

HEAVY LIQUID CHAMBER ANALYSIS OF  $\eta^0 \rightarrow 3\pi$  AND  $\eta^0 \rightarrow \gamma\gamma$  DECAY-MODESC. BAGLIN, A. BEZAGUET, B. DEGRANGE, P. MUSSET  
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From 199  $\eta^0 \rightarrow 3\pi^0 \rightarrow 6\gamma$ , 526  $\eta^0 \rightarrow \pi^+\pi^-\pi^0$  and 401  $\eta^0 \rightarrow \gamma\gamma$  events with all  $\gamma$ -rays reconstructed in a heavy liquid bubble chamber, we obtain direct and independent measurements of the branching ratios:

$R_1 = \Gamma(\eta^0 \rightarrow 3\pi^0)/\Gamma(\eta^0 \rightarrow \pi^+\pi^-\pi^0) = 1.50^{+0.15}_{-0.29}$  and  $R_2 = \Gamma(\eta^0 \rightarrow \gamma\gamma)/\Gamma(\eta^0 \rightarrow \pi^+\pi^-\pi^0) = 1.72 \pm 0.25$ . We also study the Dalitz-plots for the modes  $\eta^0 \rightarrow 3\pi^0$  and  $\eta^0 \rightarrow \pi^+\pi^-\pi^0$ , and find no significant variation of the amplitude of the decay  $\eta^0 \rightarrow 3\pi^0$ .

Previous experiments on eta meson neutral decay modes have yielded often contradictory results [1]. In most of these experiments, the number of  $\gamma$ -rays actually produced in the observed events remained unknown and the relative decay rates were fitted in order to explain some experimental distribution [2-4]. In contrast, for events selected in the present experiment [15], all the  $\gamma$ -rays from the decays  $\eta^0 \rightarrow 3\pi^0 \rightarrow 6\gamma$ ,  $\eta^0 \rightarrow \pi^+\pi^-\pi^0 (\pi^0 \rightarrow \gamma\gamma)$  and  $\eta^0 \rightarrow \gamma\gamma$  are detected by electron pair production in the Ecole Polytechnique Heavy Liquid Bubble Chamber, and their momenta are measured by the Behr-Mittner method. The branching ratios  $R_1 = \Gamma(\eta^0 \rightarrow 3\pi^0)/(\Gamma(\eta^0 \rightarrow \pi^+\pi^-\pi^0))$  and  $R_2 = \Gamma(\eta^0 \rightarrow \gamma\gamma)/\Gamma(\eta^0 \rightarrow \pi^+\pi^-\pi^0)$  are thus obtained independently of each other and free of background from other possible neutral modes, e.g.  $\pi^0\gamma\gamma$ . The  $\eta^0$  mesons were produced via  $p \rightarrow \eta^0 n$  by 950 MeV/c  $\pi^-$  from the Saclay 2 GeV proton synchrotron. The  $1 \times \frac{1}{2} \times \frac{1}{2} \text{ m}^3$

bubble chamber was filled for a part of the film with pure freon  $\text{CF}_3\text{Br}$  (11 cm radiation length) and for the other part with a mixture of propane  $\text{C}_3\text{H}_8$  and freon  $\text{CF}_3\text{Br}$  (20 cm radiation length). The experimental procedure is described further elsewhere [5].

1)  $R_1 = \Gamma(\eta^0 \rightarrow 3\pi^0)/\Gamma(\eta^0 \rightarrow \pi^+\pi^-\pi^0)$ . The sample of  $\eta^0 \rightarrow \pi^+\pi^-\pi^0$  decays has been obtained as described in ref. 6. For the numerator sample, we retain only events with all 6  $\gamma$ 's converted. The film was scanned twice for 6  $\gamma$ -events with a combined scanning efficiency of about 98% for the events surviving the cuts as described below. Events with 4  $\gamma$ 's or 5  $\gamma$ 's simulating 6  $\gamma$  production by bremsstrahlung of a conversion electron are removed by cut-offs on the angles between pairs of  $\gamma$ -rays [5]. Only 17% of actual  $\eta^0$  events are lost by this requirement which eliminates virtually all bremsstrahlung  $\gamma$ -rays as we have confirmed by inspection of the photographs. Other cuts [5] have been applied to suppress background events containing among the six a spurious  $\gamma$  from another source. The background of non 6  $\gamma$ -events in those surviving our cuts is less than 12% (90% confidence), as estimated by determining the

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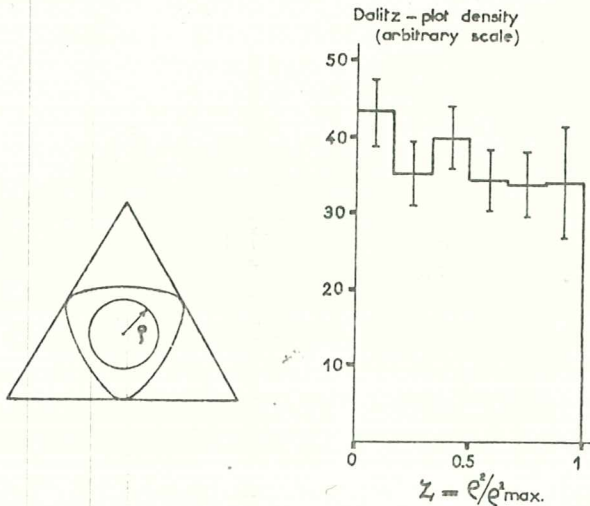


Fig. 1. Dalitz plot apparent density for  $\eta^0 \rightarrow 3\pi^0$  as a function of

$$z = \frac{2}{3} \sum_{i=1}^3 \left[ \frac{3}{M_\eta - 3m_\pi} (E_i - \frac{1}{3}M_\eta) \right]^2$$

By "apparent density", we mean that the actual density is deformed by the measurement errors and by weighting procedure. If  $\alpha'$  is the apparent slope of  $M^2 = A(1+2\alpha z)$ , Monte Carlo calculations show  $\alpha'/\alpha$  is of the order of 0.4.

probability for a given interaction to have a spurious  $\gamma$ -ray pointing to it by chance.

The production of more than  $3\pi$ 's is highly improbable at the energy chosen for the beam. Hence, any contamination of our sample of 6  $\gamma$ -events is due to non eta  $3\pi^0$  production. From direct fits of the observed  $3\pi^0$ -mass distribution to that expected from the incoherent sum of  $3\pi$ -phase space production plus eta production, and taking into account measurement errors, we find that the background of non eta  $3\pi^0$  events is less than 21% (68% confidence). We have checked previously that the estimated errors on  $\gamma$ -ray momenta account for our observed  $\pi^0 \rightarrow \gamma\gamma$  mass and  $\chi^2$  distributions [5].

Using a theoretical model [14], a similar heavy liquid bubble chamber experiment [8] concludes that the production of non eta  $\pi^0\pi^0\pi^0$  events is less favoured than the production of non eta  $\pi^+\pi^-\pi^0$  events, the latter being estimated in our experiment to be  $(13 \pm 3)\%$  of  $\eta^0 \rightarrow \pi^+\pi^-\pi^0$  events from the data of refs. 10 and 11. This order of magnitude is further confirmed by a study in our film of the production of  $\pi^+\pi^-\pi^0$  events.

Correcting for  $\gamma$  detection efficiencies ( $\sim 0.7$  per  $\gamma$  in the freon and  $\sim 0.5$  per  $\gamma$  in the mixture) and for the effects of the cuts, we find

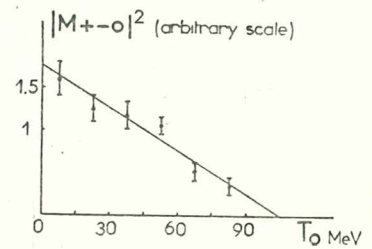


Fig. 2. Dalitz plot density for  $\eta^0 \rightarrow \pi^+\pi^-\pi^0$  as a function of  $T_0$ .

$R_1 = 15.2^{+0.17}_{-0.33}$  in the freon sample, and  $R_1 = 1.45^{+0.29}_{-0.42}$  in the mixture. The agreement between these two results is a check of our estimates of  $\gamma$  detection efficiencies since the ratio of 6  $\gamma$  to 2  $\gamma$  detection efficiencies is about three times greater in the freon than in the mixture. Combining these results we find  $R_1 = 1.50^{+0.15}_{-0.23}$ . The asymmetric error is due mainly to the non eta  $3\pi^0$  background uncertainty. This error cannot be compared directly with that of ref. 8 where the amount of non eta  $3\pi^0$  background is estimated theoretically.

Our final result confirms our preliminary value  $1.3 \pm 0.4$  presented elsewhere [7]. It agrees well with other recent experiments [3,4,8,9].

II)  $R_2 = \Gamma(\eta^0 \rightarrow \gamma\gamma)/\Gamma(\eta^0 \rightarrow \pi^+\pi^-\pi^0)$ . Our measurement of this ratio is based on only a part of our film which contains 401  $\eta^0 \rightarrow \gamma\gamma$  and 214  $\eta^0 \rightarrow \pi^+\pi^-\pi^0$  events with all  $\gamma$ -rays detected. The  $\eta^0 \rightarrow \gamma\gamma$  candidates were required to have the angle between the 2  $\gamma$ -rays projected on the scanning table  $> 90^\circ$ . This criterion eliminates 86% of  $\pi^0 \rightarrow \gamma\gamma$  events.

The main background for  $\eta^0 \rightarrow \gamma\gamma$  events is due to the reaction  $\pi^-\pi^+ \rightarrow \pi^0\pi^0$  in which two photons escape detection. This background is easily estimated and removed using distributions obtained from events with 3  $\gamma$  detected [5]. We then find  $R_2 = 1.72 \pm 0.25$ .

Combining our two independent measurements we find  $R_1 + R_2 = 3.22^{+0.29}_{-0.38}$  compared with the world average [1] deduced from hydrogen bubble chamber experiments  $\Gamma(\eta^0 \rightarrow \text{all neutrals})/\Gamma(\eta^0 \rightarrow \pi^+\pi^-\pi^0) = 3.35 \pm 0.35$ . (We do not use the fitted value given in ref.1 which already takes account of experiments on neutral modes.) Our results are thus compatible with the absence of the mode  $\eta^0 \rightarrow \pi^0\gamma\gamma$ , in agreement with recent experiments on this mode [1].

III A) Dalitz plot analysis for  $\eta^0 \rightarrow 3\pi^0$  decays. The present data on the Dalitz plot of the decay  $\eta^0 \rightarrow \pi^+\pi^-\pi^0$  are well fitted by a matrix element

linear in the kinetic energy of the neutral pion. Upper limits on possible quadratic terms have been reported [12]. Up to now, there has been no investigation of the Dalitz plot of the decay  $\eta^0 \rightarrow 3\pi^0$ , in which the linear term must vanish as a result of the symmetry among the three pions. The matrix element squared can be described to lowest order in the pion energies by a constant plus one quadratic term:  $|M_{000}|^2 = A(1 + 2\alpha z)$  where

$$z = \frac{3}{\sum_{i=1}^3 y_i^2} \text{ and } y_i = 3(E_i - \frac{1}{3}M_\eta)/(M_\eta - 3m_\pi).$$

$z$  varies from 0 to 1 and is proportional to  $\rho^2$ , where  $\rho$  is the distance to the center of the Dalitz plot.

The measurement of all  $\gamma$ -rays produced in the decay  $\eta^0 \rightarrow 3\pi^0 \rightarrow 6\gamma$  allows the reconstruction of the  $\pi^0$ 's and a study of the Dalitz plot for this decay. A priori, there are 15 ways of pairing 6 $\gamma$ 's into 3 $\pi^0$ 's. For each of these 15 hypotheses we have fitted the  $\gamma$ -ray momenta and directions to the four constraints that each of the 3 pairs of  $\gamma$ -rays must have the  $\pi^0$  mass and the six  $\gamma$ -rays must have the  $\eta^0$  mass. Each hypothesis is then given a weight proportional to its likelihood, calculated as indicated in ref. 5. Typically some three hypotheses make a significant contribution.

From our sample of 199  $\eta^0 \rightarrow 3\pi^0 \rightarrow 6\gamma$  candidates, we find  $\alpha = -0.32 \pm 0.37$  or  $-0.86 < \alpha < 0.2$  with 90% confidence. The confidence limits take account of the effect of the measurement errors on  $\gamma$ -ray momenta (calibrated by a study of observed  $\pi^0 \rightarrow \gamma\gamma$  distributions [5]), and of the effects of the fitting and of the weighting procedures. This result is compatible with a constant matrix element (fig. 1).

III B) Dalitz plot for  $\eta^0 \rightarrow \pi^+\pi^-\pi^0$ . For these events the  $\pi^0$  kinetic energy  $T_0$  in the  $\eta^0$  rest-frame is reconstructed with an average precision of 5 MeV in the two constraint unambiguous fit  $\eta^0 \rightarrow \pi^+\pi^-\pi^0 \rightarrow \pi^+\pi^-\gamma\gamma$ . Small corrections have been made to the  $T_0$  spectrum to account for ef-

fects of our cuts and to remove 13% non-resonant  $\pi^+\pi^-\pi^0$  background. These corrections vary by less than 8% across the  $T_0$  spectrum. The resulting spectrum can be fitted well by a matrix element squared linear in  $T_0$

$$|M_{+-0}|^2 = B \left[ 1 - 2a \left( \frac{3T_0}{m_\eta - 3m_\pi} - 1 \right) \right]$$

with:

$a = 0.41 \pm 0.04$  (fig. 2) in agreement with previous results [11-13].

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#### References

1. N. Barash-Schmidt et al., Rev. Mod. Phys. 41 (1969) 109 contains extensive references to experimental literature. For a recent review, see C. Baltay in Meson spectroscopy, eds. C. Baltay and A. H. Rosenfeld (Benjamin, New York, 1968) p. 95.
2. G. di Giugno et al., Phys. Rev. Letters 16 (1966) 767; M. Feldman et al., Phys. Rev. Letters 18 (1967) 868.
3. S. Buniatov et al., Phys. Letters 25B (1967) 560.
4. C. Baltay et al., Phys. Rev. Letters 19 (1967) 1495.
5. C. Baglin et al., to be published.
6. C. Baglin et al., Phys. Letters 22 (1966) 219 and 24B (1967) 637.
7. C. Baglin et al., Bull. Am. Phys. Soc. 12 (1967) 567.
8. F. W. Bullock et al., Phys. Letters 27B (1968) 402.
9. R. J. Cence et al., Phys. Rev. Letters 19 (1967) 1393.
10. C. Baltay et al., Phys. Rev. Letters 26 (1966) 1224.
11. A. Larribe et al., Phys. Letters 23 (1966) 600.
12. A. M. Cnops et al., Phys. Letters 27B (1968) 113.
13. Columbia - Berkeley - Purdue - Wisconsin - Yale collaboration, Phys. Rev. 149 (1966) 1044.
14. K. Zalewski and J. A. Danysz, Nucl. Phys. B2 (1967) 249.
15. The work on part I of this letter was done mainly in Paris and Berkeley, part II in Paris, Madrid and Strasbourg, part III in Paris and Strasbourg. The value of  $R_1$  given here supersedes those given in W. Michael, U.C. Berkeley Thesis (1967), and in G. Irwin et al., Bull. Am. Phys. Soc. 13 (1968) 1641.

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