ , has been introduced in Lemma 3.*

*Proof.* According to Lemma 5 the vector field of system (10) is a generalized rotated vector field. Hence, if this system has a limit cycle for a given values of  $\alpha$ , then by Theorem 3.5 of [6] this limit cycle can neither split nor disappear when parameter  $\alpha$  changes monotonically and also the amplitude of the limit cycle will increase or decrease monotonically.

As we know when  $\alpha = 1/(\gamma-1)$  a Hopf bifurcation occurs, therefore for  $\alpha$  slightly larger than  $1/(\gamma-1)$  system (10) has a unique limit cycle. Theorem 3.5 of [6] implies that the amplitude of the limit cycle increases when parameter  $\alpha$  increases. Moreover system (10) has no periodic orbit when  $\alpha \in (-\infty, 1/(\gamma-1)] \cup (\alpha_\gamma, +\infty)$ . Thus for each  $\gamma = 3, 4, 5, 6$ , there exists  $\alpha_\gamma^* \in (1/(\gamma-1), \alpha_\gamma]$  such that when  $\alpha$  tends to  $\alpha_\gamma^*$ , the amplitude of the limit cycle tends to infinity. This completes the proof of the lemma.  $\square$

**2.4. Phase portraits in the Poincaré disc.** By considering all information of system (1) which have been obtained until now, and taking into account that the straight line  $y = 0$  is invariant, we get the phase portraits of Figures 1 and 2 described in Theorems 1 and 2, respectively. Therefore the proofs of Theorems 1 and 2 are completed.



is equal to  $\nu$ . Moreover, if the number of non-zero members of the revised sign list is  $l$ , then,  $l - 2\nu$  is the number of the distinct real roots of polynomial  $f(x)$ .

#### ACKNOWLEDGEMENTS

We thank to the reviewers their comments and suggestions which help us to improve the presentation of this paper.

The first author is partially supported by the Ministerio de Economía, Industria y competitividad, Agencia Estatal de Investigación grant MTM2016-77278-P (FEDER), the Agència de Gestió d'Ajuts Universitaris i de Recerca grant 2017 SGR 1617, and the European project Dynamics-H2020-MSCA-RISE-2017-777911. The second author is supported by Isfahan University of Technology (IUT).

#### REFERENCES

- [1] M. J. ÁLVAREZ, A. FERRAGUT AND X. JARQUE, *A survey on the blow up technique*, Int. J. Bifurcation and Chaos. **21** (2011), 3103–3118.
- [2] J. C. ARTEŚ, J. LLIBRE AND C. VALLS, *Dynamics of the Higgins–Selkov and Selkov system*, Chaos. Sol. Frac. **114** (2018), 145–150.
- [3] P. BRECHMANN AND A. D. RENDALL, *Dynamics of the Selkov oscillator*, Mathematical Biosciences, **306** (2018), 152–159.
- [4] P. BRECHMANN AND A. D. RENDALL, *Unbounded solutions of models for glycolysis*, preprint arXiv (<https://arxiv.org/abs/2003.07140>), (2020).
- [5] H. CHEN, J. LLIBRE AND Y. TANG, *Global dynamics of a SD oscillator*, Nonlinear Dynamics, **91** (2018), 1755–1777.
- [6] H. CHEN AND Y. TANG, *Proof of Artés-Llibre-Valls’s conjectures for the Higgins–Selkov and Selkov systems*, J. Differential Equations **266** (2019), 7638–7657.
- [7] F. DUMORTIER, J. LLIBRE AND J.C. ARTÉS, *Qualitative Theory of Planar Differential Systems*, Springer Verlag, New York, 2006.
- [8] J. GUCKENHEIMER AND P. HOLMES, *Nonlinear oscillations, dynamical systems, and bifurcations of vector fields*, in: *Applied Mathematical Sciences*, vol. 42, Springer–Verlag, New York, 1986.
- [9] J. HIGGINS, *A chemical mechanism for oscillation of glycolytic intermediates in yeast cells*, Proc. Natl. Acad. Sci. (USA) **51** (1964), 989–994.
- [10] T. W. HWANG AND H. J. TSAI, *Uniqueness of limit cycles in theoretical models of certain oscillating chemical reactions*, J. Phys. A **38** (2005), 8211–8223.
- [11] L. PERKO, *Differential Equations and Dynamical Systems*, (3rd Ed., Springer, 2001).
- [12] E. E. SELKOV, *Self-oscillations in glycolysis. I. A simple kinetic model*, Eur. J. Biochem. **4** (1968), 79–86.
- [13] L. YANG, *Recent advances on determining the number of real roots of parametric polynomials*, J. Symb. Comput. **28** (1999), 225–242.
- [14] Z. ZHANG, T. DING AND W. HUANG, *Qualitative Theory of Differential Equations*, Transl. Math. Monogr. Amer. Soc., Providence, RI, 1992.

*Email address:* jllibre@mat.uab.cat

*Email address:* marzieh.mousavi@math.iut.ac.ir